UHV Techniques

As part of the course 'Characterization of Catalysts and Surfaces'

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What is Vacuum?

The term VACUUM can be used to describe these conditions:

1) Complete absence of matter (definite volume in which gases are almost absent – e.g. interstellar space)

2) Physical state in which the pressure in a definite volume is smaller than in the surroundings (e.g. smaller than the atmospheric pressure)

An example: $\rho_{gas} = 2 \cdot 10^{19} \text{ mol/cm}^3$ (atmospheric pressure) $\rho_{gas} = 10^9 \text{ mol/cm}^3$ (orbiting satellite)





Some applications of Vacuum

- Reduce the concentration of a gas below a critical level (e.g. O₂ in bulbs)
- Avoid gas-driven physico-chemical processes (e.g. experiments studying the gas-surface interaction) and increase the mean free path of particles (e.g. ion and electron spectroscopies)

• Thermal insulation

• Degasification of liquids









Vacuum classification

Vacuum	Pressure (torr)	Number Density (m ⁻³)	M.F.P. (m)	Surface Collision Freq. (m ⁻² ·s ⁻¹)	Monolayer Formation Time (s)
Atmosphere	760	2.7×10 ²⁵	7×10 ⁻⁸	3×10 ²⁷	3.3×10 ⁻⁹
Rough	10 ⁻¹ - 10⁻³	3.5×10 ¹⁹	0.05	4×10 ²¹	2.5×10 ⁻³
High	10 ⁻³ - 10⁻⁶	3.5×10 ¹⁶	50	4×10 ¹⁸	2.5
Very high	10 ⁻⁶ - 10⁻⁹	3.5×10 ¹³	50×10 ³	4×10 ¹⁵	2.5×10 ³
Ultrahigh	10 ⁻⁹ - 10⁻¹²	3.5×10 ¹⁰	50×10 ⁶	4×10 ¹²	2.5×10 ⁶ (29 days!)

Gas flow regimes

The mean free path is the average distance that a gas molecule can travel before colliding with another molecule and is determined by:

- Size of molecule (2r)
- Pressure (p)
- Temperature (T)

$$\lambda_{a} = \frac{k}{\pi\sqrt{2}} \cdot \frac{T}{(2r)^{2} p}$$

Vacuum range	Pressure in hPa (mbar)	Pressure in mmHg (Torr)	Molecules / cm ³	Molecules / m ³	Mean free path
Ambient pressure	1013	759.8	2.7 × 10 ¹⁹	2.7 × 10 ²⁵	68 nm ^[2]
Low vacuum	300 – 1	225 – 7.501×10 ⁻¹	10 ¹⁹ – 10 ¹⁶	10 ²⁵ – 10 ²²	0.1 – 100 µm
Medium vacuum	1 - 10 ⁻³	7.501×10 ⁻¹ – 7.501×10 ⁻⁴	10 ¹⁶ – 10 ¹³	10 ²² – 10 ¹⁹	0.1 – 100 mm
High vacuum	10 ⁻³ – 10 ⁻⁷	7.501×10 ⁻⁴ – 7.501×10 ⁻⁸	10 ¹³ – 10 ⁹	10 ¹⁹ – 10 ¹⁵	10 cm – 1 km
Ultra-high vacuum	10 ⁻⁷ – 10 ⁻¹²	7.501×10 ⁻⁸ – 7.501×10 ⁻¹³	10 ⁹ – 10 ⁴	10 ¹⁵ – 10 ¹⁰	1 km – 10 ⁵ km
Extremely high vacuum	<10 ⁻¹²	<7.501×10 ⁻¹³	<10 ⁴	<10 ¹⁰	>10 ⁵ km

The gas in a vacuum system can be in a **viscous state**, in a **molecular state** (or in a transition state) depending on the dimensionless parameter know as the **Knudsen number** (K_n) that is the ratio between the mean free path and the characteristic dimension of the flow channel.

$$K_n = \frac{\lambda_a}{a}$$

 $\begin{array}{l} \lambda_a \ = \ \text{mean free path} \\ a \ = \ \text{characteristic dimension of flow channel} \\ (typically \ a \ pipe \ radius) \end{array}$

Viscous Flow :





Transition Flow : 0.01 < Kn < 1.0

 $\label{eq:constraint} \mbox{Molecular Flow}: \qquad \mbox{Kn} > 1.0$





Creation of Vacuum: pumping technology

1) **Primary pumping systems**: mechanical pumps that decrease the pressure from atmospheric pressure pressures close to the ultra-high vacuum (10⁻⁶-10⁻⁸ mbar)

- Rough pumps (atmospheric pressure down to 10⁻³ mbar): membrane pumps, rotary pumps, scroll pumps, roots pumps

- Turbomolecular Pumps (from the mbar to about 10⁻⁹ mbar)

2) **UHV pumping systems:** pumps that work at low pressure and, thanks to their efficiency, allow reaching/improving the ultra-high vacuum (10⁻⁶-10⁻¹¹ mbar).

- Ion Pumps (from 10⁻⁶ mbar to 10⁻¹¹ mbar)
- Getter Pump
- Titanium Sublimation Pump

Rough pumps: membrane and rotary pumps



Membrane

Combined movement of a diaphragm (plastic-rubber) and suitable valves (check, butterfly, flap valves)



Rotary

Consists of vanes mounted to an eccentric rotor. The vanes rotate inside a cavity

Volume increased = fluid drawn into the chamber Volume decreased = fluid forced out

- Dry pumps
- Can handle gas and liquids
- The flow rate depends on the diaphragm diameter and its stroke length
- Ultimate vacuum is in the mbar (larger pumps can reach 10⁻¹ mbar)

Vanes are sealed on all edges. The rotation generates a volume expansion (gas pumping)reduction (exhaust release)

- **Oil pumps:** oil and gas are mixed inside the pump and separated externally
- Multiple stage pump can generate a good vacuum (down to 10⁻³ mbar)
- Low efficiency but possibility to pump gas, gas ullet+ dust and water

Rough pumps: scroll and roots pumps



Scroll

One of the scrolls is fixed, while the other orbits eccentrically without rotating. Compressing pockets of fluid form between the scrolls and are driven to the exhaust port

- Dry pumps
- High efficiency
- Small gas pulsation, less vibrations
- Difficult maintenance
- Ultimate vacuum is in the 10⁻² mbar range



Two 8-shaped synchronously counter-rotating rotors rotate contactlessly (small gap) in a housing

• Dry pumps

- Can generate a good vacuum (low 10⁻² mbar)
- No friction in the suction chamber, operation at high speed, large flow rate

Turbomolecular pumps



- A turbomolecular pump is used to obtain and maintain **high vacuum**.
- These pumps work on the principle that gas molecules can be given momentum in a desired direction by repeated collision with a moving solid surface.
- A rapidly spinning fan rotor (50000-100000 rpm) 'hits' gas molecules from the inlet of the pump towards the exhaust in order to create or maintain vacuum.



Ion pumps

- An ion pump is a static pump capable of reaching pressures as low as 10⁻¹¹ mbar
- Can be turned on only at pressures around (or less than) 10⁻⁴ mbar
- A strong electrical voltage (typically 3–7 kV) is applied to the anode, producing free electrons. Electrons get caught by the magnetic field and rotate around it
- Electrons hit gas molecules ionizing them
- Positively charged molecules accelerates toward the cathode (grounded) at high velocity
- The cathode is sputtered and titanium compounds deposit on the anode
- The cathode acts as a getter (e.g. adsorbs inert gases)



Rough pumps: scroll and roots pumps

Non-evaporable getter



- Static pumps helping to establish and maintain ultra-high vacuum
- Porous alloys or powder mixtures of Al, Zr, Ti, V and Fe, forming stable compounds with active gases
- Can be placed in narrow/difficult to reach spaces (particle accelerators)
- Activated by annealing to >550 K



Titanium sublimation

- Static pump helping to refine the vacuum
- Titanium filament through which a high current (typically around 40 A) is passed periodically
- Titanium sublimates and coats the surrounding chamber walls
- Components of the residual gas in the chamber which collide with the chamber wall react with titanium to form stable, solid products

Vacuum measurement



Low vacuum: Pirani (atm. pressure – 10⁻⁴ mbar)

- Two Pt filaments are the arms of a Wheatstone bridge and heated to a constant temperature
- Residual gases conduct away part of the thermal energy of the measurement filament. The amount of electrical current needed to restore its temperature is converted to a pressure readout



High vacuum: cold cathode (10⁻⁴-10⁻⁹ mbar)

2 electrodes: anode,
 cathode + permanent
 magnetic field (works like
 a ion pump!)





High-Ultrahigh vacuum: hot cathode $(10^{-4}-10^{-11} \text{ mbar})$

• 3 electrodes: filament, collector, grid

- The filament emits electrons, which are attracted to a polarized grid
- Residual gas molecules are ionized by the electrons and attracted by the collector.
 Pressure reading is determined by the electronics from the collector current.

TPD

Temperature Programmed Desorption

Temperature programmed desorption^{*}

Ultra high vacuum can be used to study the adsorption/reaction of molecules on a surface (monolayer formation time in UHV 10³-10⁶ s). Discussion based on **Langmuir ad-(de-)sorption isotherm**

Langmuir Adsorption-Desorption

- Adsorption is localized (adsorbed particles are immobile)
- Substrate surface is saturated at Θ = 1 ML (all adsorption sites occupied)
- No interactions between the adsorbed particles

*Resource for further reading: Temperature-Programmed Desorption (TPD). Thermal Desorption Spectroscopy (TDS), Sven L.M. Schroeder and Michael Gottfried, June 2002, available online.

Kinetics

$$\mathbf{r}_{des} = -\frac{d\Theta}{dt} = \mathbf{k}_{n}\Theta^{n}$$

If k (rate constant) is described by an Arrhenius eq.:

$$\mathbf{k_n} = v_n \cdot \exp\left(-\frac{\Delta E_{des}}{RT}\right)$$

The rate law is then referred to as the *Polanyi-Wigner* equation:

$$\mathbf{r}_{\mathbf{des}} = -\frac{d\Theta}{dt} = v_n \cdot \exp\left(-\frac{\Delta E_{\mathbf{des}}}{RT}\right) \cdot \Theta^n$$

- v_n: Pre-exponential factor
- n: Desorption order
- Θ: Surface coverage
- E_{des}: Activation energy for desorption

Temperature programmed desorption

A typical TPD experiment (UHV):



- Clean sample surface exposed to a precise amount of gas (usually measured in Langmuirs 1 L = 10^{-6} Torr \cdot s)
- Sample placed in front of a quadrupole mass spectrometer (QMS) and heated with a precise rate (β)
 - The quadrupole acts as a filter, separating ions with different m/z, which are then collected
 - A typical spectrum shows the intensity of a specific m/z vs. temperature

Temperature programmed desorption

If the pumping rate is faster than the desorption rate (no readsorption) a series of separated peaks can be recorded (each of them corresponding to a surface desorption process)



A TPD experiment can give important information:

- Heat of desorption
- Surface coverage (quantification of the monolayer)
- Surface reactivity (gas-substrate interaction, adsorption sites)
- Kinetics of desorption

Temperature programmed desorption

Spectral interpretation is most commonly performed using the **Polanyi-Wigner equation**

In a TPD experiment β is the heating rate, defined as $\beta = dT/dt = const$. Thus $dt = dT/\beta$ can be substituted in the equation to give



First order kinetics (molecular)







- The desorption peak areas depend on Θ
- The desorption peaks are asymmetric
- T_m constant with increasing Θ
- T_m increases with ΔE_{des}

First order kinetics: approximate evaluation of ΔE_{des}

In 1962 Redhead, assuming that activation parameters are independent of surface coverage and that desorption followed 1st order kinetics, derived a simple equation.^{*}

Solving this equation for ΔE_{des} gives:

$$\Delta E_{des} = RT_m \left[\ln \frac{\nu \cdot T_m}{\beta} - \ln \frac{\Delta E_{des}}{RT_m} \right]$$

 $\frac{\Delta E}{RT_{m}^{2}} = \frac{v}{\beta} \cdot \exp\left[-\frac{\Delta E}{RT_{m}}\right];$

The second part in the brackets is small relative to the first, and can be approximated to 3.64 (error is less than 1.5% for $10^8 < v/\beta < 10^{13} \text{ K}^{-1}$)

- T_m and β are determined experimentally
- The activation energy from a single desorption spectrum can be estimated using an approximate value for v. $v = 10^{13} \text{ s}^{-1}$ is a commonly chosen value.



As an example: in this case T_m = 117 K. Assuming v=1.0·10¹³ s⁻¹ and β=2 K/s ΔE_{des} = 29.5 kJ/mol

*P. A. Redhead, Vacuum 12, 203-211 (1962).

First order kinetics: approximate evaluation of ΔE_{des} from curves having different β

A series of spectra for the same Θ is acquired employing different $\beta = dT/dt = const$. From each spectrum, the temperature of the desorption rate maximum T_m is determined



Plotting of $ln(T_m^2/\beta)$ vs. $1/T_m$ for a series of β values provides ΔE_{des} from the slope and v from the intercept with the ordinate



Variation of the heating rate $\Theta_0 = 0.67, n = 1, 2,$ $v_n / \beta_H = 10^{11} \dots 10^{15},$ $E_d = 100 \text{ kJ/mol.}$

Second order kinetics (recombinative desoprtion)



Plotting the $\ln(T_m^2/\beta)$ vs. $1/T_m$ for a series of β values provides ΔE_{des} from the slope and (if Θ is known) v from the intercept with the ordinate

Leading edge analysis

But the activation parameters often depend on the coverage and temperature!



Habenschaden and Küppers leading edge method*

- Leading edge: O almost unchanged
- The rate of desorption is evaluated from each single leading edge

$$\mathbf{r}_{\mathbf{des}} = v_n \cdot \exp\left(-\frac{\Delta E}{RT}\right) \cdot \Theta^n \quad \longrightarrow \quad \ln(\mathbf{r}_{\mathbf{des}}) = -\frac{\Delta E}{RT} + \ln(v_n) + n \ln\Theta$$

• $ln(r_{des})$ plotted vs. 1/T. The slope gives ΔE_{des} and the intercept with y gives v

An example of TPD applied to the study of a catalytic reaction

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Laboratories and Demonstrations

A Laboratory Study of Heterogeneous Catalysis in Ultrahigh Vacuum

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Carbon monoxide oxidation on a Pt foil

The purpose of the experiment is to determine which mechanism operates in the catalytic oxidation of carbon monoxide over platinum foil. Briefly, the Langmuir-Hinshelwood (LH) mechanism [2] requires adsorption of both species involved in the reaction. The Eley-Rideal (ER) mechanism [2] involves adsorption of only one reactant with subsequent reaction by impingement of the other from the gas phase. The goal is to distinguish between these two mechanisms.



Figure 2.2: Schematic representation of the Langmuir-Hinshelwood mechanism (a) and the Eley-Rideal mechanism (b) for the catalytic oxidation of CO. Dark balls correspond to the carbon atoms and grey balls correspond to the oxygen atoms.

CO and O₂ desorption

Procedure: 50 L of gas dosed on the clean foil at RT. Sample heated from RT to 700°C with β =10 K/s



- Molecular adsorption (first order)
- Two main desorption peaks (a foil is polycristalline) at ca. 110 and 225°C

- Dissociative adsorption (recombinative desorption process)
- Multiple desorption peaks at higher temperature than for CO (larger ΔE_{des})

Temperature programmed reaction (TPR)

Procedure: 50 L of CO dosed on the clean foil at RT. Sample heated from RT to 700°C with β =10 K/s while flowing 3·10⁻⁷ Torr of O₂.



- Sample surface saturated with CO
- CO desorption peaks intensity decreases
- As the CO starts to desorb, the partial pressure of O₂ decreases and the signal of CO₂ increases

Procedure: 50 L of O₂ dosed on the clean foil at RT. Sample heated from RT to 700°C with β =10 K/s while flowing 3·10⁻⁸ Torr of CO.



- Sample surface saturated with O₂
- CO₂ is produced immediately, but its signal goes down above 300°C
- CO₂ production correlated with CO desorption



Figure 2.2: Schematic representation of the Langmuir-Hinshelwood mechanism (a) and the Eley-Rideal mechanism (b) for the catalytic oxidation of CO. Dark balls correspond to the carbon atoms and grey balls correspond to the oxygen atoms.

In case of E-R mechanism the reaction should start immediately after introducing one of the reagents (in the presence of the other adsorbed)

- In both TPR experiment CO₂ production correlates with the presence of both reagents on the sample surface
- TPR performed after CO pre-adsorption clearly demonstrates competition between the reagents for the adsorption sites (especially reaction at high temperature). CO is blocking the adsorption sites (poisoning effect), and some energy (temperature) is required to remove part of it and allow oxygen to adsorb and split
- TPR performed after O₂ pre-adsorption clearly demonstrates that adsorbed CO is necessary for CO₂ formation (no more CO₂ formed above 300°C

These model experiments support the hypothesis that a L-H mechanism operates, in good agreement with the literature



Photoelectron Spectroscon

Surve for further reading: Surface Analysis by Auger and X-ray Photoelectron Spectroscopy, D Grant, eds., IM Publications and SurfaceSpectra Ltd., 2003 Chectroscopy, Principles and Applications, S. Hüfner, eds. Springer-Verlag Berlin CC 2003.

What is XPS?



The binding energies are characteristic of specific electron orbitals in specific atoms

• XPS lines are identified by the shell from which the electron is emitted



XPS in a nut-shell

- X-ray photoelectron spectroscopy (XPS) is a classical method for the semiquantitative analysis of surface composition
- It is also referred to as Electron Spectroscopy for Chemical Analysis (ESCA)
- It is based on the photoelectric effect, i.e., emission of electron following excitation of core level electrons by photons
- It is surface sensitive because of the low inelastic mean free path of electrons
- An XPS setup consists of a X-ray source, a sample chamber and an electron analyzer
- XPS requires a source of X-rays, i.e., either from a lab-based anode or from a synchrotron
- Traditionally, XPS works only in ultrahigh vacuum because of scattering of electrons in gases
- XPS can also be performed in the mbar pressure range

This is a lab based XPS instrument



Electron energy analyzer

- The most often used type of electron kinetic energy analyzers is composed of an **electrostatic lens** and a **hemispherical analyzer**
- Electrostatic lens decelerates electrons to a **fixed (pass) energy** in the range of a few to 100 eV and at the same time focusses them on an entrance slit of the hemispheric analyzer
- Electrons travel between two concentric hemispheres with a constant potential difference and reach a detector with one dimension aligning a small kinetic energy range
- A spectrum is obtained by sweeping the lens electric field to cover specific kinetic energy ranges of the photoelectrons
- Operation of an electron analyzer requires high vacuum to avoid scattering losses of electrons and to protect the detector



The photoemission process



 $KE = hv - BE - \Phi$ for a solidKE = hv - IPfor a gas Φ : photoelectric workfunction

Fate of core hole



Why is XPS surface sensitive? XPS probe depth



Contribution of atom in depth d to PE peak:

$$I = I_0 \exp\left(\frac{-d}{\lambda \cos \theta}\right)$$

- X-ray photons can penetrate µm but...
- Only photoelectrons from the first layers can escape without energy loss
- Inelastic mean free path (λ) and probing depth strongly depends on the kinetic (and thus photon) energy
- Depth profiles can be obtained either by varying the incident photon energy (tunable x-ray source) or by varying the detection angle (θ)
- Contribution to the photoelectron signal from atoms below the surface decreases exponentially
- In normal emission 95% of the signal comes from a 3λ depth

Electron inelastic mean free path



Inelastic background



- Photoelectrons from deeper layers lose part of their energy (inelastic collisions) and are emitted with reduced KE (> BE)
- XPS spectra show characteristic "stepped" background (intensity of background towards higher BE of photoemission peak is always larger than towards lower BE)



A photoelectron spectrum in more detail





Spin-orbit splitting

- n: principal quantum number
 - I: orbital angular momentum quantum number
 - s: spin angular momentum quantum number
 - j = | I ± s |: total angular momentum quantum number
- For I = 0, **s levels** are singlets, no splitting
- For I > 0, p,d,f levels give rise to doublets. The spin angular momentum of electrons left in an orbital couple with the angular momentum vector
- The degeneracy 2j + 1 determines the possibility for parallel and anti-parallel pairing
- The ratio between the degeneracies (R), (2j₊+1)/(2j₊+1), determines the relative peak ratio of the two peak components.
- ΔE between two components = spin orbit splitting.
- Magnitude of spin-orbit splitting increases with Z and decreases with distance from nucleus (same energy level, ΔE increases with decreasing I)

Core level chemical shifts

- Position of orbitals in atom is sensitive to its chemical environment
- Chemical shift correlated with overall charge on atom (more positive charge = increased BE)
 - number of substituents
 - substituent electronegativity
 - formal oxidation state (depending upon ionicity/covalency of bonding)
- Chemical shift analysis is powerful tool for chemical composition, functional group and oxidation state analysis



Secondary structure of a spectrum

Interaction of the photoemitted electron (along its trajectory) and the remaining electrons: **final state effects (satellite peaks)**

• Shake-up. An electron of the VB can be excited. The energy of this excitation will be deducted from the kinetic energy of the photoelectron. Relevant in metals, in which the valence and conduction bands overlap, empty states are available at very low excitation energies



Simultaneous excitation of a specific wave mode in the sample (e.g. surface plasmons). The kinetic energy loss is hω_s (ω_s is the plasmon frequence), and will repeat at multiples of ω_s

Secondary structure of a spectrum

Multiplet splitting. Arises when an atom contains unpaired electrons (e.g. Cr³⁺, 3p⁶3d³). During the photoemission process, there can be coupling between the unpaired electron in the core with the unpaired electrons in the outer shell. This can create a number of final states, which will be seen in the photoelectron spectrum as a multi-peak envelope. This creates different final states, depending on the orientation of the spin of the unpaired electrons



Estimate of the signal intensity

Contribution of element A at depth d to photoemission signal



- Cross section $\sigma(KE)$ is the probability to have the photoemission event
- 10¹⁵ atoms cm⁻² (equivalent to a monolayer) lead to about 10⁻³ photoelectrons per incident photon
- Typical photon flux: 10¹² s⁻¹, leads to about 10⁹ photoelectrons s⁻¹

Ambient pressure x-ray photoelectron spectroscopy

But...XPS is historically bound to high – ultra-high vacuum

Pressure gap

• Material gap



APXPS has partially filled the gaps!*



- XPS up to 5 mbar (soft x-rays)
 and 50 mbar (tender x-rays)
 - Possibility to investigate "real" samples (powders, semiconductors)

Depth profile

If a **tunable X-ray source** (from a synchrotron) is available, a given electronic level can be excited with **varying photon energies**, resulting in varying photoelectron kinetic energies. Since the **IMFP monotonously increases with increasing kinetic energy** above about 100 eV, XPS can be used to obtain depth profiles of elements or their chemical state.



Oxidation of Au exposing a gold foil (T=373 K) to 0.3 mbar O_3 (1%) in O_2

- New component in the Au 4f spectrum, associated with cationic gold
- The appearance of a O 1s peak confirms that a Au-O bond forms
- Au^{δ+} intensity is maximum at hv=175 eV, and decreases with increasing excitation energy (surface cationic gold)

Depth profile



- The Au^{δ+}/Au⁰ ratio can be plotted vs. kinetic energy (→IMFP)
- It decreases demonstrating that the oxide stays at the surface
- The experimental data can be fitted with a function reproducing the attenuation of the photoemission signal either in a uniform oxide layer (thickness *t* straight line) or in patches of oxide overlayer (dashed line) supported on a semi-infinite metallic support
- *t*=0.3 nm suggests that a O-Au-O trilayer forms at the surface

The solid-gas interface @ NAPP (SLS-PSI)

Static chamber @ NAPP*



- Flow tube configuration
- Heated sample holder. Sample powder pressed, dispersed in ethanol and drop-casted on a gold

foil

Fast heating (e.g. from RT to stable 300°C in a few seconds)

Carbon monoxide oxidation on Pt/CeO₂

Metal NPs (Pt, Pd, Au) supported on reducible oxides showed the best performance in the low temperature oxidation of carbon monoxide^{*}



 For CeO₂: adsorption and activation of oxygen preferentially occur on the support, whereas carbon monoxide is adsorbed and supplied by the metal[#]

- Under catalytically relevant conditions: rapid reoxidation of ceria.
- Ce³⁺ generated in the catalytic cycle is difficult too short lived to be detected under steady state conditions.

Studies employing **time-resolved techniques** are required to gather more information about the evolution of the active sites and their role in the reaction mechanism

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Carbon monoxide oxidation on Pt/CeO₂



The Ce³⁺ concentration increases, thus oxygen from ceria participates in the oxidation of carbon monoxide.

- The initial rate of ceria oxidation is
 >10 times higher than that of reduction. Active Ce³⁺ is short lived
- Presence of spectator Ce³⁺ (not involved in the catalytic process) – probably formed during pretreatment of catalyst.

OPEN ISSUES

- What is the effect of the activation in hydrogen?
- Are the active sites located at the metal/oxide interface? Exploit the surface sensitivity of XPS.

- Pressure in the flow tube stabilized (root pump) at **1.0 mbar**
- Switch on/off the oxygen while acquiring (CO:O₂ = 1:4)
- Oxygen replaced by nitrogen, same pressure and stable m/z signals in ca. 10-15 seconds



Activation of the sample in hydrogen (150°C, 1 hour, p=1.0 mbar)*



 Ce 3d acquired from 2.2 to 5.0 keV photon energy range

^{*}L. Artiglia, et al. J. Phys. Chem. Lett. 2017, 8, 102.

Activation of the sample in hydrogen (150°C, 1 hour, p=1.0 mbar)



• Ce 3d acquired from 2.2 to 5.0 keV photon energy range

The peaks can be separated into 5 *d* doublets[#], associated to different final state configurations:

- v⁰ and v¹ are associated to **Ce³⁺**.
- v, v^{II} and v^{III} are associated to Ce⁴⁺.

^{*}L. Artiglia, et al. J. Phys. Chem. Lett. 2017, 8, 102.

[#]P. Burroughs, et al. J. Chem. Soc., Dalton Trans. 1976, 17, 1686-1698.



- * S. Kato, Phys.Chem.Chem.Phys., 2015, 17, 5078.
- [#] A. Cimino, et al. J. Electron Spectrosc. Relat. Phenom. 1999, 104, 1.

 Enhancement of the Ce³⁺/Ce⁴⁺ ratio at low photon energy

Higher concentration of Ce³⁺ sites is found at the surface^{*}

 On Pt free ceria (red dot) small amount of Ce³⁺ detected only at 2.2.kev

Simple model (exp. distribution) used to fit the experimental data and quantify the depth of reduction of ceria on Pt/CeO₂:[#]

$$\frac{I_o}{I_s} = \frac{F_o \rho_o \sigma_o \lambda_{oo}}{F_s \rho_s \sigma_s \lambda_{ss}} \cdot \frac{[1 - \exp(-t/\lambda_{oo} \cos \theta)]}{\exp(-t/\lambda_{so} \cos \theta)}$$

t = 1.0 nm: Ce³⁺ sites are located in the outer layers of the ceria support, oxygen does not diffuse from the bulk to the

surface.

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 A constant background of Ce³⁺ is visible throughout the experiment **Spectator sites**: formed after the activation, do not participate in the activation of oxygen