Sample Solutions to HW2

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Q1 Meniscus Adjacent to a Wilhelmy Plate

1. Interfacial pressure when $\theta > 90^{\circ}$

At the plate the meniscus, we know $p(x = 0) - p_{air} + 2\mathcal{H}\gamma = 0$. Since $\theta > \pi/2$, the mean principle radius $\mathcal{H} < 0$, therefore we have p(x = 0) > 1 atm.

2. Function f

In this question we need to solve a second-order ode equation. First we shall express the ode equation using vector representation, as the skeleton code for function dHX as follows (written in the Matlab style):

The function f you need is then to take initial values and put dHX into the ode solver:

3. Influence of initial value on the precision of result.

Ideally, at $X \to \infty$, we have both H and $\frac{dH}{dX}$ are 0. Therefore, the values we obtained from the function f, represent the error of both H and $\frac{dH}{dX}$ at $X = X_{\text{max}}$. Here we plot the absolute errors as functions of the initial guess of H_0 , with different choices of X_{max} for the ode that we study, as shown in Figure 1. As can be seen, the absolute error varies dramatically within a certain range (shown in blue shader). The width of such regime also becomes narrower when X_{max} increases. Therefore, the solution of the ODE is greatly influenced by the initial guess of H_0 .

4. Dependency of H_0 and θ

Since both errors of *H* and $\frac{dH}{dX}$ exist, to solve the ode equation, we can use the shooting method to get the correct H_0 value at a certain θ value:

The solution can be found by non-linear solvers: (e.g. fsolve in Matlab or scipy.optimize.fsolve in Python). One thing you need to take care is how to estimate the residue *R*. To achieve a more stable solution you may want to evaluate both $H(\infty)$ and $dH/dX(\infty)$ for *R*, such as: $R = [H^2(\infty) + (dH/dX)^2(\infty)]$. Using the shooting method, we get the dependency of H_0 as function of θ as shown in Figure 2. As can be expected, the height of the meniscus is larger than 0 when $\theta > 0^\circ$, and *vice versa*.

5. Height profiles

The height profile as function of *X* can be solved using the ode solver like ode45 on an array of *X* values. The result can be seen in Figure 3.



Figure 1: Absolute error of H and $\frac{dH}{dX}$ at $X = X_{max}$, when (a) $X_{max} = 3$, (b) $X_{max} = 5$, (c) $X_{max} = 7$ and (a) $X_{max} = 9$. The blue shader corresponding to the regime when the error dramatically changes. The answers here are obtained using scipy.integrate.ode interface in Python with (4)5-order Runge-Kutta algorithm.

Algorithm 1 Pseudo-code for the shooting methodprocedure SHOOTING(\hat{H}_0, θ_0) $H(0) \leftarrow \hat{H}_0; \quad \frac{dH}{dX} \leftarrow -\cot \theta_0$ $R \leftarrow ODE(H(0), \frac{dH}{dX})|_{X \to \infty}$ while R >tolerance do $H(0) \leftarrow H(0) + \delta H; \quad \frac{dH}{dX} \leftarrow \frac{dH}{dX} + \delta U$ $R \leftarrow ODE(H(0), \frac{dH}{dX})|_{X \to \infty}$ end whilereturn H(0), Rend procedure



Figure 2: (a) Hight H_0 as a function of θ at the Wilhelmy plate. (b) Residue of the estimation of $H(\infty)$ as a function of θ .



Figure 3: *X*-dependent height profile of varied contact angles.

As we expected, the meniscus near the plate is higher than 0 for $\theta_0 < 90^\circ$ (hydrophilic). On the other hand the meniscus drops below the liquid surface level for $\theta_0 > 90^\circ$.

Your plots may look different due to difference choice of solver and the conditions you used for the cutoff of "infinite" boundary. Nonetheless in your plot the range of the x axis should not exceed your cutoff value. Otherwise you will see a rise of the meniscus due to error accumulation in the ode solving process.

To get an accurate solution, some caution should be taken for the initial guess of H_0 . One trick is as follows;

- (a) Get the solution of H_0 when $\theta = 90^\circ$ (is 0!)
- (b) Use the solution of H_0 you get from step 1 as the initial guess for angles $\theta \pm \delta \theta$. $\delta \theta$ is a small angle (for instance 0.5 °)
- (c) Repeat steps 1 and 2 until you get the solution for the angle you want to study.

By making $\delta\theta$ smaller, you have higher chance to obtain a stable solution.

6. Comparing the analytical and numerical solutions

The numerical and analytical solution for cases $\theta_0 = 20^\circ$ and $\theta_0 = 75^\circ$ can be seen in Figure 4. With the offset of X_0 we see that our numerical solution is in perfect agreement with the analytical one.



Figure 4: Comparison between the numerical and analytical solutions when $\theta_0 = 20^\circ$ and $\theta_0 = 75^\circ$

Q2 Growth of Ice Crystal in Salt Water

1. Interface salt concentration

Since salt does not dissolve in solid, salt will be ejected upon ice growth, so $C_{IL} > C_{\infty}$ would occur.

2. Interface heat balance

The heat balance at interface follows

$$\rho\lambda\frac{da}{dt} = -k_L\frac{dT_L}{dx}$$

Using the temperature profiles, it follows

$$\rho\lambda\frac{da}{dt} = -k_L\frac{T_\infty - T_F(t)}{\delta}$$

Note that now it becomes an ODE with two unknowns a(t) and $T_F(t)$.

3. Interface concentration balance

The concentration balance at interface follows

$$C_{IL}(t)\frac{da}{dt} = -D_L\frac{dC_L}{dx}$$

Using the temperature profiles, it follows

$$C_{IL}(t)\frac{da}{dt} = -D_L \frac{[C_{\infty} - CIL(t)]}{\delta}$$

Note that now it becomes an ODE with two unknowns a(t) and $C_{IL}(t)$.

4. Heat and mass transfer equations

We have three unknown functions as functions of t, a(t), $T_F(t)$, and $C_{IL}(t)$, and three equations as follows.

$$\rho \lambda \frac{da}{dt} = -k_L \frac{[T_{\infty} - T_F(t)]}{\delta} \tag{1}$$

$$C_{IL}(t)\frac{da}{dt} = -D_L \frac{[C_{\infty} - C_{IL}(t)]}{\delta}$$
(2)

$$T_F(t) = T_0 - \beta C_{IL}(t) \tag{3}$$

Accordingly we can solve the ODEIVPs using matlab with the following initial values: $a(t) = 0, T_F(0) = T_W$ and $C_{IL}(0) = C_{\infty}$