

Tensor numerical modeling of the collective electrostatic potentials of many-particle systems

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Abstract

The novel tensor-structured numerical methods appeared as bridging of the algebraic tensor decompositions and the nonlinear approximation theory on separable low-rank representation of multivariate functions and operators [1, 2]. Nowadays the rank-structured tensor approach facilitates new means for numerical modeling of long-range potentials in many-particle systems. The method of grid-based assembled tensor summation of the electrostatic potentials on $(L \times L \times L)$ 3D finite lattices [4] exhibits the computational complexity of the order of $O(L)$ which is much less than $O(L^3)$ in traditional Ewald-type summation. The canonical tensor rank of the collective potential of a finite 3D rectangular lattice system (containing millions of particles) is proven to be as low as a rank for a single reference tensor for a Newton kernel. For lattices with multiple impurities the tensor rank is increased by a small factor [2].

Recent range-separated (RS) tensor format [3] applies to many-particle systems of general type. These can be the free space electrostatic potentials of large bio-molecules or the multidimensional scattered data modeled by radial basis functions. The main advantage of the RS tensor format is that the rank of the canonical/Tucker tensor representing the sum of long range contributions from all particles in the collective potential depends only logarithmically on the number of particles N . Partitioning of long and short range parts of the potentials is performed simply by sorting vectors of the generating Newton kernel. The interaction energies and forces of the many-particle system are computed by using only the long-range part of the collective potential, with representation complexity $O(n \log N)$, where n is the univariate grid size. The basic tool for calculation of the RS tensor representation is the reduced higher order SVD (RHOSVD) introduced in [5]. The representation complexity of the short range part is $O(N)$ with a small prefactor independent on the number of particles. The numerical examples are presented.

References

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