Second order regularity of transition layers in Allen-Cahn equation

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Allen-Cahn equation



One dimensional solution



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Part I. Finite Morse index solutions

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Question: Let u_{ε} be a sequence of solutions to (AC). Assume the level sets $\{u_{\varepsilon} = 0\}$ are uniformly $C^{1,\alpha}$ for some $\alpha \in (0,1)$. Can we get a uniform $C^{2,\alpha}$ regularity?

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This was used in our study on the structure of finite Morse index solutions.

Theorem (W.-Wei '17)

A finite Morse index solution of the Allen-Cahn equation

$$-\Delta u = u - u^3$$

in \mathbb{R}^2 has finitely many ends.

Solutions with finite ends

u looks like the 1d solution along each end. \Leftarrow Refined asymptotics, exponential convergence (Gui '08, Del Pino-Kowalczyk-Pacard '13).



A finiteness result for nodal domains

Let u_e be the directional derivative in *e*-direction, which satisfies the linearized equation

$$\Delta u_e = W''(u)u_e.$$

Lemma

If the Morse index is N, the number of connected components of $\{u_e \neq 0\}$ is at most 2N.

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Lemma

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• Liouville theorem for the degenerate equation

$$\operatorname{div}\left(\varphi^2 \nabla \frac{u_e}{\varphi}\right) = 0.$$

• Similar to Courant's nodal domain theorem: entire space? $n = 2 \implies \text{log cut-off functions, Ambrosio-Cabré '03...}$

Transferring finiteness information

If each end of $\{u = 0\}$ has an asymptotic direction at infinity, finiteness of nodal domains of u_e can be transformed into finiteness of ends.



Theorem

Let u be a finite Morse index solution of the Allen-Cahn equation in \mathbb{R}^2 . For all x large,

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Theorem (Schoen '83)

Given a three dimensional manifold M (with some curvature bounds). Let Σ be a stable immersed minimal surface in a ball $B_R(p) \subset M$ with $\partial \Sigma \subset \partial B_R(p)$. Then

$$\sup_{B_{R/2}(p)\cap\Sigma}|A_{\Sigma}|^2\leq \frac{C}{R^2}.$$

Sternberg-Zumbrun inequality

Stability
$$\Leftrightarrow \int |\nabla \varphi|^2 |\nabla u|^2 \ge \int \varphi^2 \left[|\nabla^2 u|^2 - |\nabla |\nabla u||^2 \right].$$

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If |∇u| ≠ 0,

$$\frac{|\nabla^2 u|^2 - |\nabla|\nabla u||^2}{|\nabla u|^2} = |\mathcal{A}|^2 + |\nabla_T \log |\nabla u||^2,$$

where A is the second fundamental form of level sets $\{u = const.\}$ and ∇_T is the tangential derivatives along these level sets.

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• Simons inequality for this curvature term? Not found yet. Seems to be a common difficulty in semilinear problems.

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- Find y_k satisfying

$$|A(y_k)| \ge |A(x_k)|, \qquad |A(y_k)||y_k| \ge k,$$

$$|A(x)| \leq 2|A(y_k)|, \quad \forall x \in B_{k|A(y_k)|^{-1}}(y_k).$$

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• Let $\varepsilon_k := |A(y_k)|$ and define $u_k(x) := u(y_k + \varepsilon_k^{-1}x)$. $|y_k| \to +\infty$ and $\varepsilon_k \to 0 \iff$ Locally close to 1D solution, by stable De Giorgi for n = 2.

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- In $B_k(0)$, u_k is a stable solution of (AC) with parameter ε_k .
- The curvature of $\{u_{\varepsilon_k} = 0\}$ is uniformly bounded, and it equals 1 at the origin.

Theorem (W.-Wei '18)

Let u_{ε} be a sequence of stable solutions to (AC) such that $\{u_{\varepsilon} = 0\}$ are uniformly $C^{1,\beta}$ for some $\beta \in (0,1)$. If $n \leq 10$, then $\{u_{\varepsilon} = 0\}$ are uniformly bounded in $C^{2,\alpha}$ for any $\alpha \in (0,1)$. Moreover, the mean curvature is of the order $O(\varepsilon^{\alpha})$.

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Chodosh-Mantoulidis '18 has obtained the same result in dimension 3, which was used in their study of min-max minimal surfaces in three manifolds (Multiplicity one conjecture of Marques-Neves, existence of infinitely many minimal surfaces in generic case). In general the stability condition is necessary \Longleftarrow Clustering interfaces, Toda system.

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Theorem

Let u_{ε} be a sequence of the Allen-Cahn equation in $B_1(0)$, with $\{u_{\varepsilon} = 0\}$ given by the graph of a uniformly $C^{1,\beta}$ functions f_{ε} for some $\beta \in (0,1)$. Then f_{ε} are uniformly bounded in $C^{2,\alpha}_{loc}(B_1^{n-1})$ for any $\alpha \in (0,1)$.

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Caffarelli and Córdoba '06 and Savin '09: Lipschitz or flat \implies uniform $C^{1,\alpha}$.

Part II. Proof of $C^{2,\alpha}$ estimates

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Clustering interfaces

There could be more and more connected components of $\{u_{\varepsilon} = 0\}$, which can collapse to the same limit as $\varepsilon \to 0$.



Curvature bound on $\{u_{\varepsilon} = 0\} \Longrightarrow \{u_{\varepsilon} = 0\}$ locally represented by graphs $\cup_k \Gamma_{k,\varepsilon}$, where

$$\Gamma_{k,\varepsilon} = \{x_2 = f_{k,\varepsilon}(x_1)\}, \quad \cdots < f_{k-1,\varepsilon} < f_{k,\varepsilon} < f_{k+1,\varepsilon} < \cdots.$$

The cardinality of index set could go to infinity.

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Interaction between different interfaces has the form

$$\operatorname{div}\left(\frac{\nabla f_{k,\varepsilon}}{\sqrt{1+|\nabla f_{k,\varepsilon}|^2}}\right) = \frac{A}{\varepsilon} \left[e^{-\frac{\sqrt{2}}{\varepsilon} \left(f_{k,\varepsilon} - f_{k-1,\varepsilon}\right)} - e^{-\frac{\sqrt{2}}{\varepsilon} \left(f_{k+1,\varepsilon} - f_{k,\varepsilon}\right)} \right] + h.o.t.$$

Infinite dimensional Lyapunov-Schmidt reduction of Del Pino, Kowalczyk and Wei.

Obstruction to $C^{2,\alpha}$ estimates of $f_{k,\varepsilon}$

$$\Delta f_{k,\varepsilon} = \frac{A}{\varepsilon} e^{-\frac{\sqrt{2}}{\varepsilon} \left(f_{k,\varepsilon} - f_{k-1,\varepsilon} \right)} - \frac{A}{\varepsilon} e^{-\frac{\sqrt{2}}{\varepsilon} \left(f_{k+1,\varepsilon} - f_{k,\varepsilon} \right)}.$$
$$f_{k+1,\varepsilon} - f_{k,\varepsilon} \ge \frac{\sqrt{2} \left(1 + \alpha \right)}{2} \varepsilon |\log \varepsilon| - C\varepsilon \Longrightarrow f_{k,\varepsilon} \in C^{2,\alpha}.$$

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Obstruction to $C^{2,\alpha}$ estimates of $f_{k,\varepsilon}$

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On the other hand, if

$$f_{k+1,\varepsilon} - f_{k,\varepsilon} \leq \frac{\sqrt{2}}{2} \varepsilon |\log \varepsilon| + C \varepsilon,$$

define the blow up sequence

$$\widetilde{f}_{k,arepsilon}(x):=rac{1}{arepsilon}f_{k,arepsilon}\left(arepsilon^{rac{1}{2}}x
ight)-rac{\sqrt{2}lpha}{2}|\logarepsilon|.$$

They converge to an entire solution of the Toda system

$$\Delta f_k = e^{-\sqrt{2}(f_k - f_{k-1})} - e^{-\sqrt{2}(f_{k+1} - f_k)}, \quad \text{in } \mathbb{R}^{n-1}.$$

Example I: Two end solutions of Agudelo-del Pino-Wei

For n ≥ 10, the Liouville equation (= Two component Toda system)

$$\Delta f = e^{-\sqrt{2}f}$$

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- For λ small ($\iff f_{\lambda}(0) \gg 1$), there exists a monotone (in λ) family of solutions u^{λ} to the unscaled (AC). \implies stable.
- Let $u_{\varepsilon}(x) := u^{\varepsilon^{1/2}} (\varepsilon^{-1}x)$. Its nodal set $\{u_{\varepsilon} = 0\}$ is given by the graph of

$$f_{\varepsilon}(x) \approx \varepsilon f\left(\varepsilon^{-rac{1}{2}}x
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which satisfies

$$|
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eq 0.$$

Example II: Multiple end solutions in \mathbb{R}^2

Del Pino-Kowalczyk-Pacard-Wei '10 : Unstable solutions with $\{u = 0\}$ close to the graph of Toda solutions:

 $f_k''(x) = e^{-\sqrt{2}(f_k(x) - f_{k-1}(x))} - e^{-\sqrt{2}(f_{k+1}(x) - f_k(x))}, \quad x \in \mathbb{R}, \quad 1 \le k \le Q.$



$$\sum_{k} \int |\nabla \eta_{k}|^{2} \geq \frac{\sqrt{2}A}{\varepsilon^{2}} \sum_{k} \int (\eta_{k} - \eta_{k-1})^{2} e^{-\frac{\sqrt{2}}{\varepsilon} \left(f_{k,\varepsilon} - f_{k-1,\varepsilon}\right)} - h.o.t..$$

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• Uniform $C^{2,\alpha}$ estimates of clustering interfaces does not hold

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 ⇒ Existence of entire stable solutions of Toda system

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Uniform C^{2,α} estimates of clustering interfaces does not hold
 ⇒ Existence of entire stable solutions of Toda system
 Liouville theorem: No stable solution of Toda in ℝⁿ for n ≤ 9.

Theorem

For any n, there exists a universal constant η such that, if (f_k) is a stable solution to the Toda lattice

$$\Delta f_k = e^{-\sqrt{2}(f_k - f_{k-1})} - e^{-\sqrt{2}(f_{k+1} - f_k)}$$
 in $B_1 \subset \mathbb{R}^n$,

then

$$\int_{B_1} e^{-\sqrt{2}(f_k - f_{k-1})} \leq \eta(n) \implies \sup_{B_{1/2}} e^{-\sqrt{2}(f_k - f_{k-1})} \leq \frac{1}{2}.$$

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Applying this ε -regularity to suitable rescalings of Toda system constructed from (AC), gives a decay estimate on $e^{-\frac{\sqrt{2}}{\varepsilon}(f_{k,\varepsilon}-f_{k-1,\varepsilon})}$ in shrinking balls, leading finally to

$$e^{-rac{\sqrt{2}}{arepsilon}(f_{k,arepsilon}-f_{k-1,arepsilon})}\lesssimarepsilon^{1+lpha}, \quad ext{ in the interior}.$$

Without the stability condition, for the Liouville equation

$$\Delta f = e^{-f},$$

this ε -regularity has been proved by Brezis-Merle '91 in 2 dimension and F. Da Lio '08 in 3 dimension. Higher dimensions are not known.

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Our proof relies essentially on the stability condition, which gives us an integral estimate (as in Farina '07)

$$\int_{B_r} e^{-2\sqrt{2}(f_k-f_{k-1})} \leq Cr^{-2} \int_{B_{2r}} e^{-\sqrt{2}(f_k-f_{k-1})}.$$

Part III. Derivation of the Toda system

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$$G(x_1, x_2) = 0 \iff \begin{cases} \prod_{F} \circ G(x_1, x_2) = 0, \\ \prod_{F^{\perp}} \circ G(x_1, x_2) = 0. \end{cases}$$

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- Substituting the second equation into the first one \implies an equation defined on *E*.

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- If *E* is finite dimensional: finite dimensional reduction.

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- Solve the second equation by Implicit Function Theorem.
- Substituting the second equation into the first one \implies an equation defined on *E*.
- If *E* is finite dimensional: finite dimensional reduction.
- Even if *E* is infinite dimensional, sometimes the reduction problem is still a good one.

• Finite dimensional reduction method has been used by many authors to construct solutions of nonlinear PDEs: gluing method.

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- We already have a solution u_{ε} of the Allen-Cahn equation.
- In order to get estimates on level sets, a good equation satisfied by these level sets is needed.
- Decouple the single equation (AC) into two: one is the equation for level sets (on the kernel space E), the other one is on E[⊥] which enjoys good a priori estimates.

Starting assumptions: close to a canonical configuration by assuming $C^{1,1}$ regularity.

Around $\Gamma_{\varepsilon} := \{u_{\varepsilon} = 0\}$, u_{ε} looks like

$$g_*(y,z) := g\left(\frac{z-h_{\varepsilon}(y)}{\varepsilon}\right),$$

where $z = \text{dist}_{\Gamma_{\varepsilon}}$ is the signed distance to Γ_{ε} and y denotes a point on Γ_{ε} . \implies Introduction of Fermi coordinates w.r.t. Γ_{ε} .

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A small perturbation in normal direction is also needed: introduction of h_{ε} .

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$$\Delta = \Delta_z - H\partial_z + \partial_{zz},$$

 Δ_z is the Beltrami-Laplace operator on {dist_{$\Gamma_{\varepsilon}} = z$ }. *H* is the mean curvature of {dist_{$\Gamma_{\varepsilon}} = z$ }.</sub></sub> Fermi coordinates: x = y + zN(y), where $y \in \Gamma_{\varepsilon}$ is the nearest point and N the unit normal vector, $z = \text{dist}_{\Gamma_{\varepsilon}}$. Well defined in an $O(\delta)$ neighborhood of Γ_{ε} , where δ depends only on $\sup |A_{\Gamma_{\varepsilon}}|$.

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$$\Delta d = \sum_{i} \frac{k_i}{1 - k_i d} = H + O(|A|^2),$$

where k_i are principal curvatures of Γ_{ε} .

 h_{ε} is a function defined on Γ_{ε} , which must be introduced so that the orthogonal condition holds:

$$\int_{-\infty}^{+\infty} \left(u_{\varepsilon}(y,z) - g_*(y,z)
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For each $y \in \Gamma_{\varepsilon}$,

$$\min_{t\in\mathbb{R}}\int_{-\infty}^{+\infty}\left|u_{\varepsilon}(y,z)-g\left(\frac{z-t}{\varepsilon}\right)\right|^{2}dz$$

is attained at a unique point $h_{\varepsilon}(y)$.

The error between $u_arepsilon$ and g_* , $\phi_arepsilon:=u_arepsilon-g_*$, satisfies

$$\varepsilon^{2}\left(\Delta_{0}\phi_{\varepsilon}+\frac{\partial^{2}\phi_{\varepsilon}}{\partial z^{2}}\right)=W''(g_{*})\phi_{\varepsilon}+\varepsilon\left[H_{\varepsilon}(y)+\Delta h_{\varepsilon}(y)\right]g'_{*}+h.o.t.$$

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• ϕ_{ε} is controlled by h.o.t. which are quadratic in ϕ_{ε} and $H_{\varepsilon} \Longrightarrow$ decay estimate starting from a good position.

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The parallel part gives the equation satisfied by {u_ε = 0}:

$$H_{\varepsilon}=O\left(\varepsilon
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Then standard elliptic estimates on minimal surface equation gives the $C^{2,\alpha}$ estimate.

$$\begin{cases} \Delta_{\mathbb{R}^{n-1}\phi} + \partial_{zz}\phi = W''(g(z))\phi + a(y)g'(z) + E, \\ \int_{-\infty}^{+\infty} \phi(y,z)g'(y,z)dz = 0, \quad \forall y \in \mathbb{R}^{n-1}. \end{cases}$$

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Main tool: Nondegeneracy of 1D solution, i.e. kernel of the linearized operator is spanned by g'.

Spectral gap: the second eigenvalue is positive \Longrightarrow

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$$a(y)\int_{-\infty}^{+\infty}g'(z)^2dz=-\int_{-\infty}^{+\infty}E(y,z)g'(z)dz.$$
Starting assumptions: close to a canonical configuration. Around $\Gamma_{k,\varepsilon}$, u_{ε} looks like

$$g_{k,arepsilon} := g\left(rac{{\mathsf{dist}}_{{\mathsf{\Gamma}}_{k,arepsilon}}-h_{k,arepsilon}}{arepsilon}
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where dist_{\Gamma_{k,\varepsilon}} is the signed distance to $\Gamma_{k,\varepsilon}$. \Longrightarrow As before, we need to

- introduce Fermi coordinates w.r.t. $\Gamma_{k,\varepsilon}$;
- introduce a small perturbation h_{k,ε} to obtain the orthogonal condition.

Interaction between transition layers

Approximate solution: near Γ_k ,



Decoupling in clustering interfaces

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 The parallel part (w.r.t. g'_{k,ε}) gives the Toda system with remainder terms of higher order (quadratic in φ_ε, e^{-√2/ε}(f_{α+1,ε}-f_{α,ε}) etc.)

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- The parallel part (w.r.t. $g'_{k,\varepsilon}$) gives the Toda system with remainder terms of higher order (quadratic in ϕ_{ε} , $e^{-\frac{\sqrt{2}}{\varepsilon}(f_{\alpha+1,\varepsilon}-f_{\alpha,\varepsilon})}$ etc.)
- ϕ is mainly controlled by the interaction term:

$$\|\phi_{\varepsilon}\|_{\mathcal{C}^{2,\alpha}} \lesssim \varepsilon^{1-\alpha} + \frac{1}{\varepsilon} \sup\left[\mathrm{e}^{-\frac{f_{k+1,\varepsilon} - f_{k,\varepsilon}}{\varepsilon}} + \mathrm{e}^{-\frac{f_{k,\varepsilon} - f_{k-1,\varepsilon}}{\varepsilon}} \right].$$

Reduction of the stability condition again

• In the stability condition for the Allen-Cahn equation,

$$\int \varepsilon |\nabla \psi|^2 \geq \int \frac{1}{\varepsilon} W''(u_{\varepsilon}) \psi^2,$$

choose (as in Agudelo-Del Pino-Wei '16)

$$\psi := \sum_{k} \eta_{k} g'_{k,\varepsilon},$$

where $\eta_k \in C_0^{\infty}(\Gamma_{k,\varepsilon})$.

• Good decomposition: $u_{\varepsilon} = g_* + \phi_{\varepsilon}$, i.e. good estimates on ϕ_{ε} \implies main order terms are

$$\sum_{k} \int |\nabla \eta_{k}|^{2} \geq \frac{\sqrt{2}A}{\varepsilon^{2}} \sum_{k} \int (\eta_{k} - \eta_{k-1})^{2} e^{-\frac{\sqrt{2}}{\varepsilon} \left(f_{k,\varepsilon} - f_{k-1,\varepsilon}\right)}.$$

• Reduction method is basically a (partial) linearization procedure: when the solution u_{ε} is close to $\sum_{i} g_{i,\varepsilon}$, the ϕ_{ε} equation is almost a linearized one;

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- The reduced problem is still a nonlinear one in most cases, but its complexity is decreased;
- Long range interaction in these phase field models. This is helpful for the construction of Jacobi fields (Chodosh-Mantoulidis '18).

Thanks for your attention!