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## Data Sheet

Journal:
Report Ref.: RSUQ-2021-004
Arxiv Ref.: https://arxiv.org/abs/2107.00394 [stat.CO]

DOI:
Date submitted: July 1st, 2021

Date accepted:

# Global sensitivity analysis using derivative-based sparse Poincaré chaos expansions 

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July 16, 2021


#### Abstract

Variance-based global sensitivity analysis, in particular Sobol' analysis, is widely used for determining the importance of input variables to a computational model. Sobol' indices can be computed cheaply based on spectral methods like polynomial chaos expansions (PCE). Another choice are the recently developed Poincaré chaos expansions (PoinCE), whose orthonormal tensor-product basis is generated from the eigenfunctions of one-dimensional Poincaré differential operators. In this paper, we show that the Poincaré basis is the unique orthonormal basis with the property that partial derivatives of the basis form again an orthogonal basis with respect to the same measure as the original basis. This special property makes PoinCE ideally suited for incorporating derivative information into the surrogate modelling process. Assuming that partial derivative evaluations of the computational model are available, we compute spectral expansions in terms of Poincaré basis functions or basis partial derivatives, respectively, by sparse regression. We show on two numerical examples that the derivative-based expansions provide accurate estimates for Sobol' indices, even outperforming PCE in terms of bias and variance. In addition, we derive an analytical expression based on the PoinCE coefficients for a second popular sensitivity index, the derivative-based sensitivity measure (DGSM), and explore its performance as upper bound to the corresponding total Sobol' indices.


## 1 Introduction

Computer models simulating physical phenomena and industrial systems are commonly used in engineering and safety studies, for prediction, validation or optimisation purposes. These numerical models often take as inputs a high number of physical parameters, whose values are variable or not perfectly known, creating the need for uncertainty quantification on model
computations (Smith, 2014). Uncertainty quantification typically becomes more challenging the higher the input dimension is (curse of dimensionality). In this situation, global sensitivity analysis (GSA) is an invaluable tool that allows the analyst to rank the relative importance of each input of the model and to detect non-influential inputs (Borgonovo and Plischke, 2016; Razavi et al., 2021). Most often relying on a probabilistic modeling of the model input variables, GSA tries to explain model output uncertainties on the basis of model input uncertainties, accounting for the full range of variation of the variables.

A well-known and widely used GSA method is Sobol' analysis (Sobol', 1993), which relies on the functional ANOVA (analysis of variance) decomposition (Efron and Stein, 1981). For a squareintegrable model and independent input variables, Sobol' analysis determines which part of the model output variance can be attributed to each input and to each interaction between inputs. The overall contribution of each input, including interactions with other inputs, is provided by the total Sobol' index (Homma and Saltelli, 1996). Sobol' indices can be estimated efficiently using various Monte Carlo-based techniques as well as metamodel-based techniques (Prieur and Tarantola, 2017). The latter save on expensive model evaluations by first performing a small number of model runs, which are used to compute an accurate approximation to the original model - the meta- or surrogate model - from which the Sobol' indices are finally computed (Fang et al., 2006; Le Gratiet et al., 2017).
One of the most popular and powerful metamodelling methods is the polynomial chaos expansion (PCE) (Xiu and Karniadakis, 2002). PCE represents the model in a specific basis consisting of polynomials that are orthonormal with respect to the input distribution. Orthogonal polynomial systems have been studied throughout the last century and they have many useful properties (see, e.g., Szegö (1939) and Simon (2010)). One particular strength of PCE is that once it is computed, it easily gives all the variance-based quantities defined through the ANOVA decomposition, and in particular the Sobol' indices at all orders (Sudret, 2008). In practice, the expansion cannot use infinitely many terms and must be truncated. Among the many approaches available to compute the expansion coefficients, sparse regression techniques combined with adaptive basis selection appear to be especially promising (see Lüthen et al. (2021b,a) for an overview). Here, a small number of terms is selected which is able to best represent the computational model based on the available model evaluations.

In some practical situations, partial derivatives of the model output with respect to each input are easily accessible, for example by algorithmic differentiation of the numerical model in the reverse (adjoint) mode (Griewank and Walther, 2008). This technique allows for computing all partial derivatives of the model output at a cost independent of the number of input variables. Since PCEs are such a well-established metamodelling tool, there have been many efforts to leverage the additional information contained in model derivatives to improve the performance of PCE. The idea of including derivative information into sparse regression problems, often called gradient-enhanced $\ell^{1}$-minimization, is tested by Jakeman et al. (2015) for one numerical example with uniform inputs, and analyzed theoretically and numerically by Peng et al. (2016) for Hermite PCE. Both report favorable results. Roderick et al. (2010) and Li et al. (2011) apply polynomial regression (PCE) in the context of nuclear engineering. They include derivative
information into the least-squares regression formulation and observe that most polynomial families are not orthogonal with respect to the $H^{1}$ inner product, which destroys the orthogonality of the regression matrix. To alleviate this issue, Guo et al. (2018) develop a preconditioning procedure for gradient-enhanced sparse regression with certain polynomial families, with the goal of improving the orthogonality properties of the regression matrix. In all these approaches, the inclusion of derivative information is not straightforward, but requires specific polynomial families (in particular Hermite) or elaborate techniques like preconditioning, because the partial derivatives of a PCE basis do in general not form an orthogonal system. Gejadzea et al. (2019) have derived derivative-enhanced projection methods to compute the PCE coefficients but, as usual with projection methods, only a small maximal degree of the PCE can be reached.

On a different note, the availability of model derivatives has implications also for GSA. The so-called Derivative-based Global Sensitivity Measures (DGSM) are computed by integrating the squared partial derivatives of the model output over the domain of the inputs. These indices have been shown to be efficiently estimated by sampling techniques (as Monte Carlo or quasi-Monte Carlo) as well as from PCE (Sudret and Mai, 2015), and have been proven to be an excellent screening technique (i.e., detecting all the non-influential inputs among a large number), see e.g. the review in Kucherenko and Iooss (2017). Indeed, the interpretation of DGSM indices is straightforward due to their inequality relationship with Sobol' indices: multiplied with the associated Poincaré constant, DGSM indices provide an upper bound of the total Sobol' index (Roustant et al., 2017), regardless of the input probability distribution.
Another way to utilize model derivatives, which solves the issues present for the polynomial chaos formulation, and naturally provides sharp lower bounds as well as upper bounds on total Sobol' indices, is to compute Poincaré chaos expansions (Roustant et al., 2020), which we will abbreviate by PoinCE in the sequel. Similar to PCE, PoinCE is a spectral expansion in terms of an orthonormal basis whose elements are eigenfunctions of the so-called Poincaré differential operator. The eigenfunctions are in general non-polynomial, except for the special case of the Gaussian distribution, where they coincide with the Hermite polynomials. The key property of PoinCE is that the partial derivatives of the basis form again an orthogonal basis with respect to the input distribution. This allows to conveniently expand the derivative of the computational model in terms of partial derivatives of the basis (PoinCE-der), which yields another estimator for partial variances and Sobol' indices. If the partial derivatives of the model have smaller variability than the model itself, the estimates based on the model derivatives might be more accurate. This makes PoinCE(-der) an efficient tool for screening (Roustant et al., 2020).

Our present contribution to the field of generalized chaos expansions and GSA is two-fold. On the theoretical side, we provide a proof that the Poincaré basis is in fact characterized uniquely as the orthonormal basis which remains an orthogonal basis (w.r.t. the same probability measure) after differentiation. Furthermore, we show how PoinCE naturally generalizes an analytical formula for DGSM originally developed for Hermite PCE (Sudret and Mai, 2015), which implies that PoinCE simultaneously and efficiently provides lower and upper bounds to all partial variances. On the computational side, we improve on Roustant et al. (2020), which introduced projectionbased Poincaré chaos and demonstrated that small Sobol' indices were approximated particularly
well by the derivative expansion. In this contribution, we compute PoinCE by sparse regression, thus generalizing the powerful and cost-effective sparse PCE methodology to non-polynomial functions. We explore the performance of PoinCE as an estimator for partial variances (upper and lower bounds) and compare it to standard PCE.

This paper is organized as follows. Section 2 revisits the mathematical foundations of PoinCE and presents several analytical results related to Sobol' indices and DGSM. Section 3 explains the computation of PoinCE basis functions, and the sparse regression methodology adapted from PCE to PoinCE. The methodology is applied in Section 4, where two example problems are investigated to demonstrate its performance for sensitivity analysis and screening. Finally, we summarize our conclusions in Section 5.

## 2 Mathematical background

### 2.1 Orthonormal bases in $L^{2}$

In this section, we recall some important facts about orthonormal bases in $L^{2}(E, \mu)$ where $E$ is a subset of $\mathbb{R}^{d}$ and $\mu$ is a probability measure on $E$. We first outline the general theory in Section 2.1.1. The particular cases of polynomial and Poincaré bases in several dimensions are developed in Sections 2.1.2 and 2.1.3.

### 2.1.1 General theory

To begin with, recall that $L^{2}(E, \mu)$ endowed with the inner product

$$
\begin{equation*}
\langle f, g\rangle=\int_{E} f(x) g(x) \mu(d x), \quad \text { for } f, g \in L^{2}(E, \mu) \tag{1}
\end{equation*}
$$

is a Hilbert space. Recall that a sequence of functions $\left(\Phi_{\alpha}\right)_{\alpha \in \mathcal{I}}$ is an orthonormal system in $L^{2}(E, \mu)$ if it satisfies the two following assumptions:

1) For all $\alpha \neq \alpha^{\prime},\left\langle\Phi_{\alpha}, \Phi_{\alpha^{\prime}}\right\rangle=0$, (orthogonality)
2) For all $\alpha,\left\langle\Phi_{\alpha}, \Phi_{\alpha}\right\rangle=1$ (unit norm).

An orthonormal system in $L^{2}(E, \mu)$ is called complete if the closure of the span generated by $\left(\Phi_{\alpha}\right)$ is $L^{2}(E, \mu)$. In this case, the system $\left(\Phi_{\alpha}\right)$ is called an Hilbertian or orthonormal basis of $L^{2}(E, \mu)$ and for any function $f \in L^{2}(E, \mu)$ the following expansion holds:

$$
\begin{equation*}
f=\sum_{\alpha}\left\langle\Phi_{\alpha}, f\right\rangle \Phi_{\alpha}, \quad(\mu \text { almost surely }) . \tag{2}
\end{equation*}
$$

When used to represent random variables in terms of a basis of uncorrelated random variables, such an expansion is often called chaos expansion in the uncertainty quantification literature (Wiener, 1938; Ghanem and Spanos, 1991; Ernst et al., 2012).

An archetype example of chaos expansion is given by the so-called Fourier expansion. This corresponds to the case where the set $E=[0,1]$ is endowed with the Lebesgue measure and we have for $\alpha \in \mathbb{Z}$,

$$
\Phi_{\alpha}(x)=\sqrt{2} \cos (2 \pi \alpha x) \text { if } \alpha<0, \Phi_{0}(x)=1, \text { and } \Phi_{\alpha}(x)=\sqrt{2} \sin (2 \pi \alpha x) \text { if } \alpha>0 .
$$

In this frame, any square-integrable function $f$ may be expanded as

$$
f(x)=a_{0}+\sqrt{2} \sum_{\alpha>0}\left(a_{\alpha} \cos (2 \pi \alpha x)+b_{\alpha} \sin (2 \pi \alpha x)\right) .
$$

Here, for all $\alpha \in \mathbb{Z}_{*}$,

$$
a_{0}=\int_{0}^{1} f(x) d x, \quad a_{\alpha}=\sqrt{2} \int_{0}^{1} f(x) \cos (2 \pi \alpha) d x, b_{\alpha}=\sqrt{2} \int_{0}^{1} f(x) \sin (2 \pi \alpha) d x
$$

In the following we describe two particular chaos types, namely the classical polynomial chaos and the recently developed Poincaré chaos.

### 2.1.2 Polynomial chaos

A classical family of chaos expansions on an interval $E$ of $\mathbb{R}$ endowed with a probability measure are polynomial chaos expansions (PCE) given by orthonormal polynomial bases. A well-known example is the Hermite expansion for which the set $E$ is the whole line $\mathbb{R}$ endowed with the standard Gaussian distribution. In this example, for $\alpha \in \mathbb{N}, \Phi_{\alpha}=H_{\alpha}$ is the Hermite polynomial of degree $\alpha$. The first Hermite polynomials are

$$
H_{0}(x)=1, H_{1}(x)=x, H_{2}(x)=\frac{x^{2}-1}{2}, H_{3}(x)=\frac{x^{3}-3 x}{6}, H_{4}(x)=\frac{x^{4}-6 x^{2}+3}{24} \quad(x \in \mathbb{R}) .
$$

When considering a product of intervals $E=E_{1} \times E_{2} \times \cdots \times E_{d}$ endowed with a product probability measure $\mu=\mu_{1} \otimes \cdots \otimes \mu_{d}$ there is a canonical way to build a Hilbertian basis from a collection of marginal Hilbertian bases. Indeed, for $i=1, \ldots, d$ assume that $\left(\Phi_{\alpha_{i}}^{(i)}\right)$ is a Hilbertian basis of $L^{2}\left(E_{i}, \mu_{i}\right)$. Then, setting $\boldsymbol{\alpha}:=\left(\alpha_{1}, \ldots, \alpha_{d}\right)$ and $\Phi_{\boldsymbol{\alpha}}:=\prod_{i=1}^{d} \Phi_{\alpha_{i}}^{(i)}$ we obtain that $\left(\Phi_{\boldsymbol{\alpha}}\right)$ is an orthonormal basis of $L^{2}(E, \mu)$.

### 2.1.3 Poincaré chaos

The Poincaré basis is another example of an orthonormal basis of $L^{2}(\mu)$, consisting of functions that admit weak derivatives, i.e. that belong to $H^{1}(\mu)=\left\{f \in L^{2}(\mu)\right.$ s.t. $\left.f^{\prime} \in L^{2}(\mu)\right\}$. Recall that $H^{1}(\mu)$, endowed with the norm $\|f\|_{H^{1}(\mu)}^{2}=\|f\|^{2}+\left\|f^{\prime}\right\|^{2}$, is a Hilbert space. Like in the case of polynomial chaos, the multivariate Poincaré basis is obtained by tensorization of 1D orthonormal bases. Thus it is enough to describe the 1D case, and we now fix $d=1$. This short summary is based on Roustant et al. (2017) in which more details can be found. We assume that:

Assumption 1. The probability measure $\mu$ is supported on a bounded interval ( $a, b$ ) and admits a density of the form $\rho=e^{-V}$, where $V$ is continuous and piecewise $C^{1}$ on $[a, b]$ with respect to the Lebesgue measure.

This assumption is sufficient to guarantee the existence of a Poincaré basis. On the topological side, it implies that the Hilbert space $L^{2}(\mu)\left(\right.$ resp. $\left.H^{1}(\mu)\right)$ is equal to $L^{2}(a, b)\left(\right.$ resp. $\left.H^{1}(a, b)\right)$, with an equivalent norm. Indeed, $\mu$ is a bounded perturbation of the uniform measure on $[a, b]$, meaning that the pdf $\rho$ is bounded from below and above by strictly positive constants (by continuity of $V$ on the compact support $[a, b]$ ).

Theorem 1 (1D Poincaré basis). Under Assumption 1, there exists an orthonormal basis $\left(\varphi_{\alpha}\right)_{\alpha \geq 0}$ of $L^{2}(\mu)$ such that for all $f \in H^{1}(\mu)$ and for all integer $\alpha \geq 0$, we have:

$$
\begin{equation*}
\left\langle f^{\prime}, \varphi_{\alpha}^{\prime}\right\rangle=\lambda_{\alpha}\left\langle f, \varphi_{\alpha}\right\rangle, \tag{3}
\end{equation*}
$$

where $\left(\lambda_{\alpha}\right)_{\alpha \geq 0}$ is an increasing sequence that tends to infinity:

$$
0=\lambda_{0}<\lambda_{1}<\lambda_{2}<\cdots<\lambda_{\alpha} \underset{\alpha \rightarrow \infty}{\longrightarrow}+\infty .
$$

Here, the inner product $\langle\cdot, \cdot\rangle$ is the one on $L^{2}(\mu)$ as defined in (1). The basis functions $\varphi_{\alpha}$ are unique up to a sign change, and form the so-called Poincaré basis. Notice that $\varphi_{0}$ is a constant function equal to $\pm 1$; by convention, we choose $\varphi_{0}=1$.
Furthermore, the Poincaré basis functions are the eigenfunctions of the differential operator

$$
L(f)=f^{\prime \prime}-V^{\prime} f^{\prime}
$$

i.e. satisfy $L(f)=-\lambda f$, subject to Neumann conditions $f^{\prime}(a)=f^{\prime}(b)=0$. The $\left(\lambda_{\alpha}\right)_{\alpha \geq 0}$ are the corresponding eigenvalues. Finally, $\alpha^{-2} \lambda_{\alpha} \rightarrow \pi^{2}$ when $\alpha$ tends to infinity, and for all $\alpha \in \mathbb{N}^{\star}$, the eigenfunction $\varphi_{\alpha}$ has exactly $\alpha$ zeros in $(a, b)$.

Proof. The main part of the Theorem can be found in Roustant et al. (2017) or Bakry et al. (2014). The two last assertions come by rewriting the differential equation $f^{\prime \prime}-V^{\prime} f^{\prime}=-\lambda f$ in the Sturm-Liouville form

$$
-\left(p f^{\prime}\right)^{\prime}+q f=\lambda w f
$$

with $p=w=\rho$ and $q=0$. Then, by the Sturm-Liouville theory (see e.g. Zettl (2010, Theorem 4.3.1, (1), (6) and (7))), we have that $\alpha^{-2} \lambda_{\alpha} \rightarrow \pi^{2}$ when $\alpha$ tends to infinity, and for all $\alpha \in \mathbb{N}^{\star}$, the eigenfunction $\varphi_{\alpha}$ has exactly $\alpha$ zeros in $(a, b)$.

The Poincaré basis shares some similarity with both the polynomial chaos and the Fourier basis in terms of oscillations: by Theorem 1, the higher the order of the eigenvalue, the more oscillating the corresponding eigenfunction.
For some specific cases, the Poincaré basis is known analytically. For instance, for the uniform distribution, the Poincaré basis is a kind of Fourier basis (see e.g. Roustant et al. (2020, §4)). Otherwise it has to be computed numerically, e.g., by a finite element technique (see Section 3.1).

Note that Assumption 1 is a convenient sufficient condition which guarantees the existence of a Poincaré basis. It is satisfied for a large range of truncated probability distributions. The set of probability distributions for which the Poincaré chaos exists is larger, but not well known. For instance, the Poincaré chaos is defined for the Gaussian distribution, and then coincides with
polynomial chaos, corresponding to Hermite polynomials. This is the only case where Poincaré chaos and polynomial chaos coincide (Bakry et al., 2014, §2.7). On the other hand, Poincaré chaos is not defined for the Laplace distribution, since the eigenvalues of the associated operator do not form a countable set (Bakry et al., 2014, §4.4.1).

The Poincaré basis is useful for sensitivity analysis. First, it is linked to the Poincaré inequality

$$
\begin{equation*}
\operatorname{Var}_{\mu}(f) \leq C_{P}(\mu) \int f^{\prime 2} d \mu \tag{4}
\end{equation*}
$$

which holds for all functions $f \in H^{1}(\mu)$ under the assumptions on $\mu$. Indeed, the smallest constant $C_{P}(\mu)$ such that (4) is satisfied is equal to $C_{P}(\mu)=1 / \lambda_{1}$, and choosing $f=\varphi_{1}$ corresponds to the equality case (Roustant et al., 2017). Roughly speaking, the Poincaré basis function associated to the first non-zero eigenvalue gives the best "control" of the variability of $f$ by its derivative (in the $L^{2}$ sense). A second appealing property for the analysis of variance is that the derivatives of the Poincaré basis remain orthogonal functions:

Proposition 1. Under Assumption 1, the sequence $\left(\frac{1}{\sqrt{\lambda_{\alpha}}} \varphi_{\alpha}^{\prime}\right)_{\alpha \geq 1}$ is an orthonormal basis of $L^{2}(\mu)$.

Proof. The orthonormality of the sequence is a consequence of (3) by choosing $f=\varphi_{\beta}$, with $\beta \in \mathbb{N}^{*}$. It remains to show that the system is dense in $L^{2}(\mu)$, or equivalently, that its orthogonal is null. Let thus $f \in L^{2}(\mu)$ such that

$$
\left\langle f, \varphi_{\alpha}^{\prime}\right\rangle=0, \quad \text { for all } \alpha \geq 1
$$

As explained when stating Assumption $1, L^{2}(\mu)\left(\right.$ resp. $\left.H^{1}(\mu)\right)$ is equal to $L^{2}(a, b)\left(\right.$ resp. $\left.H^{1}(a, b)\right)$, with an equivalent norm. Now, there exists $g \in H^{1}(\mu)$ such that $f=g^{\prime}$. Indeed, let us define $g$ by $g(x)=g(a)+\int_{a}^{x} f(t) d t$. As $f \in L^{2}(\mu)=L^{2}(a, b)$, then $g$ belongs to $H^{1}(a, b)=H^{1}(\mu)$, and $g^{\prime}=f$. Then we have

$$
\left\langle g^{\prime}, \varphi_{\alpha}^{\prime}\right\rangle=0, \quad \text { for all } \alpha \geq 1
$$

By (3), we obtain $\left\langle g, \varphi_{\alpha}\right\rangle=0$ for all $\alpha \geq 1$ (as $\lambda_{\alpha}>0$ for $\alpha \geq 1$ ). As the functions $\varphi_{\alpha}$ form an orthonormal basis of $L^{2}(\mu)$ with $\varphi_{0}=1$, this implies that $g$ is a constant function, and finally $f=0$. The proof is completed.

In fact, the property in Proposition 1, i.e., that the derivatives of the Poincaré basis form again an orthogonal basis in $L^{2}(\mu)$, uniquely characterizes the Poincaré basis:

Proposition 2. Under Assumption 1, Poincaré bases are the only orthonormal bases $\left(\varphi_{\alpha}\right)$ of $L^{2}(\mu)$ in $H^{1}(\mu)$ such that $\left(\varphi_{\alpha}^{\prime}\right)$ is an orthogonal basis of $L^{2}(\mu)$.

This result seems difficult to find in the literature. German-speaking readers can find a similar proposition in Mikolas (1955), stated in the frame of Sturm-Liouville theory for twicedifferentiable functions satisfying boundary conditions. We provide below a proof based on Hilbertian arguments.

As a corollary, if there exists a basis different from the Poincaré basis for which derivatives form an orthogonal system, then that system is not dense in $L^{2}(\mu)$. As an example, for the uniform probability measure on $[0,2 \pi]$, consider the usual Fourier basis formed by $\{\cos (n x), \sin (n x): n \geq$ $0\}$ (up to multiplicative constants). Taking derivatives results in the same set of functions (up to multiplicative constants) - except for the constant function $\cos (0 x)=1$. Thus, the derivatives form an orthogonal system which covers the orthogonal of constant functions in $L^{2}(\mu)$, which is a strict subspace of $L^{2}(\mu)$. Meanwhile, the Poincaré basis for this probability measure is formed by functions proportional to $\cos \left(\frac{n}{2} x\right)$ for $n \geq 0$. Proposition 2 guarantees that all functions of $L^{2}(\mu)$, including the constant functions, are spanned by the derivatives. Indeed, this is explained intuitively by the presence of half-frequencies: when $n$ is odd, the functions $\cos \left(\frac{n}{2} x\right)$ are not orthogonal to 1 .

Proof of Proposition 2. The fact that a Poincaré basis remains an orthogonal basis by derivation has been proved in Proposition 1. Conversely, let $\left(\varphi_{\alpha}\right)_{\alpha \geq 0}$ be a system of $H^{1}(\mu)$, with $\varphi_{0}=1$, such that $\left(\varphi_{\alpha}\right)$ is an orthonormal basis of $L^{2}(\mu)$ and $\left(\varphi_{\alpha}^{\prime}\right)_{\alpha \geq 1}$ is an orthogonal basis of $L^{2}(\mu)$. Let us first prove that $\left(\varphi_{\alpha}\right)$ is an orthogonal basis of $H^{1}(\mu)$. The orthogonality is a direct consequence of the definition of the inner product of $H^{1}(\mu)$ :

$$
\left\langle\varphi_{\alpha}, \varphi_{\beta}\right\rangle_{H^{1}(\mu)}=\left\langle\varphi_{\alpha}, \varphi_{\beta}\right\rangle_{L^{2}(\mu)}+\left\langle\varphi_{\alpha}^{\prime}, \varphi_{\beta}^{\prime}\right\rangle_{L^{2}(\mu)}=\left(1+\left\|\varphi_{\alpha}^{\prime}\right\|_{L^{2}(\mu)}^{2}\right) \delta_{\alpha, \beta} .
$$

Let us prove that $\left(\varphi_{\alpha}\right)$ is dense in $H^{1}(\mu)$. As explained when stating Assumption $1, L^{2}(\mu)$ (resp. $H^{1}(\mu)$ ) is equal to $L^{2}(a, b)$ (resp. $H^{1}(a, b)$ ), with an equivalent norm. Hence, it is equivalent to prove that $\left(\varphi_{\alpha}\right)$ is dense in $H^{1}(a, b)$. Now, let $f$ be in $H^{1}(a, b)$. As $\left(\varphi_{\alpha}^{\prime}\right)$ is dense in $L^{2}(a, b)$ (equivalently in $L^{2}(\mu)$ ), then $f^{\prime}$ expands as $f^{\prime}=\sum_{\alpha \in \mathbb{N}} c_{\alpha} \varphi_{\alpha}^{\prime}$. In $H^{1}(a, b)$ each function is equal to the primitive of its derivative, hence we have:

$$
\begin{aligned}
\left|f(t)-f(a)-\sum_{\alpha=1}^{N} c_{\alpha}\left(\varphi_{\alpha}(t)-\varphi_{\alpha}(a)\right)\right| & =\left|\int_{a}^{t}\left(f^{\prime}(x)-\sum_{\alpha=1}^{N} c_{\alpha} \varphi_{\alpha}^{\prime}(x)\right) d x\right| \\
& \leq(b-a)\left\|f^{\prime}-\sum_{\alpha=1}^{N} c_{\alpha} \varphi_{\alpha}^{\prime}\right\|_{L^{2}(a, b)}
\end{aligned}
$$

where the inequality comes from the Cauchy-Schwarz inequality. We deduce that $\| f-f(a)-$ $\sum_{\alpha=1}^{N} c_{\alpha}\left(\varphi_{\alpha}-\varphi_{\alpha}(a)\right) \|_{L^{2}(a, b)} \rightarrow 0$ when $N$ tends to infinity. Together with $f^{\prime}=\sum_{\alpha \in \mathbb{N}} c_{\alpha} \phi_{\alpha}^{\prime}$, this implies that $\left\|f-f(a)-\sum_{\alpha=1}^{N} c_{\alpha}\left(\varphi_{\alpha}-\varphi_{\alpha}(a)\right)\right\|_{H^{1}(a, b)} \rightarrow 0$. As $\varphi_{0}=1$, this proves that $\left(\varphi_{\alpha}\right)$ is dense in $H^{1}(a, b)$, which was to be proved.
Now, let us fix $\alpha \geq 0$. Consider the linear form $L_{\alpha}$ defined on $H^{1}(\mu)$ by $L_{\alpha}(f)=\left\langle f^{\prime}, \varphi_{\alpha}^{\prime}\right\rangle_{L^{2}(\mu)}$. The Cauchy-Schwarz inequality gives $\left|L_{\alpha}(f)\right| \leq\left\|f^{\prime}\right\|_{L^{2}(\mu)}\left\|\varphi_{\alpha}^{\prime}\right\|_{L^{2}(\mu)} \leq\|f\|_{H^{1}(\mu)}\left\|\varphi_{\alpha}^{\prime}\right\|_{L^{2}(\mu)}$. This proves that $L_{\alpha}$ is continuous. Hence, by the Riesz representation theorem, there exists a unique $\zeta_{\alpha} \in H^{1}(\mu)$ such that for all $f \in H^{1}(\mu), L_{\alpha}(f)=\left\langle f, \zeta_{\alpha}\right\rangle_{H^{1}(\mu)}$, i.e. $\left\langle f^{\prime}, \varphi_{\alpha}^{\prime}\right\rangle_{L^{2}(\mu)}=\left\langle f, \zeta_{\alpha}\right\rangle_{H^{1}(\mu)}$. Choosing $f=\varphi_{\beta}$ with $\beta \neq \alpha$, we obtain by orthogonality of $\left(\varphi_{\alpha}^{\prime}\right)$ that for all $\beta \neq \alpha$, $\left\langle\varphi_{\beta}, \zeta_{\alpha}\right\rangle_{H^{1}(\mu)}=0$. As $\left(\varphi_{\beta}\right)_{\beta \geq 0}$ is an orthogonal basis of $H^{1}(\mu)$, this implies that $\zeta_{\alpha}$ is collinear to $\varphi_{\alpha}$, i.e., there exists $\tilde{\lambda}_{\alpha} \in \mathbb{R}$ such that $\zeta_{\alpha}=\tilde{\lambda}_{\alpha} \varphi_{\alpha}$. Thus, for all $f \in H^{1}(\mu)$, we have $\left\langle f^{\prime}, \varphi_{\alpha}^{\prime}\right\rangle_{L^{2}(\mu)}=\tilde{\lambda}_{\alpha}\left\langle f, \varphi_{\alpha}\right\rangle_{H^{1}(\mu)}$. Choosing $f=\varphi_{\alpha}$, we get $\tilde{\lambda}_{\alpha}=\frac{\left\|\varphi_{\alpha}^{\prime}\right\|_{L^{2}(\mu)}^{2}}{1+\left\|\varphi_{\alpha}^{\prime}\right\|_{L^{2}(\mu)}^{2}}$, which belongs to
$[0,1)$. Finally, we obtain that $\left\langle f^{\prime}, \varphi_{\alpha}^{\prime}\right\rangle_{L^{2}(\mu)}=\lambda_{\alpha}\left\langle f, \varphi_{\alpha}\right\rangle_{L^{2}(\mu)}$, where $\lambda_{\alpha}=\frac{\tilde{\lambda}_{\alpha}}{1-\lambda_{\alpha}}$ is a non-negative real number. As it is true for all $f$ in $H^{1}(\mu)$ and all $\alpha \in \mathbb{N}$, this implies, by uniqueness of the Poincaré basis (under Assumption 1), that $\left(\varphi_{\alpha}\right)_{\alpha \geq 0}$ is a Poincaré basis.

In higher dimensions, we assume that for all $i=1, \ldots, d$, the probability measure $\mu_{i}$ satisfies Assumption 1, and we denote by $\left(\varphi_{i, \alpha_{i}}\right)_{\alpha_{i} \geq 0}$ the sequence of 1-dimensional Poincaré basis functions, and by $\left(\lambda_{i, \alpha_{i}}\right)_{\alpha_{i} \geq 0}$ the sequence of associated eigenvalues. The Poincaré chaos basis is then defined by the tensor product $\Phi_{\alpha}=\varphi_{1, \alpha_{1}} \otimes \cdots \otimes \varphi_{d, \alpha_{d}}$. Using the properties of $L^{2}$ bases, (3) thus implies that for all $f \in H^{1}(\mu)$, for all $i=1, \ldots, d$ :

$$
\begin{equation*}
\left\langle\frac{\partial f}{\partial x_{i}}, \frac{\partial \Phi_{\alpha}}{\partial x_{i}}\right\rangle=\lambda_{i, \alpha_{i}}\left\langle f, \Phi_{\alpha}\right\rangle . \tag{5}
\end{equation*}
$$

Similarly, applying Proposition 1, we get:
Proposition 3. Under Assumption 1, for all $i=1, \ldots, d$, the sequence $\left(\frac{1}{\sqrt{\lambda_{i, \alpha_{i}}}} \frac{\partial \Phi_{\alpha}}{\partial x_{i}}\right)_{\alpha, \alpha_{i} \geq 1}$ is an orthonormal basis of $L^{2}(\mu)$.

### 2.2 Variance-based indices, derivative-based indices

We first recall the definition of variance-based sensitivity indices, which quantify the importance of each input variable in terms of function response variability.

Let $f$ be a real-valued function defined on $E=E_{1} \times \cdots \times E_{d} \subseteq \mathbb{R}^{d}$. The uncertainty of the inputs is represented by a random vector $\boldsymbol{X}=\left(X_{1}, \ldots, X_{d}\right)^{T}$ with probability measure $\mu$ on $E$. We further assume that the $X_{i}$ 's are independent and that $f(\boldsymbol{X})$ belongs to $L^{2}(E, \mu)$. Denoting by $\mu_{i}$ the marginal distribution of $X_{i}$ on $E_{i}(i=1, \ldots, d)$, we then have $\mu=\mu_{1} \otimes \cdots \otimes \mu_{d}$. In this framework, $f(\boldsymbol{X})$ can be decomposed uniquely as a sum of terms of increasing complexity

$$
\begin{equation*}
f(\boldsymbol{X})=f_{0}+\sum_{1 \leq i \leq d} f_{i}\left(X_{i}\right)+\sum_{1 \leq i<j \leq d} f_{i, j}\left(X_{i}, X_{j}\right)+\cdots+f_{1, \ldots, d}\left(X_{1}, \ldots, X_{d}\right) \tag{6}
\end{equation*}
$$

under centering conditions $\mathbb{E}\left[f_{I}\left(X_{I}\right)\right]=0$ and non-overlapping conditions $\mathbb{E}\left[f_{I}\left(X_{I}\right) \mid X_{J}\right]=0$, for all sets $I \subseteq\{1, \ldots, d\}$ and all strict subsets $J$ of $I$. We have used the set notation $X_{I}$ to represent the subvector of $\boldsymbol{X}$ obtained by selecting the coordinates belonging to $I$. These conditions imply that all the terms of (6) are orthogonal, leading to the variance decomposition

$$
\begin{equation*}
\operatorname{Var} f(\boldsymbol{X})=\sum_{1 \leq i \leq d} \operatorname{Var} f_{i}\left(X_{i}\right)+\sum_{1 \leq i<j \leq d} \operatorname{Var} f_{i, j}\left(X_{i}, X_{j}\right)+\cdots+\operatorname{Var} f_{1, \ldots, d}\left(X_{1}, \ldots, X_{d}\right) \tag{7}
\end{equation*}
$$

Due to this property, the functional decomposition (6) is often called ANOVA (ANalysis Of VAriance) decomposition. Originating from Hoeffding (1948), it was revisited by Efron and Stein (1981), Antoniadis (1984) and Sobol' (1993). For a given set $I \subseteq\{1, \ldots, d\}$, we call the corresponding term of (7) partial variance (denoted $D_{I}$ ), and call its normalized version Sobol' index (denoted $S_{I}$ ):

$$
D_{I}=\operatorname{Var}\left(f_{I}\left(X_{I}\right)\right), \quad S_{I}=\frac{D_{I}}{D}
$$

where $D=\operatorname{Var} f(\boldsymbol{X})$ is the overall variance (total variance). In particular, for $i \in\{1, \ldots, d\}$, the first-order Sobol' index $S_{i}$ corresponds to the proportion of variance of $f(\boldsymbol{X})$ explained by $X_{i}$ only. In order to include also the interactions of $X_{i}$ with the other variables, the total partial variance and the total Sobol' index are defined by

$$
D_{i}^{\mathrm{tot}}=\sum_{I \supseteq\{i\}} \operatorname{Var}\left(f_{I}\left(X_{I}\right)\right), \quad S_{i}^{\mathrm{tot}}=\frac{D_{i}^{\mathrm{tot}}}{D}
$$

Note that practitioners also call the (total) partial variances unnormalized (total) Sobol' indices. In the sequel, we will use these two words interchangeably.
The total Sobol' index can be used for screening. Indeed, under mild conditions, if $S_{i}^{\text {tot }}=0$ then the function $f$ does not depend on $x_{i}$ over $E$ (in the pointwise sense).

When the derivatives are available, a global sensitivity index can be obtained by integration. The so-called derivative-based sensitivity measure (DGSM) index of $f$ with respect to $X_{i}$ is defined by (see e.g. Sobol and Gresham (1995)):

$$
\begin{equation*}
\nu_{i}=\mathbb{E}\left[\left(\frac{\partial f}{\partial x_{i}}(\boldsymbol{X})\right)^{2}\right]=\int_{\mathbb{R}^{d}}\left(\frac{\partial f}{\partial x_{i}}(\boldsymbol{x})\right)^{2} d \mu(\boldsymbol{x})=\left\|\frac{\partial f}{\partial x_{i}}\right\|^{2} . \tag{8}
\end{equation*}
$$

Contrarily to variance-based indices, DGSM are not associated to a variance decomposition. Nevertheless, they can be used for screening. Indeed, under mild conditions, $\nu_{i}=0$ implies that $f$ does not depend on $x_{i}$ over $E$.

### 2.3 Chaos expansion serving sensitivity analysis

One main advantage of using an orthonormal basis for sensitivity analysis is that, once the expansion has been obtained, the variance-based indices can be computed in a straightforward way as a sum of squared coefficients (Sudret, 2006, 2008). More precisely, let $f$ be in $L^{2}(\mu)$, and let $\left(\Phi_{\alpha}\right)_{\alpha \in \mathbb{N}^{d}}$ be a multivariate orthonormal basis obtained by tensorization as described in Section 2.1. The expansion of $f$ in this basis is given by

$$
\begin{equation*}
f=\sum_{\alpha \in \mathbb{N}^{d}} c_{\alpha} \Phi_{\alpha} \tag{9}
\end{equation*}
$$

By using the orthonormality we obtain the expression of the total variance

$$
\begin{equation*}
D=\sum_{\alpha \neq \mathbf{0}} c_{\boldsymbol{\alpha}}^{2} \tag{10}
\end{equation*}
$$

The expression of the total Sobol' index $S_{i}^{\text {tot }}$ is obtained by only considering the terms of the decomposition (9) that contain the variable $x_{i}$, i.e. such that $\alpha_{i} \geq 1$. Hence, we have $S_{i}^{\text {tot }}=\frac{D_{i}^{\text {tot }}}{D}$ with $D_{i}^{\text {tot }}$ the total partial variance

$$
\begin{equation*}
D_{i}^{\text {tot }}=\sum_{\alpha, \alpha_{i} \geq 1} c_{\alpha}^{2} . \tag{11}
\end{equation*}
$$

The first-order Sobol' index $S_{i}^{1}$ relies on the terms that include $x_{i}$ only, i.e., $S_{i}^{1}=\frac{D_{i}^{1}}{D}$ with

$$
\begin{equation*}
D_{i}^{1}=\sum_{\substack{\alpha, \alpha_{i} \geq 1, \alpha_{j}=0 \text { for } j \neq i}} c_{\alpha}^{2} . \tag{12}
\end{equation*}
$$

Let us now consider the case where the gradient of $f$ is available. The Poincaré basis is particularly suited to this situation. Indeed, we can derive in a straightforward way expressions of both variance-based and derivative-based indices, involving the derivatives of $f$. Due to orthonormality, the coefficients of the basis expansion in (9) are given by the projection of $f$ onto the associated basis element:

$$
\begin{equation*}
c_{\alpha}=\left\langle f, \Phi_{\alpha}\right\rangle . \tag{13}
\end{equation*}
$$

From now on, let $\left(\Phi_{\alpha}\right)_{\alpha}$ denote the Poincaré basis. Combining (5) and (13), and assuming that $\alpha_{1} \geq 1, c_{\alpha}$ can be written using the partial derivatives w.r.t variable $X_{1}$ (Roustant et al., 2020):

$$
\begin{equation*}
c_{\alpha}=\left\langle f, \Phi_{\alpha}\right\rangle=\frac{1}{\lambda_{1, \alpha_{1}}}\left\langle\frac{\partial f}{\partial x_{1}}, \frac{\partial \Phi_{\alpha}}{\partial x_{1}}\right\rangle=\frac{1}{\lambda_{1, \alpha_{1}}}\left\langle\frac{\partial f}{\partial x_{1}}, \frac{\partial \varphi_{1, \alpha_{1}}}{\partial x_{1}} \otimes \varphi_{2, \alpha_{2}} \otimes \cdots \otimes \varphi_{d, \alpha_{d}}\right\rangle \tag{14}
\end{equation*}
$$

and equivalently using partial derivatives w.r.t variable $X_{i}$ if $\alpha_{i} \geq 1$. Thus, (10), (11) and (12) can also be computed using the various partial derivatives of $f$. Whereas the theoretical expressions are equal, their estimators have different properties. For example, if the integral is evaluated by Monte Carlo simulation, the expression whose integrand has smaller variance will be more accurate. We describe in Section 3 the computation of the expansion coefficients by regression, and we empirically compare the two estimation procedures in Section 4.

Furthermore, DGSM can be computed directly from the Poincaré expansion. More precisely, we have the following proposition.

Proposition 4 (DGSM formula for Poincaré chaos). Let $f \in H^{1}(\mu)$. Let $f=\sum_{\alpha} c_{\alpha} \Phi_{\alpha}$ be the expansion of $f$ in the Poincaré chaos basis, with $c_{\alpha}=\left\langle f, \Phi_{\alpha}\right\rangle$. Then the DGSM index of $f$ with respect to $X_{i}$ is equal to:

$$
\begin{equation*}
\nu_{i}=\sum_{\alpha, \alpha_{i} \geq 1} \lambda_{i, \alpha_{i}}\left(c_{\alpha}\right)^{2} . \tag{15}
\end{equation*}
$$

Proof. Write $f=\sum_{\alpha} c_{\alpha} \Phi_{\alpha}$. Then by Proposition 3, we get

$$
\frac{\partial f}{\partial x_{i}}=\sum_{\alpha, \alpha_{i} \geq 1} c_{\alpha} \frac{\partial \Phi_{\alpha}}{\partial x_{i}}
$$

where we can constrain the sum to multi-indices $\boldsymbol{\alpha}$ such that $\alpha_{i} \geq 1$, since $\varphi_{i, \alpha_{i}}=1$ for $\alpha_{i}=0$. Now, using again the orthogonality of Poincaré basis derivatives (Proposition 3), it follows that

$$
\nu_{i}=\left\|\frac{\partial f}{\partial x_{i}}\right\|^{2}=\sum_{\alpha, \alpha_{i} \geq 1}\left(c_{\alpha}\right)^{2}\left\|\frac{\partial \Phi_{\alpha}}{\partial x_{i}}\right\|^{2}=\sum_{\alpha, \alpha_{i} \geq 1} \lambda_{i, \alpha_{i}}\left(c_{\alpha}\right)^{2} .
$$

Formula (15) extends a previous result given by Sudret and Mai (2015) when all the $\mu_{i}$ are standard Gaussian. Indeed, in that case, Poincaré chaos coincides with polynomial chaos, and $\lambda_{i, \alpha_{i}}=\alpha_{i}$.

Using the expressions provided in (11) and (15) and an inequality derived by Lamboni et al. (2013), we obtain lower and upper bounds to total partial variances as follows:

$$
\begin{equation*}
\sum_{\boldsymbol{\alpha} \in \mathcal{A}, \alpha_{i} \geq 1}\left(c_{\boldsymbol{\alpha}}\right)^{2} \leq D_{i}^{\mathrm{tot}} \leq C_{P}\left(\mu_{i}\right) \nu_{i}=\sum_{\boldsymbol{\alpha} \in \mathbb{N}^{d}, \alpha_{i} \geq 1} \frac{\lambda_{i, \alpha_{i}}}{\lambda_{i, 1}}\left(c_{\boldsymbol{\alpha}}\right)^{2} \tag{16}
\end{equation*}
$$

## 3 Computation of sparse Poincaré expansions

Let $f \in H^{1}(\mu, E)$ be a computational model defined on the input space $E \subset \mathbb{R}^{d}$, with independent input random variables and with the input probability measure $\mu$ admitting a probability density function $\rho$ fulfilling Assumption 1 for each of the marginals. In the remainder of this paper, we assume that $\rho$ is known. We also assume that we are provided with an i.i.d. sample from the input distribution and with the corresponding model evaluations and model gradient values at each of the points.

With Poincaré expansion (PoinCE) we denote the expansion of the computational model onto the Poincaré basis

$$
\begin{equation*}
f(\boldsymbol{x})=\sum_{\boldsymbol{\alpha}} c_{\boldsymbol{\alpha}} \Phi_{\boldsymbol{\alpha}}(\boldsymbol{x}) \tag{17}
\end{equation*}
$$

and with Poincaré derivative expansion in direction $i$ (PoinCE-der- $i$ ) the expression

$$
\begin{equation*}
\frac{\partial f}{\partial x_{i}}(\boldsymbol{x})=\sum_{\boldsymbol{\alpha}, \alpha_{i} \geq 1} c_{\boldsymbol{\alpha}}^{\partial, i} \frac{\partial \Phi_{\boldsymbol{\alpha}}}{\partial x_{i}}(\boldsymbol{x}) \tag{18}
\end{equation*}
$$

or the equivalent expansion using normalized basis derivatives that have unit norm in $L^{2}(\mu)$. Note that (18) is the partial derivative of (17) w.r.t. variable $X_{i}$. Because the zeroth order basis function of a Poincaré basis is the constant function, basis terms for which $\alpha_{i}=0$ have zero partial derivative w.r.t. $X_{i}$ and are not included in (18). While in theory by Equation (14), the two expressions (17) and (18) provide identical coefficients for corresponding basis elements, i.e., $c_{\boldsymbol{\alpha}}=c_{\boldsymbol{\alpha}}^{\partial, i}$ for $\boldsymbol{\alpha} \in\left\{\boldsymbol{\alpha}^{\prime} \in \mathcal{A}: \alpha_{i}^{\prime} \geq 1\right\}$, in practice they will not coincide when estimated from a data set of finite size. This will be investigated in Section 4 for a number of numerical examples. In this section, we describe how such expansions are computed in practice: this concerns the computation of the Poincaré basis functions, the choice of truncation, the location of the sampled points, and the method for computing the coefficients. The implementation relies on and integrates into the UQLab framework (Marelli and Sudret, 2014).

### 3.1 Implementation of Poincaré basis functions

As described in Section 2.1.3, Poincaré basis functions are tensor products of univariate Poincaré basis functions. Each 1D basis consists of the eigenfunctions of the Poincaré differential operator associated with the respective marginal distribution (Theorem 1).

A Poincaré basis is guaranteed to exist for marginal distributions fulfilling Assumption 1 and for the Gaussian distribution. Other distributions have to be transformed or truncated to allow for a Poincaré basis. Since an isoprobabilistic transformation to standard variables can be highly nonlinear (Torre et al., 2019; Oladyshkin and Nowak, 2012), we opt for truncation: if the distribution is not Gaussian and has (one- or two-sided) unbounded support, we truncate it to its $10^{-6}$ - and $\left(1-10^{-6}\right)$-quantiles, respectively. ${ }^{1}$

We consider here only standard parametric families of probability densities (bounded and unbounded), although a Poincaré basis can be computed for any input distribution which after truncation fulfills Assumption 1. In particular, without any changes to the methodology PoinCE could be used in a data-driven framework (Torre et al., 2019) by computing the Poincaré basis for a dimensionwise kernel density estimate of the input distribution (assuming independence) given the available data.

As can be seen from applying the change-of-variables formula for a linear transformation to (3), the eigenvalues of the Poincare differential operator scale with the inverse of the squared support interval length. To avoid numerical difficulties, we therefore linearly transform (i.e., shift and rescale) parametric families to standard parameters using

- their bounds in the case of uniform, beta, triangular;
- their location and scale parameter in the case of Gaussian, Gumbel, Gumbel-min, Laplace, logistic;
- their (inverse) scale parameter in the case of exponential, gamma, Weibull, lognormal.

In the current implementation, distributions not belonging to this group of families are not being rescaled.

For standard uniform $(\mathcal{U}([-0.5,0.5])$ and standard Gaussian $(\mathcal{N}(0,1))$ marginals, the Poincaré basis can be analytically computed and is given by the Fourier (cosine) basis and the Hermite polynomial basis, respectively (Roustant et al., 2020). Therefore, in the special case of uniform or Gaussian marginals, we always (after rescaling) use the analytical solution.

For all other marginals, the Poincaré basis is computed numerically using linear finite elements. We use a fine uniform grid within the bounds and piecewise linear functions with local support, commonly called 'hat' functions. Using the weak formulation of the eigenvalue problem of the Poincaré differential operator given in (3), we arrive at the shifted generalized eigenvalue problem

$$
\begin{equation*}
\boldsymbol{K} \boldsymbol{a}^{(n)}=\left(\lambda_{n}+1\right) \boldsymbol{M} \boldsymbol{a}^{(n)} \tag{19}
\end{equation*}
$$

as described in Roustant et al. (2017, section 4.3), where the eigenvector $\boldsymbol{a}^{(n)}$ denotes the vector of coefficients used to express eigenfunction $\varphi_{n}$ in terms of 'hat' functions. Here $\boldsymbol{M}$ is the mass

[^0]matrix, and $\boldsymbol{K}$ is the sum of mass- and stiffness matrix. After solving this problem using Matlab's builtin function eigs, we interpolate the discrete eigenvectors with piecewise cubic splines, prescribing zero derivatives at the interval boundaries. Then, the basis derivatives are computed using centered finite differences. While more sophisticated techniques (in particular, Hermitian $C^{1}$ elements) could of course be used to improve this numerical computation procedure, it is accurate enough for our purposes of demonstrating the usefulness of PoinCE. Eigenfunctions and eigenfunction derivatives are scaled to have unit norm with respect to the measure $\mu_{i}$.

### 3.2 Choice of the basis truncation

In practice, the series in (17) and (18) cannot include an infinite number of terms, but must be truncated to a finite expansion. We denote by $\mathcal{A} \subset \mathbb{N}^{d}$ the subset of multi-indices that are included in the expansion. For PCE, $\mathcal{A}$ is typically chosen to include terms up to a certain degree $p$, resulting in the so-called total degree basis

$$
\begin{equation*}
\mathcal{A}^{p}=\left\{\boldsymbol{\alpha} \in \mathbb{N}^{d}: \sum_{i=1}^{d}\left|\alpha_{i}\right| \leq p\right\} . \tag{20}
\end{equation*}
$$

To further restrict the number of terms used in the expansion, another common truncation method is hyperbolic truncation (Blatman and Sudret, 2011)

$$
\begin{equation*}
\mathcal{A}^{p, q}=\left\{\boldsymbol{\alpha} \in \mathbb{N}^{d}:\|\boldsymbol{\alpha}\|_{q} \leq p\right\} . \tag{21}
\end{equation*}
$$

with the $\ell^{q}$-(quasi-) norm $\|\boldsymbol{\alpha}\|_{q}=\left(\sum_{i=1}^{d} \alpha_{i}^{q}\right)^{\frac{1}{q}}$ for $q \in(0,1]$.
Since the Poincaré basis is in general not polynomial, the concept of polynomial degree cannot be used to characterize the basis functions. Instead, we use the natural order of the basis functions corresponding to the increasing sequence of Poincaré eigenvalues, which also corresponds to an increasing number of oscillations (Theorem 1 ; recall that the $n$th eigenfunction has $n$ zeros). Therefore, we use the PCE terminology "degree" also for PoinCE. In particular, a degree of $\alpha_{i}=0$ denotes the constant basis function $\varphi_{i, 0}\left(x_{i}\right)=1$ associated to the eigenvalue $\lambda_{i, 0}=0$.
Often, in practice it is not known which degree is needed for a given problem. While in theory the expansion is more accurate the larger the total degree is, in practice accuracy is limited by the number of available sample points, since the quality of the regression solution (see Section 3.3) depends on the ratio of sample points to basis elements. In that case, a successful strategy consists of applying degree adaptivity, i.e., choosing the best degree for the expansion by crossvalidation (Blatman and Sudret, 2011; Lüthen et al., 2021a). This procedure is computationally inexpensive, since it only requires a new surrogate model fit for each new total degree, but no additional model evaluations. We apply leave-one-out (LOO) cross-validation together with a modification factor introduced by Chapelle et al. (2002) as in Blatman and Sudret (2011).

Both hyperbolic truncation and degree adaptivity contribute to the sparsity of the resulting expansion by identifying a suitable subset of basis functions necessary for a good approximation. Sparsity is a successful concept in regression-based PCE (Lüthen et al., 2021b). Denote by $P=|\mathcal{A}|$ the number of basis elements in the truncated expansion. $\mathcal{A}$, also called candidate
basis, contains the basis elements available for approximation. We describe below how sparse regression further selects only a subset of $\mathcal{A}$ to be active, i.e., have a nonzero coefficient. The final expansion might (and indeed often will) have less than $P$ active terms.

### 3.3 Computation of the coefficients by sparse regression

For computing the coefficients of an orthogonal expansion as in (17) and (18), there exist two main approaches. One is projection: the model $f$ is projected onto the basis functions, see (2). The resulting integral may be evaluated by Monte Carlo (MC) simulation, as done by Roustant et al. (2020) for Poincaré chaos, or by (sparse) quadrature methods (Le Martre et al., 2002; Matthies and Keese, 2005; Gerstner and Griebel, 1998). However, note that in general MC converges slowly, while quadrature is heavily affected by the curse of dimensionality.

The second approach is regression, introduced for PCE by Blatman and Sudret (2008). Here, after choosing an experimental design $\mathcal{X}=\left\{\boldsymbol{x}^{(1)}, \ldots, \boldsymbol{x}^{(N)}\right\}$ of input points, (17) is discretized as

$$
\begin{equation*}
y \approx \Psi c \tag{22}
\end{equation*}
$$

where $\boldsymbol{y}=\left(f\left(\boldsymbol{x}^{(1)}\right), \ldots, f\left(\boldsymbol{x}^{(N)}\right)\right)^{T}$ is the vector of model evaluations, $\boldsymbol{\Psi} \in \mathbb{R}^{N \times P}$ is the regression matrix with entries $\Psi_{k j}=\Phi_{j}\left(\boldsymbol{x}^{(k)}\right)$ where $j$ refers to an enumeration of the multivariate basis $\left(\Phi_{\boldsymbol{\alpha}}\right)_{\boldsymbol{\alpha} \in \mathcal{A}}$, and $\boldsymbol{c}$ is the vector of expansion coefficients. The discretization of (18) is analogous, with a vector

$$
\begin{equation*}
\boldsymbol{y}_{\partial, i}=\left(\frac{\partial f}{\partial x_{i}}\left(\boldsymbol{x}^{(1)}\right), \ldots, \frac{\partial f}{\partial x_{i}}\left(\boldsymbol{x}^{(N)}\right)\right)^{T} \tag{23}
\end{equation*}
$$

containing model partial derivatives and a regression matrix $\boldsymbol{\Psi}_{\partial, i}$ with entries

$$
\begin{equation*}
\Psi_{k j}^{\partial, i}=\frac{1}{\sqrt{\lambda_{i, \alpha_{j}(i)}}} \frac{\partial \Phi_{j}}{\partial x_{i}}\left(\boldsymbol{x}^{(k)}\right) \tag{24}
\end{equation*}
$$

where $\boldsymbol{\alpha}_{j}(i)$ denotes the $i$ th component of the $j$ th basis element characterized by the multi-index $\boldsymbol{\alpha}_{j}$ (see also Proposition 3).
The regression problem can be solved by ordinary least squares as

$$
\begin{equation*}
\hat{\boldsymbol{c}}=\arg \min _{\boldsymbol{c}}\|\boldsymbol{\Psi} \boldsymbol{c}-\boldsymbol{y}\|_{2}^{2} \tag{25}
\end{equation*}
$$

provided that enough model evaluations are available - at least $N \geq P$, or better $N \geq k P$ with $k=2,3$ to avoid overfitting. Due to the rapid growth of the total-degree basis with increasing dimension and degree, this requirement on model evaluations is often too restrictive for real-world problems.

To avoid this problem, sparse regression can be used, which regularizes the problem by encouraging solutions with few nonzero coefficients (Candès and Wakin, 2008; Kougioumtzoglou et al., 2020). An example is $\ell^{1}$-minimization:

$$
\begin{equation*}
\hat{\boldsymbol{c}}=\arg \min _{\boldsymbol{c}}\|\boldsymbol{\Psi} \boldsymbol{c}-\boldsymbol{y}\|_{2}^{2}+\lambda\|\boldsymbol{c}\|_{1} . \tag{26}
\end{equation*}
$$

The $\ell^{1}$-norm penalizes the coefficient vector so that sparse solutions are preferred. The sparse regression formulation allows for accurate solutions even in the case $N<P$. There exist many sparse regression methods utilizing different formulations of the sparse regression problem, see e.g. Lüthen et al. (2021b) for an overview of available sparse regression solvers in the context of PCE. In this work, we use the sparse solver Hybrid Least Angle Regression (Hybrid-LARS) (Blatman and Sudret, 2011; Marelli et al., 2021).

A result by Candès and Plan (2011) on sparse recovery emphasizes the importance of isotropy of the row distribution of the regression matrix, i.e., the requirement that for a row $\boldsymbol{a}=$ $\left(\Phi_{\boldsymbol{\alpha}_{1}}(\boldsymbol{x}), \ldots, \Phi_{\boldsymbol{\alpha}_{P}}(\boldsymbol{x})\right)$ of the regression matrix $\boldsymbol{\Psi}$ it holds that $\mathbb{E}\left[\boldsymbol{a}^{T} \boldsymbol{a}\right]=I_{P}$, where $I_{P}$ is the identity matrix of size $P$, and the expectation is with respect to the distribution of the experimental design points. If the experimental design points are chosen to follow the input distribution, the distributions of regression matrix rows for Poincaré as well as for normalized Poincaré derivative expansions are isotropic by construction due to orthonormality of the bases w.r.t. the input distribution. To improve the space-filling property of the experimental design, we use Latin Hypercube Sampling (LHS) (McKay et al., 1979) with maximin distance optimization.

### 3.4 Coefficients and Sobol' indices for Poincaré derivative expansions

Due to the rescaling of the input distribution described in Section 3.1, and the normalization of the PoinCE-der- $i$ basis elements to unit norm, the computed regression coefficients first have to be rescaled with the appropriate factors before they can be postprocessed to partial variances: if the vector $\hat{\boldsymbol{c}}^{\partial, i}$ is the regression solution using normalized (see Proposition 1) and rescaled (see Section 3.1) basis partial derivatives, then the regression solution to (18) is the vector $\tilde{\boldsymbol{c}}^{\partial, i}$ with entries

$$
\begin{equation*}
\tilde{c}_{\boldsymbol{\alpha}}^{\partial, i}=\frac{\gamma_{i}}{\sqrt{\lambda_{i, \alpha_{i}}}} \hat{c}_{\boldsymbol{\alpha}}^{\partial, i} \quad \text { for } \boldsymbol{\alpha} \in\left\{\boldsymbol{\alpha}^{\prime} \in \mathcal{A}: \alpha_{i}^{\prime} \geq 1\right\} \tag{27}
\end{equation*}
$$

where $\gamma_{i}$ is the factor by which the $i$ th input was divided in order to get to standard parameters. Theoretically, $\tilde{c}_{\alpha}^{\partial, i}$ is equal to the PoinCE solution $c_{\alpha}$ from (17), however when estimated from a data set of finite size they will in general not coincide.

We now explain how to compute the total variance from all $d$ Poincaré derivative expansions. By construction, the $i$ th Poincaré derivative expansion (PoinCE-der- $i,(18)$ ) only provides coefficients corresponding to the basis elements from the set $\left\{\boldsymbol{\alpha} \in \mathcal{A}: \alpha_{i} \geq 1\right\}$, since the partial derivatives of the basis elements $\left\{\boldsymbol{\alpha} \in \mathcal{A}: \alpha_{i}=0\right\}$ w.r.t. $X_{i}$ are zero and therefore no coefficient value can be computed for those elements. The set $\left\{\boldsymbol{\alpha} \in \mathcal{A}: \alpha_{i} \geq 1\right\}$ is sufficient for computing partial variances for variable $i$ as in (11) and (12), but not enough for computing the total variance (10), which requires all coefficients $c_{\boldsymbol{\alpha}}, \boldsymbol{\alpha} \in \mathcal{A}$, and which is needed for normalizing the partial variances to Sobol' indices. To compute the total variance from PoinCE-der expansions, we therefore aggregate the coefficients of all $d$ PoinCE-der- $i$ expansions into one vector $\tilde{\boldsymbol{c}}^{\partial}{ }^{\text {avg }}$ as follows:

$$
\begin{equation*}
\tilde{c}_{\boldsymbol{\alpha}}^{\partial \text {,avg }}=\frac{1}{\#\left\{i: \alpha_{i} \geq 1\right\}} \sum_{i: \alpha_{i} \geq 1} \tilde{c}_{\boldsymbol{\alpha}}^{\partial, i} \quad \text { for each } \boldsymbol{\alpha} \in \mathcal{A} \backslash \mathbf{0}, \tag{28}
\end{equation*}
$$

i.e., every PoinCE-der- $i$ expansion which computed a coefficient value for the basis element with index $\boldsymbol{\alpha}$ contributes equally to the averaged value. It follows that in theory, the averaged coefficient $\tilde{c}_{\alpha}^{\partial, \text { avg }}$ is equal to the PoinCE solution $c_{\alpha}$ from (17), too. It can therefore be used to estimate the total variance according to (10).
Let $\tilde{\boldsymbol{c}}_{\boldsymbol{\alpha}}^{\partial \text { avg }}=\left(c_{\boldsymbol{\alpha}}^{\partial \text { avg }}\right)_{\boldsymbol{\alpha} \in \mathcal{A} \backslash \mathbf{0}}$ be in the form of a column vector in $\mathbb{R}^{(P-1) \times 1}$. In order to use the averaged PoinCE-der expansion also as a surrogate model, we estimate the remaining coefficient $c_{0}^{\partial \text { avg }}$ corresponding to the constant term by ordinary least-squares on the residual $\boldsymbol{y}_{\text {res }}$ :

$$
\begin{align*}
\boldsymbol{y}_{\mathrm{res}} & =\boldsymbol{y}-\boldsymbol{\Psi}\binom{0}{\tilde{\boldsymbol{c}}_{\alpha}^{\partial \text { avg }}},  \tag{29}\\
c_{0}^{\partial, \mathrm{avg}} & =\frac{1}{N} \sum_{k=1}^{N} \boldsymbol{y}_{\mathrm{res}}^{(k)} \tag{30}
\end{align*}
$$

## 4 Numerical results

We investigate the performance of PoinCE (both based on model evaluations and on derivatives) on two numerical examples. The focus of our study is on Sobol' sensitivity analysis, but we also investigate DGSM-based upper bounds to partial variances, validation error (relative meansquared error) and sparsity. Our implementation is based on UQLab (Marelli and Sudret, 2014) and integrates into its PCE module (Marelli et al., 2021).

We use the following estimation techniques to compute the Sobol' indices of the models:

- PCE-LARS: PCE computed by LARS (with generalized polynomial chaos adapted to the respective input) (Blatman and Sudret, 2011; Marelli et al., 2021)
- PoinCE-MC / PoinCE-der-MC: MC-based computation using the Poincaré basis/the Poincaré partial derivative basis as in Roustant et al. (2020)
- PoinCE-LARS / PoinCE-der-LARS: Poincaré expansion and Poincaré derivative expansion computed by LARS (see Section 3.3)

Partial variances are normalized to Sobol' indices using the total variance. For PCE-LARS and PoinCE-LARS, the total variance is computed from the expansion coefficients as in (10). For PoinCE-MC and PoinCE-der-MC, we use the sample variance as done by Roustant et al. (2020). For PoinCE-der-LARS, the total variance is obtained by the procedure detailed in Section 3.4. The DGSM-based upper bound to the total partial variances is computed from (16) using the coefficients of the PoinCE derivative expansions as described in Section 3.4. Note that the inequalities in (16) are analytical bounds that do not necessarily hold for the estimated quantities. For uniform and Gaussian input variables, the analytical expression for the Poincaré basis functions is used, while for all others, the basis functions are computed numerically using a resolution of $10^{3}$ points for the uniform grid within the given bounds (see Section 3.1).

### 4.1 Dyke cost model

Our first application is a simplified analytical model computing the cost associated to a dyke that is to be constructed along a stretch of river to prevent flooding (Iooss and Lemaître, 2015). Its output is the cost in million euros given by

$$
\begin{equation*}
Y=\mathbb{1}_{S>0}+\left[0.2+0.8\left(1-\exp ^{-\frac{1000}{S^{4}}}\right)\right] \cdot \mathbb{1}_{S \leq 0}+\frac{1}{20}\left(8 \cdot \mathbb{1}_{H_{d} \leq 8}+H_{d} \cdot \mathbb{1}_{H_{d}>8}\right) \tag{31}
\end{equation*}
$$

where $S$ is the maximal annual overflow and $H_{d}$ is the dyke height. Here, the first term represents the cost of the consequences of a flooding event, the second describes the maintenance costs, and the third is associated to the construction cost. $S$ is computed from the river characteristics detailed in Table 1 via the 1D Saint-Venant equations under several simplifying assumptions as follows:

$$
\begin{equation*}
S=\left(\frac{Q}{B K_{s} \sqrt{\frac{Z_{m}-Z_{v}}{L}}}\right)^{0.6}+Z_{v}-H_{d}-C_{b} \tag{32}
\end{equation*}
$$

The full model has 8 input variables, of which $Q, K_{s}, Z_{v}$ and $H_{d}$ are important, and $C_{b}, Z_{m}, L$ and $B$ are unimportant (see also last two columns of Table 1).

Table 1: Input variables to the dyke cost model and reference values of first-order and total Sobol' indices

| Input | Function | Unit | Distribution | $S_{i}$ | $S_{i}^{\text {tot }}$ |
| :--- | :--- | :---: | :--- | :--- | :--- |
| $Q$ | Maximal annual flowrate | $\mathrm{m}^{3} / \mathrm{s}$ | Gumbel $\mathcal{G}(1013,558)$ <br> truncated to $[500,3000]$ | 0.358 | 0.483 |
| $K_{s}$ | Strickler coefficient | - | Gaussian $\mathcal{N}\left(30,8^{2}\right)$ <br>  <br> $Z_{v}$ | River downstream level | m |
| $Z_{m}$ | River upstream level | m | Triangular $\mathcal{T}(49,51)$ | 0.156 | 0.252 |
| $H_{d}$ | Dyke height | m | Uniform $\mathcal{U}([7,9])$ | 0.167 | 0.223 |
| $C_{b}$ | Bank level | m | Triangular $\mathcal{T}(55,56)$ | 0.003 | 0.008 |
| $L$ | Length of river stretch | m | Triangular $\mathcal{T}(4990,5010)$ | 0.119 | 0.177 |
| $B$ | River width | m | Triangular $\mathcal{T}(295,305)$ | 0.029 | 0.040 |

The dyke cost model has been used by Roustant et al. (2020) to demonstrate the performance of projection-based PoinCE. We compare the new regression-based methods PoinCE-LARS and PoinCE-der-LARS with the projection-based counterparts PoinCE-MC and PoinCE-der-MC, and additionally with the standard PCE method PCE-LARS. The projection-based estimates use a basis of total degree 2, while the regression-based estimates use degree adaptivity with a degree of up to 5 (remember that for PoinCE, the degree corresponds to the ordering of the eigenfunctions by the magnitude of the eigenvalues). The experimental design is sampled by LHS with maximin distance optimization. For each size of the experimental design, we perform 50 independent repetitions.

We display results only for three input variables: the most important variable $Q$, the lowimportance variable $C_{b}$, and the unimportant variable $B$. The results for the remaining input
variables can be found in Appendix A.

### 4.1.1 Comparison of MC-based and regression-based computation of PoinCE(der)

First we investigate the two different ways to compute PoinCE: projection-based as in Roustant et al. (2020) versus sparse regression-based as described in Section 3. Figures 1 and 2 show estimates for first-order and total Sobol' indices. We observe that in all cases the regressionbased estimates have a smaller variance than the corresponding projection-based estimates. Also, the median of the regression-based estimates is closer to the true Sobol' index value than the median of the projection-based estimates. Note that while the regression-based estimates use a degree-adaptive basis of $p \leq 5$, the projection-based estimates use a fixed degree of only $p=2$. While this choice introduces a certain bias to the projection-based estimates, a larger value for $p$ leads to unfeasibly large variance for those estimates. This is because the coefficients of higher-order terms cannot be estimated precisely with few experimental design points, which makes the overall estimate less precise.

We also observe that regression-based estimates are often clustered around the true Sobol' index already for very small experimental design sizes. While projection computes every coefficient independently, regression finds coefficients which jointly approximate the data as well as possible. Therefore, if an important term is missing from the expansion, projection will underestimate the true variance, while regression can somewhat compensate for the missing term by adjusting the remaining coefficients. Since regression generally leads to more precise estimates than projection, in the remainder of this paper we only show regression-based PoinCE estimates.

Furthermore, as already observed by Roustant et al. (2020), PoinCE-der estimates for Sobol' indices have a smaller variance than PoinCE estimates. In the case of projection-based estimates, this is the case if the derivative has a smaller variance than the original model. In the case of regression, the explanation might be that PoinCE-i-der has to compute less coefficients than PoinCE for the same number of experimental design points ( $\left\{\boldsymbol{\alpha} \in \mathcal{A}: \alpha_{i}>0\right\}$ vs. $\mathcal{A}$ ), which can result in a more precise estimate of the true coefficient values.

### 4.1.2 Comparison of regression-based PoinCE-der with PCE and the DGSM-based upper bound

Next, we investigate the performance of regression-based PoinCE compared to state-of-the-art PCE, and the usefulness of the DGSM-based upper bound to partial variances derived in (16). The corresponding results, unnormalized ${ }^{2}$ estimates for first-order and total Sobol' indices, are displayed in Figs. 3 and 4. Because PoinCE-der achieves more accurate estimates than PoinCE, we compute the DGSM-based upper bound using the PoinCE-der- $i$ coefficients. For total Sobol' indices, we also include a precise Monte Carlo estimate for the DGSM-based upper bound (using $10^{7}$ derivative samples) computed from (8) and the second inequality of (16).

[^1]

Figure 1: Comparison of PoinCE estimates of first-order Sobol' indices for the dyke cost model. Degree $p=2$ for the MC-based estimates and $p \leq 5$ (degree-adaptive) for the regression-based estimates. Results for the remaining variables are displayed in Fig. 9 in the appendix.


Figure 2: Comparison of PoinCE estimates of total Sobol' indices for the dyke cost model. Degree $p=2$ for the MC-based estimates and $p \leq 5$ (degree-adaptive) for the regression-based estimates. Results for the remaining variables are displayed in Fig. 10 in the appendix.

We make the following observations: the PCE-LARS estimates are generally very similar to the PoinCE-LARS estimates, but the latter often have a slightly larger range. The similarity might be because both rely on model evaluations only. However, the respective basis functions have a very different shape (for inputs that do not follow a Gaussian distribution). In particular, the PoinCE basis functions by construction obey Neumann boundary conditions, i.e., have zero derivative on the boundary.

As observed before for normalized indices, PoinCE-der performs better than PoinCE: the median is closer to the true value, and the range is smaller. This effect is especially pronounced for lowimportance variables. In 8 dimensions, a PoinCE-der expansion of degree 5 has 495 terms while the total-degree basis of PCE and PDO has 1287 terms. This means that a PoinCE-der expansion has to estimate less than half of the coefficients. PoinCE-der generally gives a tighter "lower bound" than PCE (but note that the estimates are not guaranteed to be a lower bound). By construction (16), the DGSM-based upper bound estimate is larger than or equal to the corresponding total Sobol' index estimate. However, it would be an upper bound to the true Sobol' index value only if the full infinite expansion was used. This is visible in Figs. 4 and 12: for some inputs, the upper bound estimate almost coincides with the Sobol' index estimate, and is smaller than the true Sobol' index.

For some inputs, such as $K_{s}$ and especially $H_{d}$ (see Fig. 12), the DGSM-based upper bound is not tight. An explanation for the large gap between the Sobol' index and the upper bound for $H_{d}$, using (16), might be that $H_{d}$ is responsible for a kink in the model (through the last term of (31)) which is difficult to approximate by differentiable functions. Therefore, its expansion needs high-order terms. From (16) it follows that even if the PoinCE-der found the true coefficients, the estimated upper bound is not tight due to the quickly-growing eigenvalues.

Finally, we see that in many cases the gap between the estimated upper bound estimate and its true value is larger than the gap between Sobol' index estimate and its true value. The reason might be that some higher-order terms are still missing from the considered expansion. Due to the eigenvalue factor involved in the estimate of the upper bound (16), this has a larger influence on the upper bound than on the Sobol' index estimate.

### 4.1.3 Comparison of total variance, relative mean-squared error, and sparsity

To estimate Sobol' indices precisely, it is crucial to have a good estimate for the total variance. For PCE-LARS, PoinCE-LARS, and PoinCE-der-LARS, this value can directly be computed from the expansion coefficients, while for PoinCE computed by projection, we are using the empirical variance, as detailed in the beginning of Section 4. In Fig. 5a we display the scatter of the variance estimates ( 50 replications). The empirical estimate has the largest variation, while PoinCE-der-LARS has the smallest. PCE-LARS and PoinCE-LARS underestimate the total variance more than PoinCE-der-LARS. This is likely one reason for the good performance of PoinCE-der for the estimation of Sobol' indices: a more accurate total variance leads to more accurate Sobol' indices.


Figure 3: Estimates of unnormalized first-order Sobol' indices for the dyke cost model ( $p \leq 5$ ). Boxplots: in grey the PCE-based estimates and in black the DGSM-based upper bound from (16). The dashed line ("True value") denotes a high-precision estimate for the unnormalized first-order Sobol' index. Results for the remaining input variables can be found in Fig. 11 in Appendix A.


Figure 4: Estimates of unnormalized total Sobol' indices for the dyke cost model ( $p \leq 5$ ). Boxplots: in grey the PCE-based estimates and in black the DGSM-based upper bound from (16). Lines: the dashed line ("True value") denotes a high-precision estimate for the unnormalized total Sobol' index, while the dotted line ("UB true value") is a MC-based high-precision estimate for the DGSM-based upper bound. Results for the remaining input variables can be found in Fig. 12 in Appendix A.

Interestingly, while PoinCE performs well for the estimation of Sobol' indices, this is not true for the generalization error, given by the relative mean-squared error

$$
\begin{equation*}
\operatorname{RelMSE}=\frac{\mathbb{E}_{X}\left[\left(f(X)-f^{\text {surr }}(X)\right)^{2}\right]}{\operatorname{Var}_{X}[f(X)]} \tag{33}
\end{equation*}
$$

with the surrogate model $f^{\text {surr }}$. The RelMSE is computed by Monte Carlo integration on a validation set of size $10^{6}$ sampled from the input distribution $\mu$. In Fig. 5b we display boxplots of estimates for the generalization error on a validation set of size $10^{6}$ (mean-squared error normalized by the variance of the validation set). PoinCE-der attains a smaller relative MSE than PoinCE. PCE shows faster convergence behavior than both, and attains a smaller relative MSE than PoinCE. PoinCE-der performs better than PCE for the two small experimental design sizes, which shows that the information brought by derivatives might be especially useful when data is scarce.

Finally, we display the number of nonzero coefficients of each expansion in Fig. 5c. PCE and PoinCE have a similar number of active coefficients, while PoinCE-der has considerably more active coefficients. This corresponds to a better validation error only for the smallest experimental design size (Fig. 5b). A possible explanation is the following: using the derivative information to compute PoinCE-der, the same amount of data is used for a (truncated) expansion containing less terms than the PoinCE expansion has (since some terms, which do not contain the variable for which the partial derivative is taken, drop out). In this way, more coefficients can be computed.


Figure 5: Dyke cost model: Comparison of PCE and PoinCE(-der) with respect to the following metrics: estimation of total variance, relative mean-squared error, and number of nonzero coefficients in the respective expansions.

### 4.2 Mascaret data set

Our second application focuses on a phenomenological and industrial simulation model, called Mascaret (Goutal et al., 2012), based on a 1D solver of the Saint Venant equations and aiming at computing water height for river flood events. The studied case, taken from Petit et al. (2016)
and also studied in Roustant et al. (2017), is the French Vienne river in permanent regime whose uncertain input data concern flowrate, several physical parameters and geometrical data (transverse river profiles). 37 independent inputs have then been considered as random variables:

- 12 Strickler coefficients of the main channel $K_{s, c}^{i}$, uniform in [20,40];
- 12 Strickler coefficients of the flood plain $K_{s, p}^{i}$, uniform in [10, 30];
- 12 slope perturbations $d Z^{i}$, standard Gaussian with bounds [ $-3,3$ ];
- 1 discharge value $Q$, Gaussian with zero mean and standard deviation 50 , bounds [ $-150,50$.

The derivatives of the model output with respect to these 37 inputs have been efficiently (with a cost independent of the number of inputs) computed by using the adjoint model of Mascaret obtained by automatic differentiation (Griewank and Walther, 2008). A large-size Monte Carlo sample ( $n=20000$ ) is available from the study of Petit et al. (2016). This data set contains all the values of the 37 inputs, the water height as output and the 37 partial derivatives of the output (one derivative with respect to each input). Note that this sample, which has a very large size, has been obtained during a research work for a demonstrative purpose. In industrial practice, the aim is to use the minimal possible sample size: it is expected to use methods able to deal with sample sizes of the order of one hundred.

Previous studies on this data set (Petit et al., 2016; Roustant et al., 2017) have identified 32 of the 37 inputs as noninfluential. In our study, we display results for the 5 remaining inputs ( $K_{s, c}^{11}$, $K_{s, c}^{12}, d Z^{11}, d Z^{12}$, and $Q$ ) and for one of the noninfluential inputs $\left(K_{s, c}^{1}\right)$. We choose a basis with hyperbolic truncation using $q=0.5$, and degree adaptivity $p=1,2, \ldots, 8$. We analyze several experimental design sizes ranging from 30 to 300 . For each experimental design size, we run 30 replications, sampling the design randomly without replacement from the given full data set. "True" values for Sobol' indices and total variance are computed from a PCE using all 20000 points.

### 4.2.1 Comparison of regression-based PoinCE(der) with PCE and the DGSMbased upper bound

Estimates of first-order and total Sobol' indices are displayed in Figs. 6 and 7. We display results for regression-based PCE, PoinCE, and PoinCE-der. In addition, we display the upper bound computed as in (16), computed based on PoinCE-der coefficients and normalized by the PoinCE-der total variance. We observe that for the non-influential variable $K_{s, c}^{1}$ (and indeed all other 31 non-influential variables), derivative-based PoinCE correctly identify a total and first-order Sobol' index of 0 . Overall, PoinCE and PCE show very similar results, with PoinCE having slightly larger variance in a few cases. For some variables such as $d Z^{11}$ and $Q$, the DGSM-based estimate of the upper bound almost coincides with the PoinCE-der estimate. Overall, we observe that PoinCE-der estimates have smaller variance than PCE and PoinCE for the important variables $K_{s, c}^{11}, d Z^{11}$, and $Q$, even already for 30 experimental design points.

For low-importance variables such as $K_{s, c}^{12}$ and $d Z^{12}$, PoinCE-der correctly identifies a value away from zero already for the smallest experimental design, while half of the PCE and PoinCE estimates are zero. For small experimental design sizes, the PoinCE-der estimates also have a smaller bias than the PCE and PoinCE estimates. Sometimes the PoinCE-der estimates seem to systematically over- or underestimate the true Sobol' index by a small amount. However, note that the "true" value was computed by a PCE (based on all 20000 points), and might therefore itself be slightly inaccurate.


Figure 6: First-order Sobol' indices for the Mascaret data set (30 replications). "True" values computed from a PCE using all 20000 points.

### 4.2.2 Comparison of total variance and relative mean-squared error

In Fig. 8a we display various estimates for the total variance. We observe again that PoinCE-der yields an estimate with smaller variance and less bias than PoinCE and PCE. PoinCE and PCE both have smaller variance than the empirical estimate, but generally underestimate the total variance.

While PoinCE-der estimates Sobol' indices and total variance well, we observe in Fig. 8b showing the relative MSE that PCE and PoinCE are performing better as global surrogate models: their model approximation error is for large experimental designs almost an order of magnitude better than for PoinCE-der.


Figure 7: Total Sobol' indices for the Mascaret data set (30 replications). "True" values computed from a PCE using all 20000 points.


Figure 8: Mascaret data set. Comparison of PCE and PoinCE(-der) with respect to the following metrics: estimation of total variance, and relative mean-squared error. "True" value of total variance computed from a PCE using all 20000 points.

## 5 Conclusion

In this paper we studied PoinCE, an expansion in terms of the Poincaré basis, which is an orthonormal basis of $L^{2}(\mu)$ with the unique property that all its partial derivatives form again an orthogonal basis for the same space. We provided a proof of this property as well as a few analytical results as direct consequences. In particular, we showed how upper and lower bounds for partial variances can be obtained analytically from PoinCE coefficients.

We described the computation of PoinCE and Poincaré derivative expansions by sparse regression and applied the method to two numerical examples. We found that while PoinCE does not outperform PCE in terms of validation error, it can be advantageous for estimating Sobol' indices in the low-data regime. PoinCE is therefore a valuable tool if model derivatives are cheaply available (e.g., by automatic differentiation or as a by-product of the simulation). Taking partial derivatives reduces the size of the truncated basis especially for high-dimensional, low-order total-degree bases, which gives an advantage to derivative-based PoinCE over expansions relying on model evaluations.

Future work on the topic of PoinCE will investigate the simultaneous use of model evaluations and derivatives for the computation of the coefficients, and compare to the related topic of gradient-enhanced PCE.

## Acknowledgments

This paper is a part of the project "Surrogate Modeling for Stochastic Simulators (SAMOS)" funded by the Swiss National Science Foundation (Grant \#200021_175524), whose support is gratefully acknowledged. Part of this research was conducted within the frame of the Chair in Applied Mathematics OQUAIDO, gathering partners in technological research (BRGM, CEA, IFPEN, IRSN, Safran, Storengy) and academia (CNRS, Ecole Centrale de Lyon, Mines SaintEtienne, University of Grenoble, University of Nice, University of Toulouse) around advanced methods for computer experiments. Support from the ANR-3IA Artificial and Natural Intelligence Toulouse Institute is gratefully acknowledged.

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## A Additional results

In Figs. 9 to 12, we show additional results for the dyke cost model, namely Sobol' index estimates (normalized and unnormalized) for the remaining five input dimensions. For the corresponding discussion, see Section 4.1.


Figure 9: Comparison of PoinCE estimates of first-order Sobol' indices for the dyke cost model. Degree $p=2$ for the MC-based estimates and $p \leq 5$ (degree-adaptive) for the regression-based estimates. See also Fig. 1 in the main part of the paper.


Figure 10: Comparison of PoinCE estimates of total Sobol' indices for the dyke cost model. Degree $p=2$ for the MC-based estimates and $p \leq 5$ (degree-adaptive) for the regression-based estimates. See also Fig. 2 in the main part of the paper.


Figure 11: Estimates of unnormalized first-order Sobol' indices for the dyke cost model $(p \leq 5)$.
Boxplots: in grey the PCE-based estimates and in black the DGSM-based upper bound from (16). The dashed line ("True value") denotes a high-precision estimate for the unnormalized first-order Sobol' index. See also Fig. 3.


Figure 12: Estimates of unnormalized total Sobol' indices for the dyke cost model ( $p \leq 5$ ). Boxplots: in grey the PCE-based estimates and in black the DGSM-based upper bound from (16). Lines: the dashed line ("True value") denotes a high-precision estimate for the unnormalized total Sobol' index, while the dotted line ("UB true value") is a MC-based high-precision estimate for the DGSM-based upper bound. See also Fig. 4.


[^0]:    ${ }^{1}$ One might argue that this can distort the results obtained with PoinCE, especially in the tails. It is true that this truncation introduces a small error. However, as all such methods, PoinCE by design approximates accurately mainly the bulk, not the tails (for this, specialized techniques like subset simulation shall be used). Furthermore, in practical applications it is a modelling choice how to represent the input distribution. Choosing an unbounded parametric distribution is common, but not necessarily the most sensible choice, since for virtually every quantity in the real world there is an upper bound that cannot be exceeded.

[^1]:    ${ }^{2}$ We show unnormalized indices because the DGSM-based upper bound is not normalized.

