

# Parallels between Earthquakes, Financial crashes and epileptic seizures



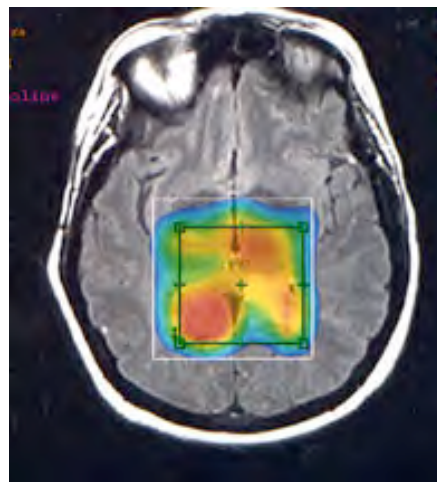
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D. Sornette

## Critical Phenomena in Natural Sciences

Chaos, Fractals,  
Selforganization and Disorder:  
Concepts and Tools

**First edition  
2000**

**Second  
enlarged edition  
2004**

 Springer

**DIDIER SORNETTE**

Princeton  
University  
Press  
Jan. 2003



Critical Events in  
Complex Financial Systems

Malevergne · Sornette



Extreme Financial Risks

Y. Malevergne  
D. Sornette

# Extreme Financial Risks

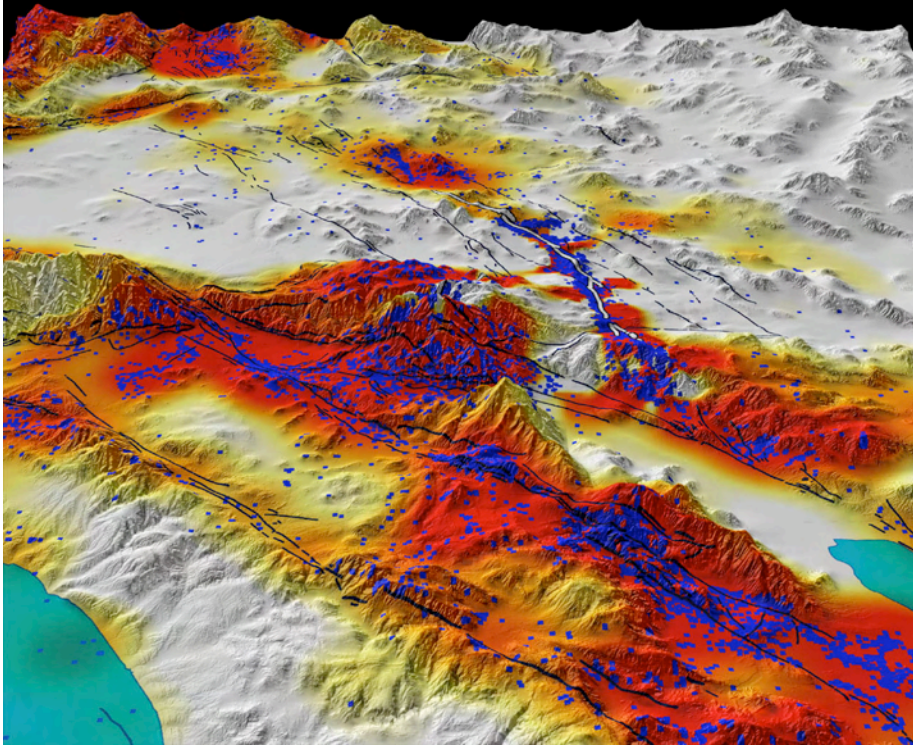
From Dependence  
to Risk Management

(November 2005)

 Springer

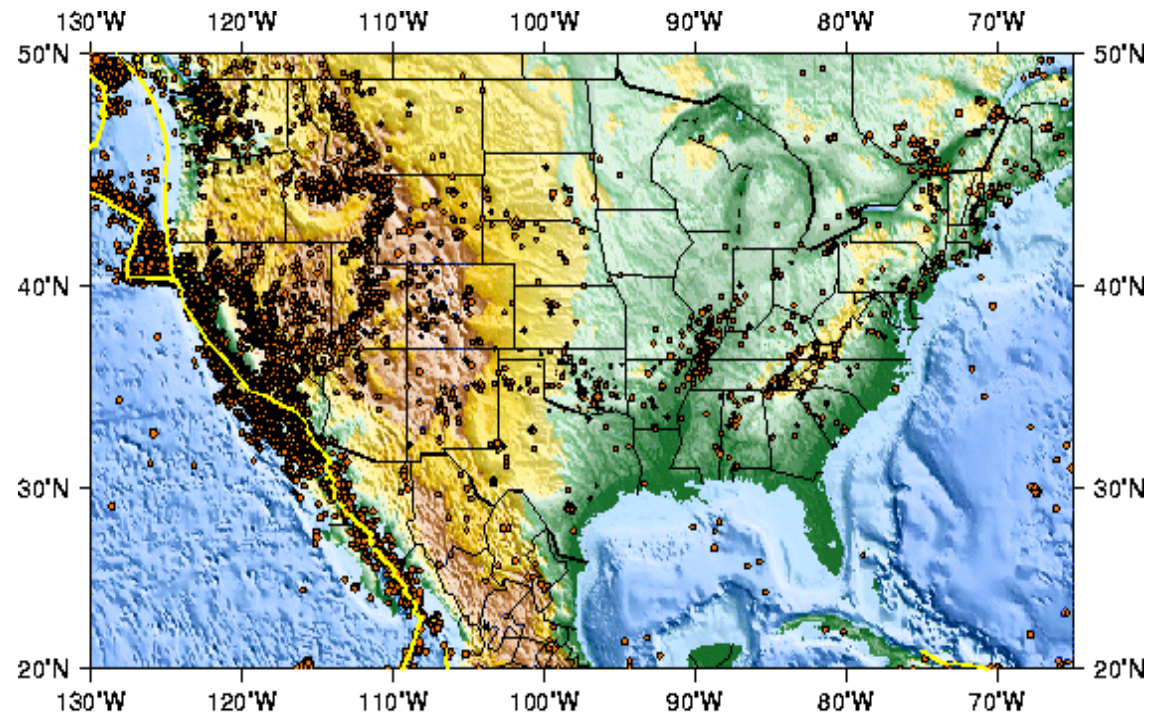
# Earthquake Conversations

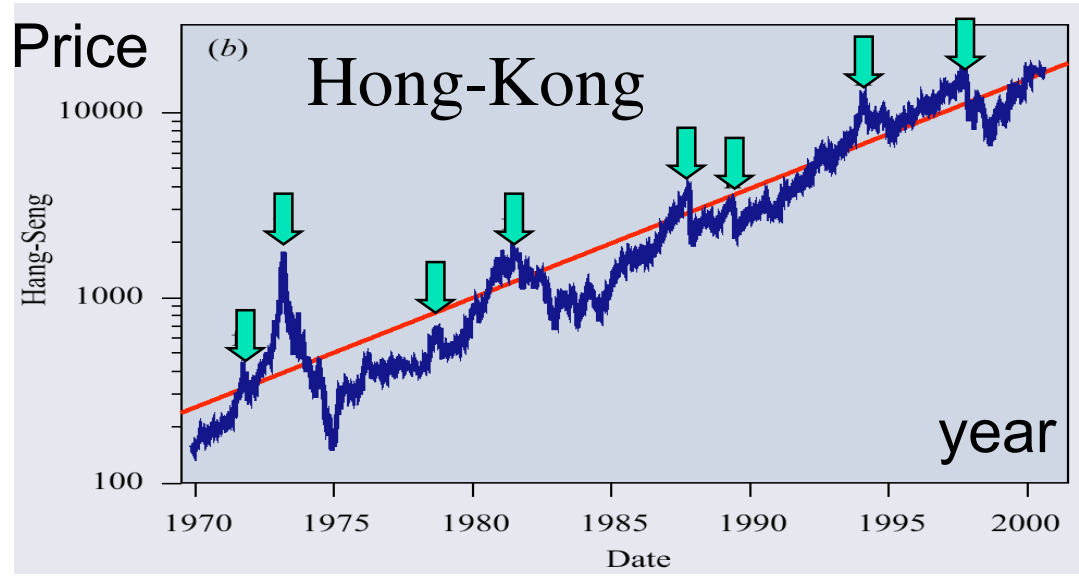
Ross S. Stein  
U.S. Geological Survey



Epidemic processes by  
word-of-mouth,  
sentiment, convention...

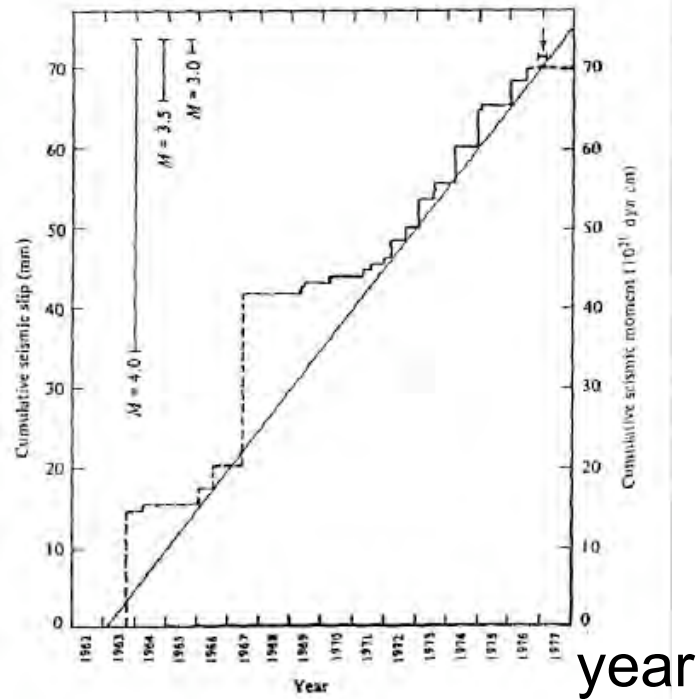






Red line is 13.8% per year: but The market is never following the average growth; it is either super-exponentially accelerating or crashing

### Cumulative Slip



Cumulative moment and seismic slip in a zone of the Calaveras fault (1962-77)

## Statistical laws of seismicity

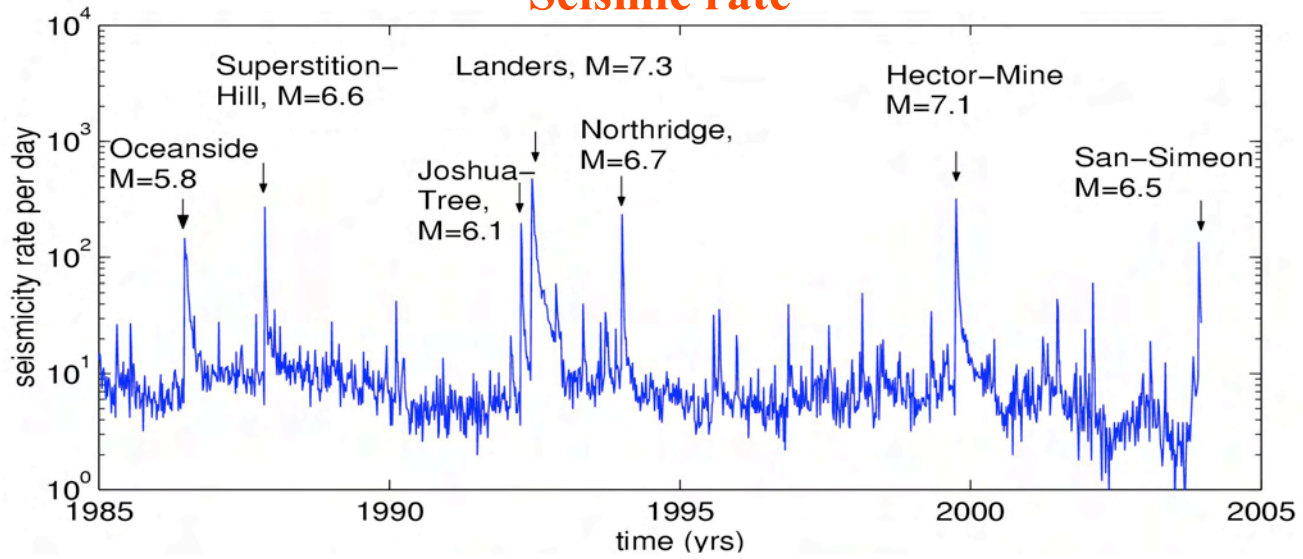
- Gutenberg-Richter law:  $\sim 1/E^{1+\beta}$  (with  $\beta \approx 2/3$ )
- Omori law  $\sim 1/t^p$  (with  $p \approx 1$  for large earthquakes)
- Productivity law  $\sim E^a$  (with  $a \approx 2/3$ )
- PDF of fault lengths  $\sim 1/L^2$
- Fractal/multifractal structure of fault networks  $\zeta(q), f(\alpha)$
- PDF of seismic stress sources  $\sim 1/s^{2+\delta}$  (with  $\delta \geq 0$ )
- Distribution of inter-earthquake times
- Distribution of seismic rates

# Stylized facts of financial markets

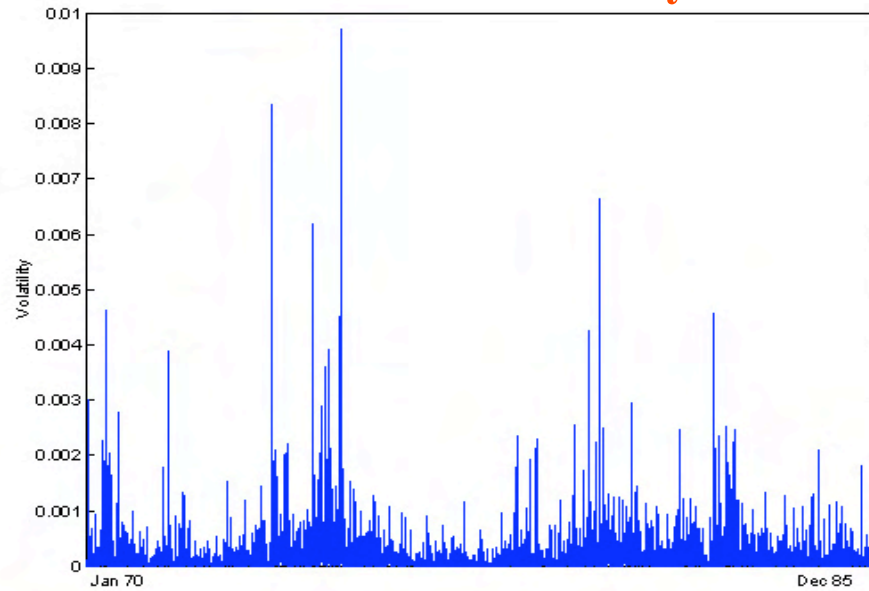
- Heavy-tail pdf of returns
- Omori law and Long-memory of volatility
- Price impact function      Price  $\sim V^\beta$  with  $\beta=0.2-0.6$
- Pareto distribution of wealth
- Multifractal structure of returns
- PDF of news' sizes?
- Distribution of inter-shock times
- Distribution of limit order sizes
- “Leverage” effect

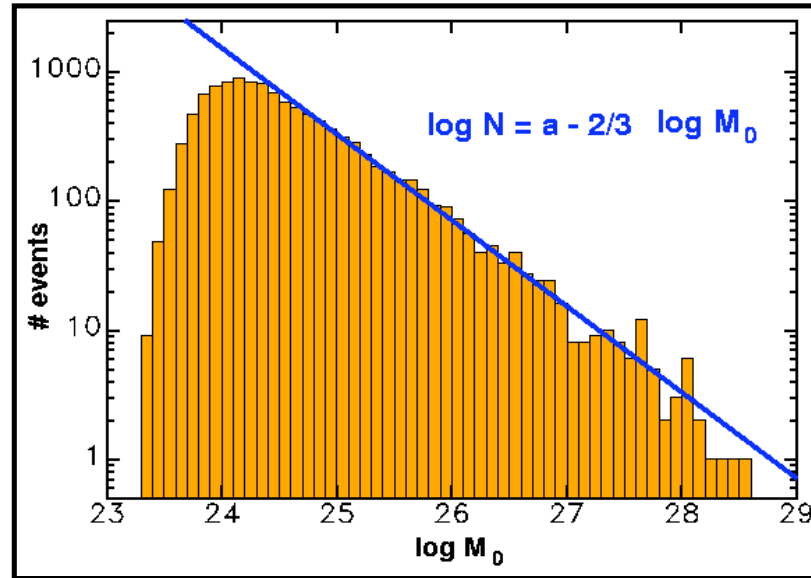


## Seismic rate

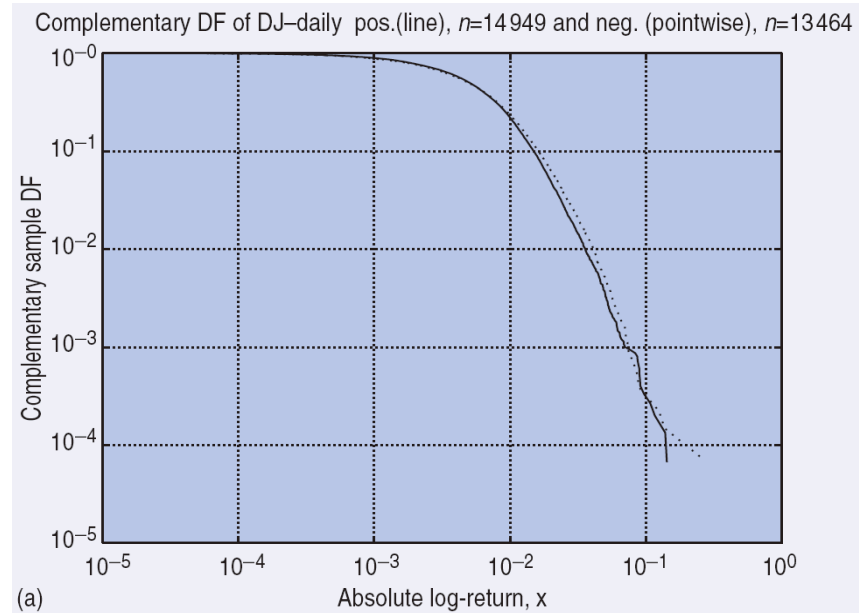


## Financial Volatility

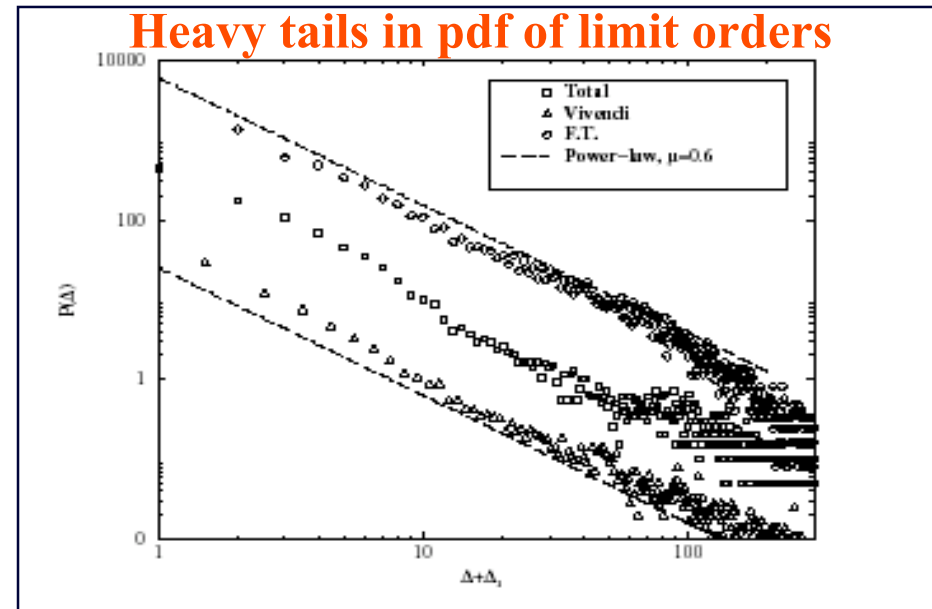
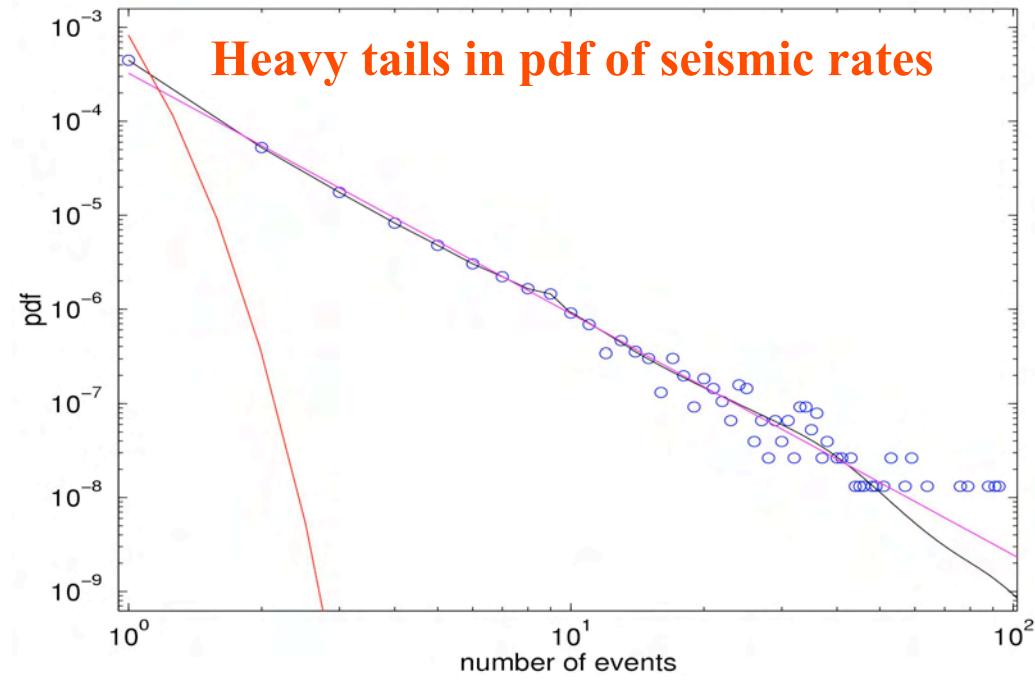




**Heavy tails in  
pdf of earthquakes  
 $b=2/3$**

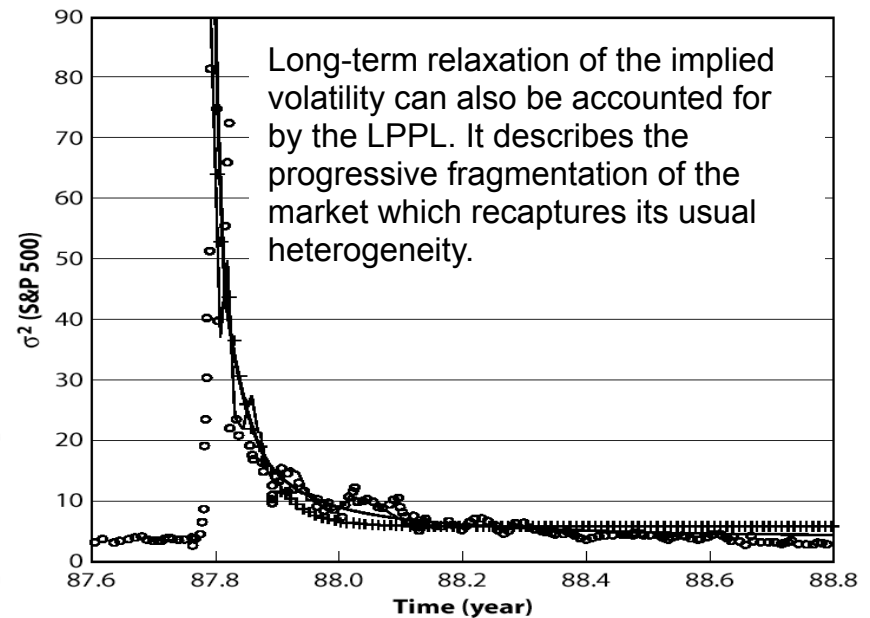
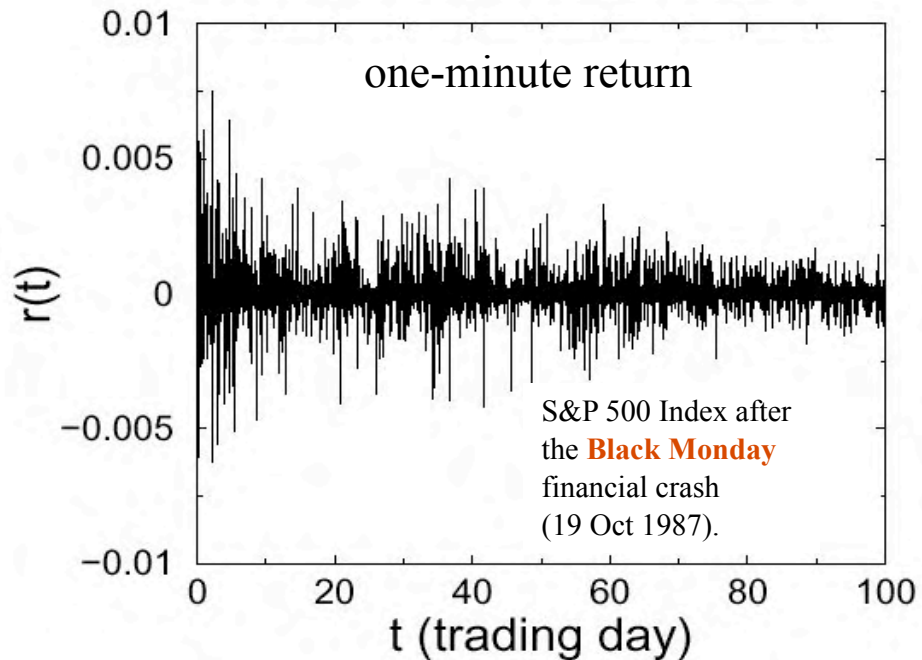
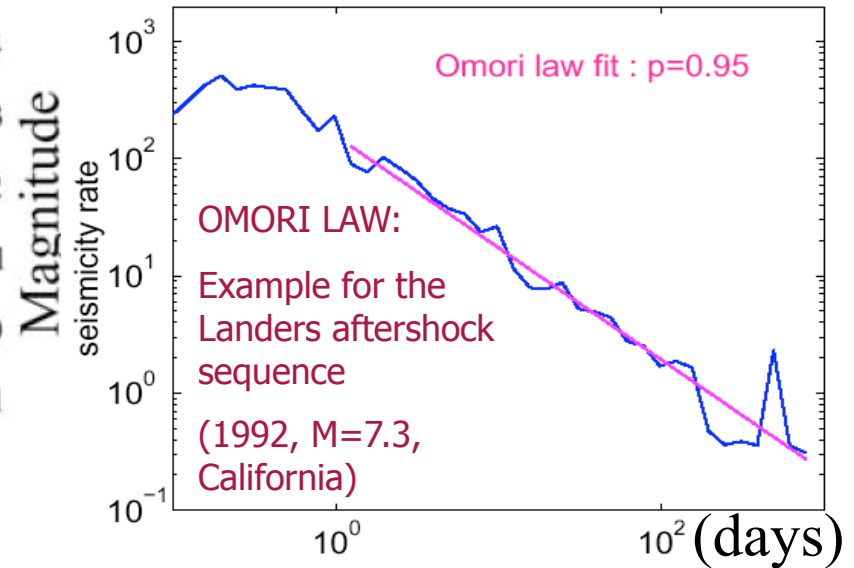
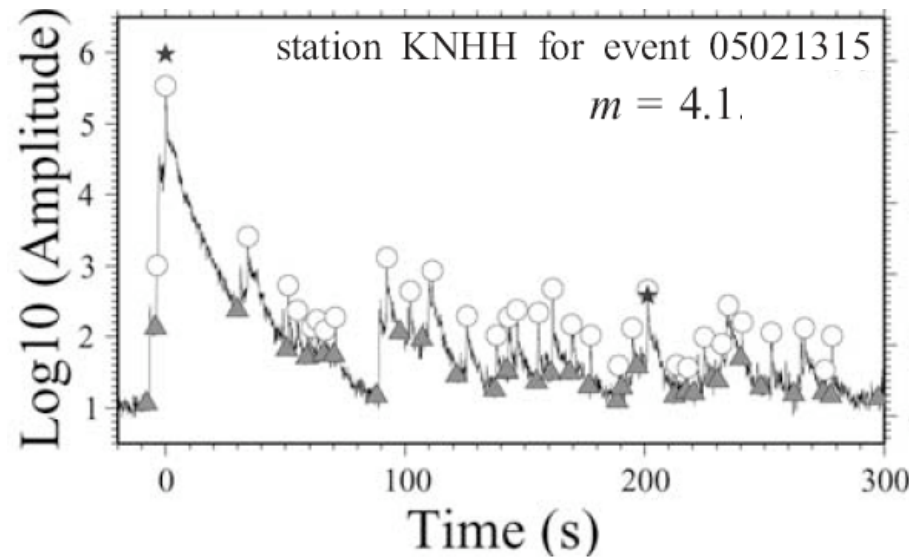


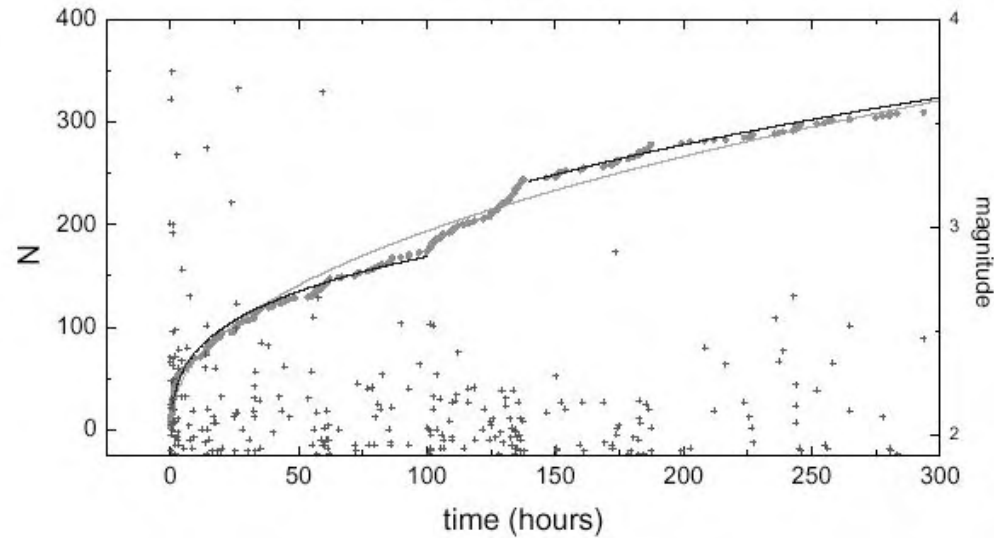
**Heavy-tails of  
price changes  
 $b=3$**



Peng et al.

JOURNAL OF GEOPHYSICAL RESEARCH, VOL. 112, B03306, doi:10.1029/2006JB004386, 2007

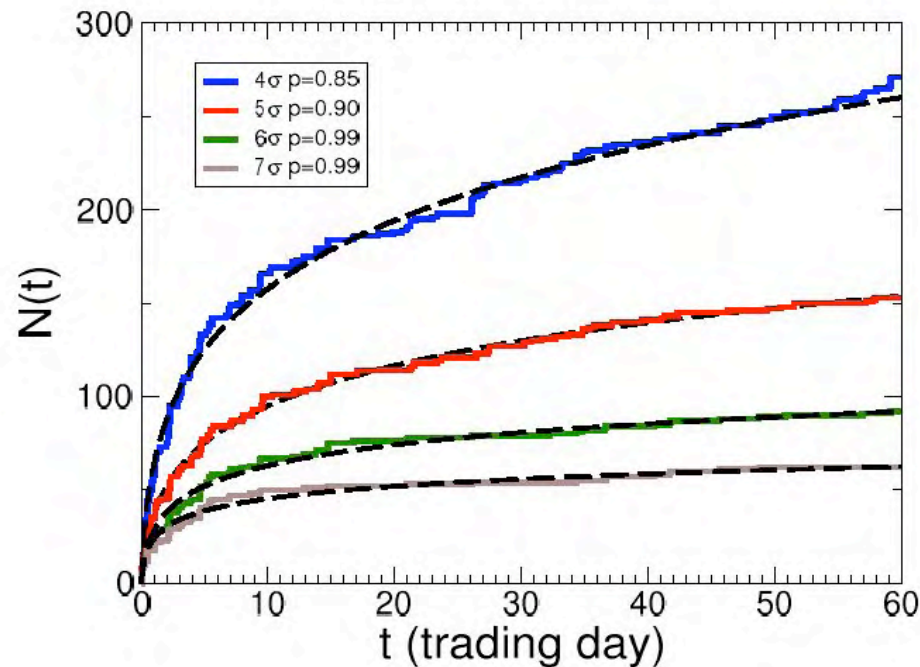




Cumulative number of aftershocks in the earthquake occurring in eastern Pyrenees on February 18, 1996 (from Moreno *et al.*, J. of Geophys. Res., **106 B4**, 6609-6619 (2001))

$$n(t) \propto t^{-p}; \quad N(t) = \int_0^t n(s) ds$$

1987



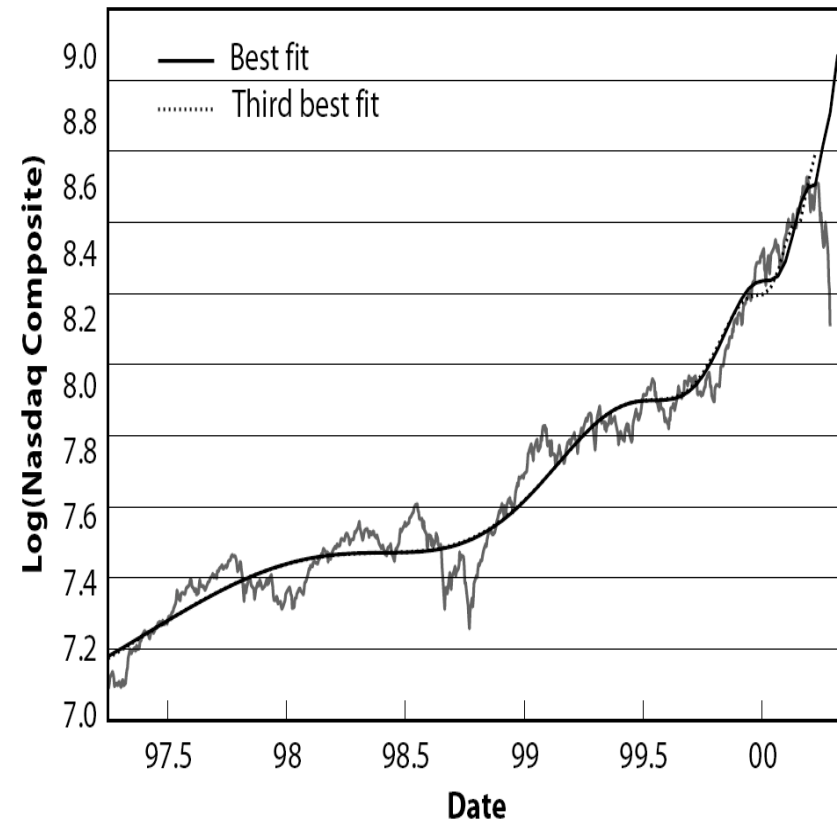
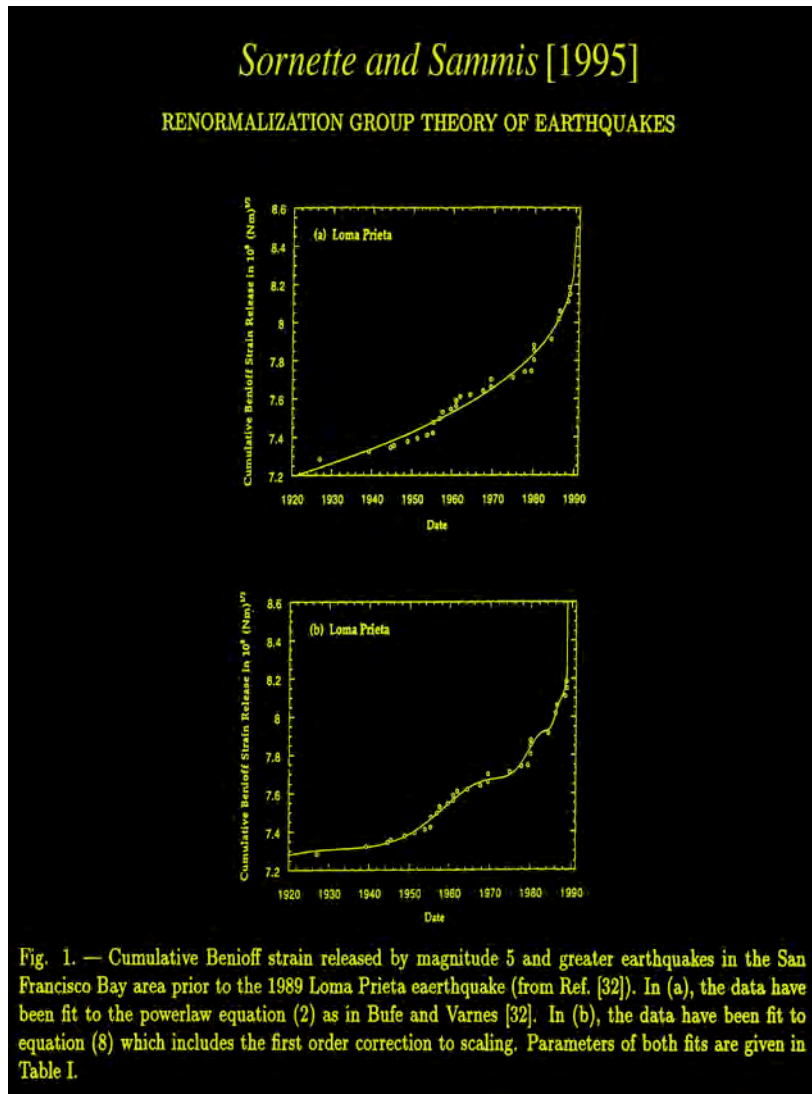
$$N(t) = K[(t + \tau)^{1-p} - \tau^{1-p}] / (1-p)$$

Oct. 1987 crash:  
Cumulative number of S&P500 index returns exceeding a given threshold  $n\sigma$

†Lillo and Mantegna, PRE **68**, 016119 (2003)

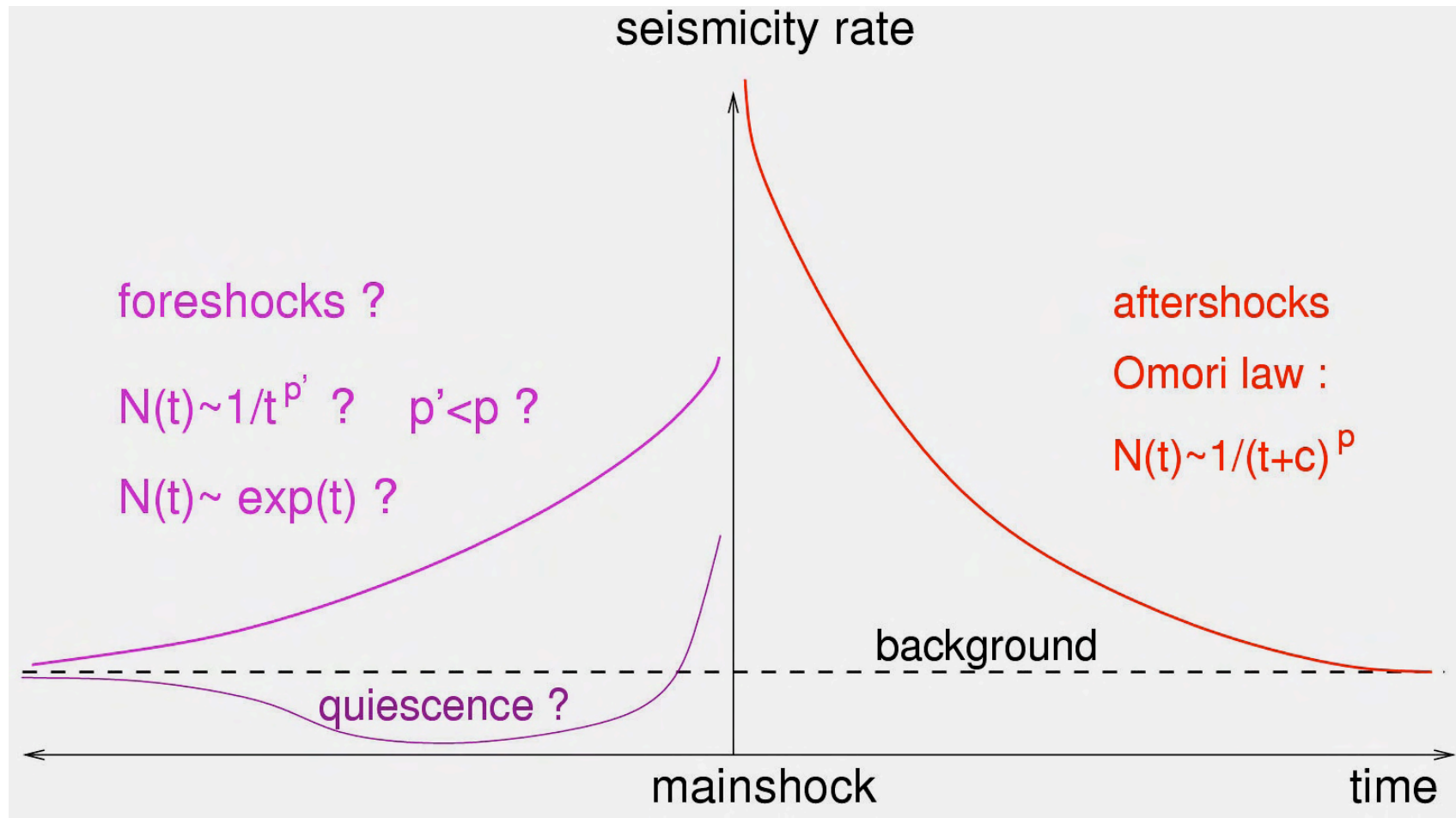
# Critical earthquakes? Critical crashes?

## THE NASDAQ CRASH OF APRIL 2000

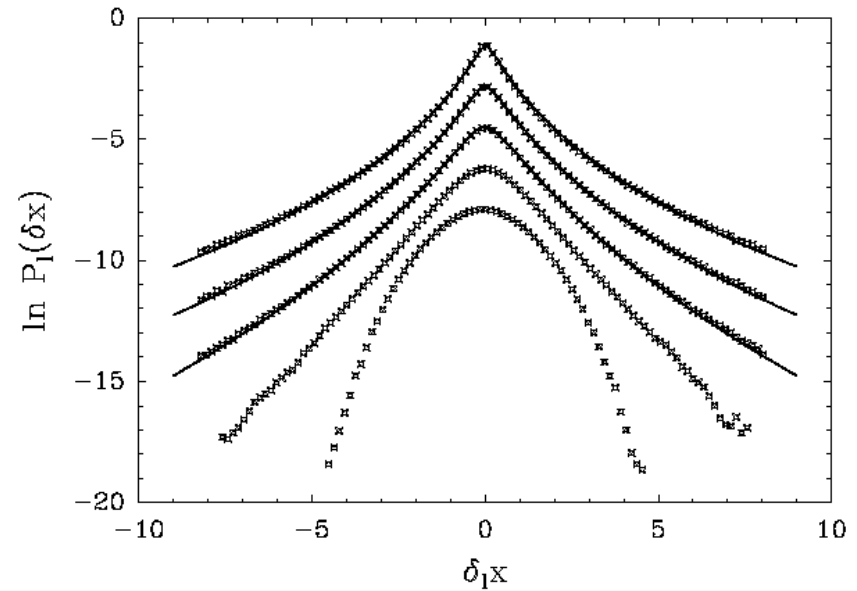
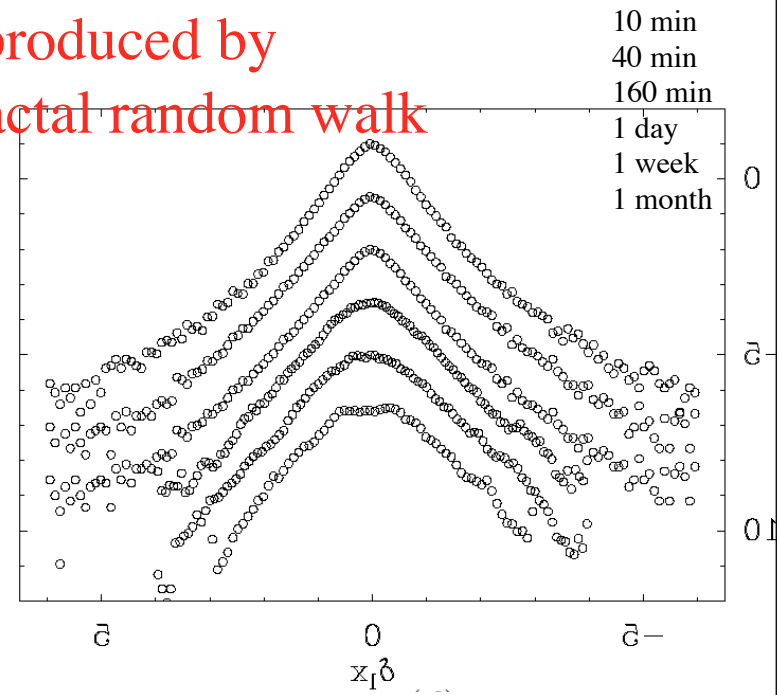
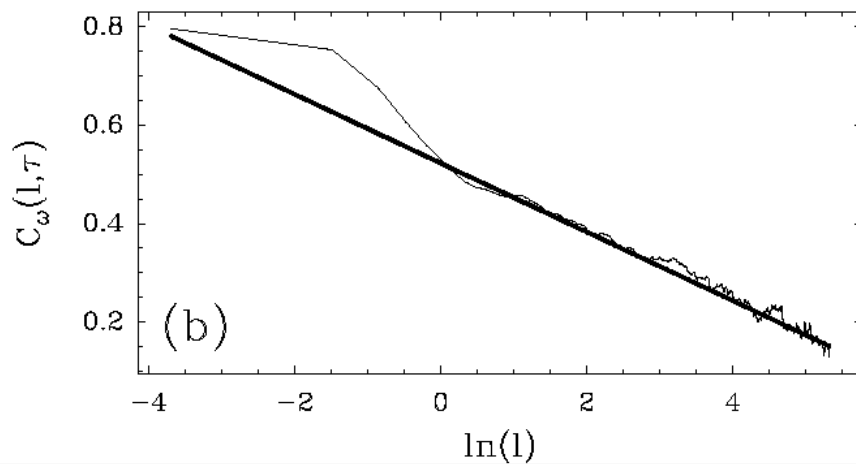
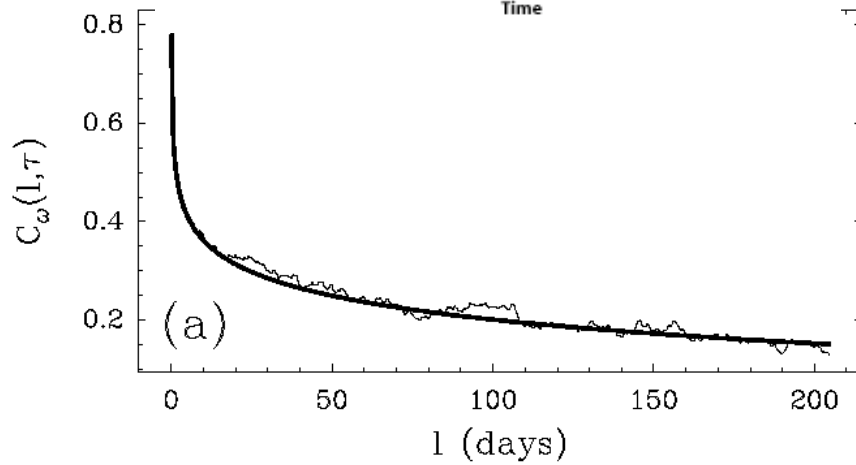
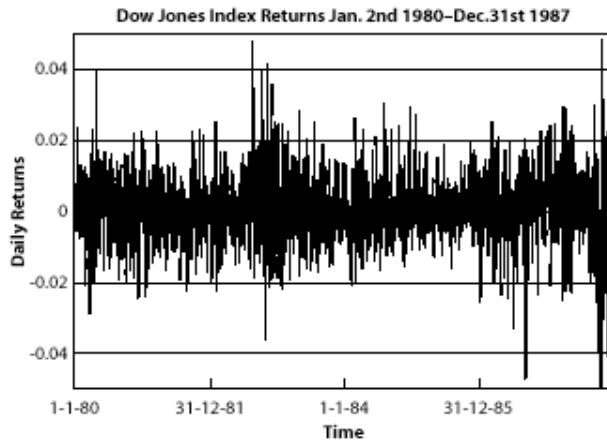


# Generic multifractality

## Endogenous versus Exogenous responses

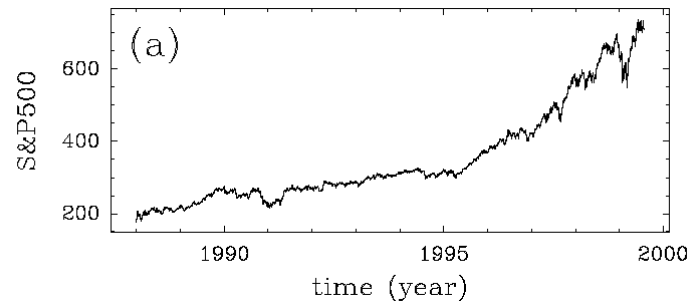


Stylized facts in financial markets  
 well-reproduced by  
 Multifractal random walk

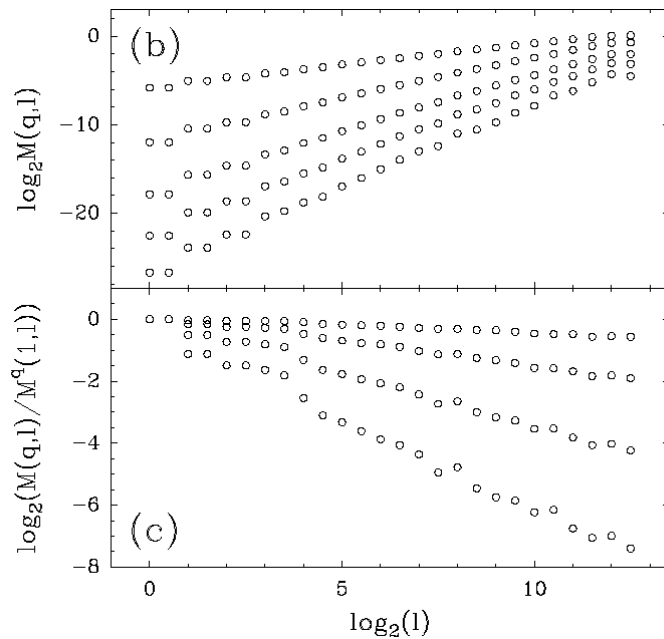




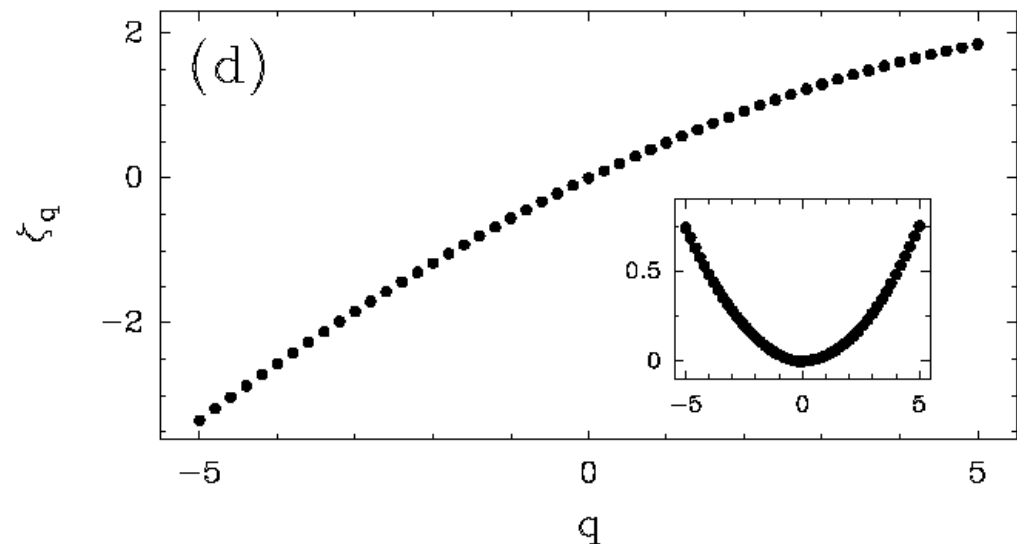
Multifractality:  $\langle [\delta_\tau X(t)]^q \rangle = a(q) \tau^{\zeta(q)}$ , for  $\tau < T$ .

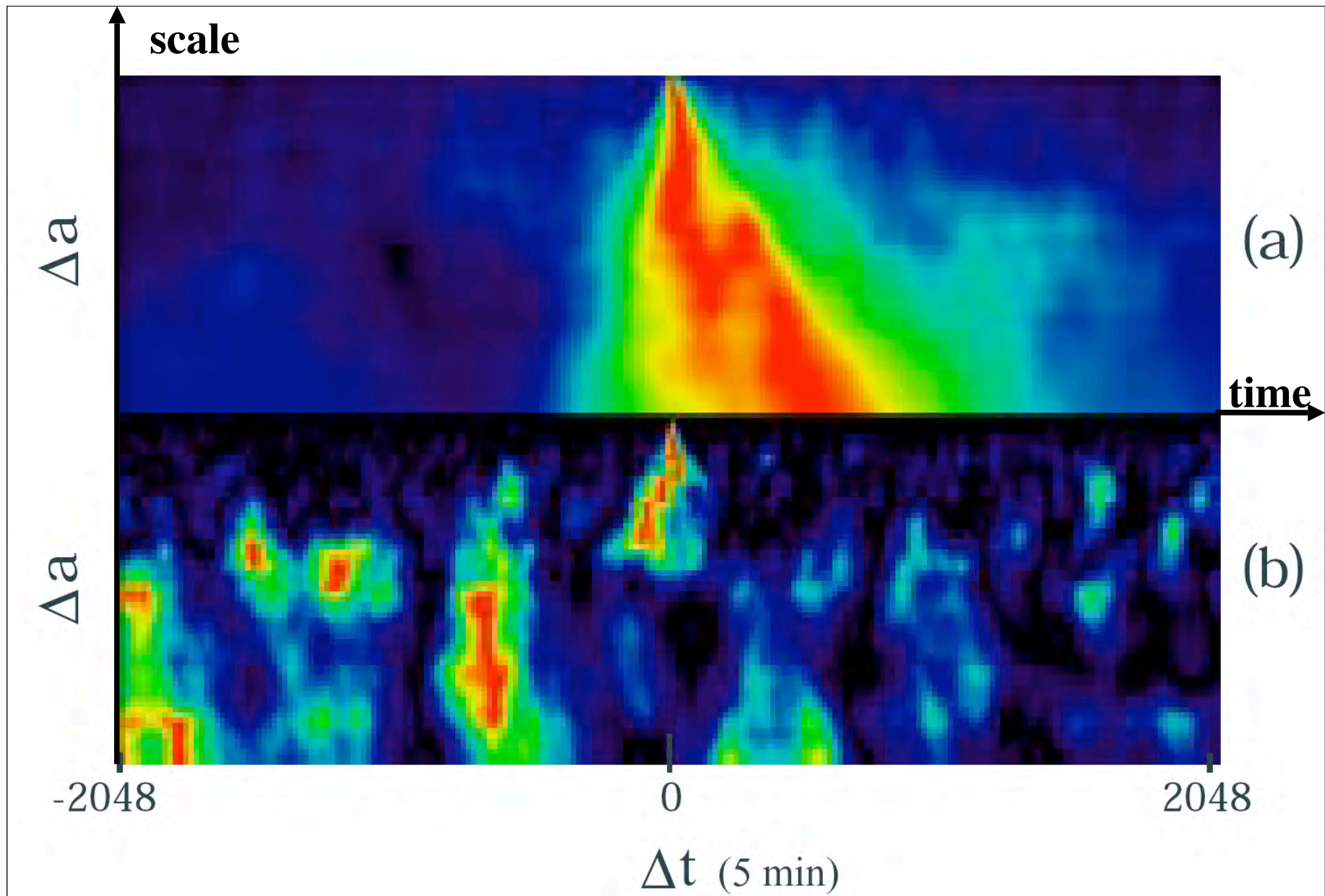


$$\zeta(q) = \left(1 + \frac{\lambda^2}{2}\right)q - \frac{\lambda^2}{2}q^2$$



$\lambda^2 = -\zeta''(0)$  is the so-called *intermittency coefficient*.

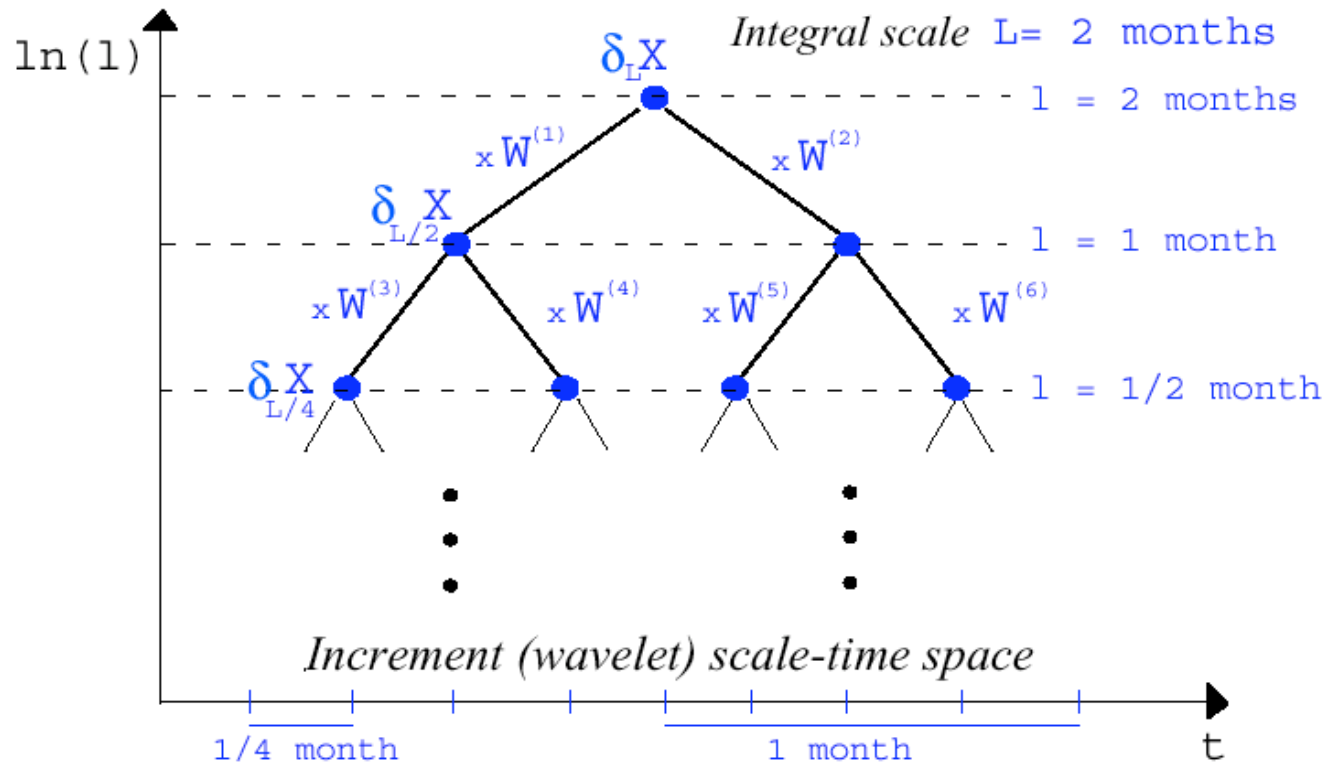




# The multiplicative cascade paradigm

$$\delta_{\lambda l} X(\lambda t) = \lambda^H \delta_l X(t) = W_\lambda \delta_l X(t)$$

- $\mathcal{W}$ -cascades (wavelet cascade)



## The Multifractal Random Walk (MRW) model

$$r_{\Delta t}(t) = \epsilon(t) \cdot \sigma_{\Delta t}(t) = \epsilon(t) \cdot e^{\omega_{\Delta t}(t)}$$

$$\mu_{\Delta t} = \frac{1}{2} \ln(\sigma^2 \Delta t) - C_{\Delta t}(0)$$

$$C_{\Delta t}(\tau) = \text{Cov}[\omega_{\Delta t}(t), \omega_{\Delta t}(t + \tau)] = \lambda^2 \ln \left( \frac{T}{|\tau| + e^{-3/2} \Delta t} \right)$$

$$\omega_{\Delta t}(t) = \mu_{\Delta t} + \int_{-\infty}^t d\tau \eta(\tau) K_{\Delta t}(t - \tau)$$

$\omega_{\Delta t}(t)$  is Gaussian with mean  $\mu_{\Delta t}$  and variance  $V_{\Delta t} = \int_0^{\infty} d\tau K_{\Delta t}^2(\tau) = \lambda^2 \ln \left( \frac{T e^{3/2}}{\Delta t} \right)$

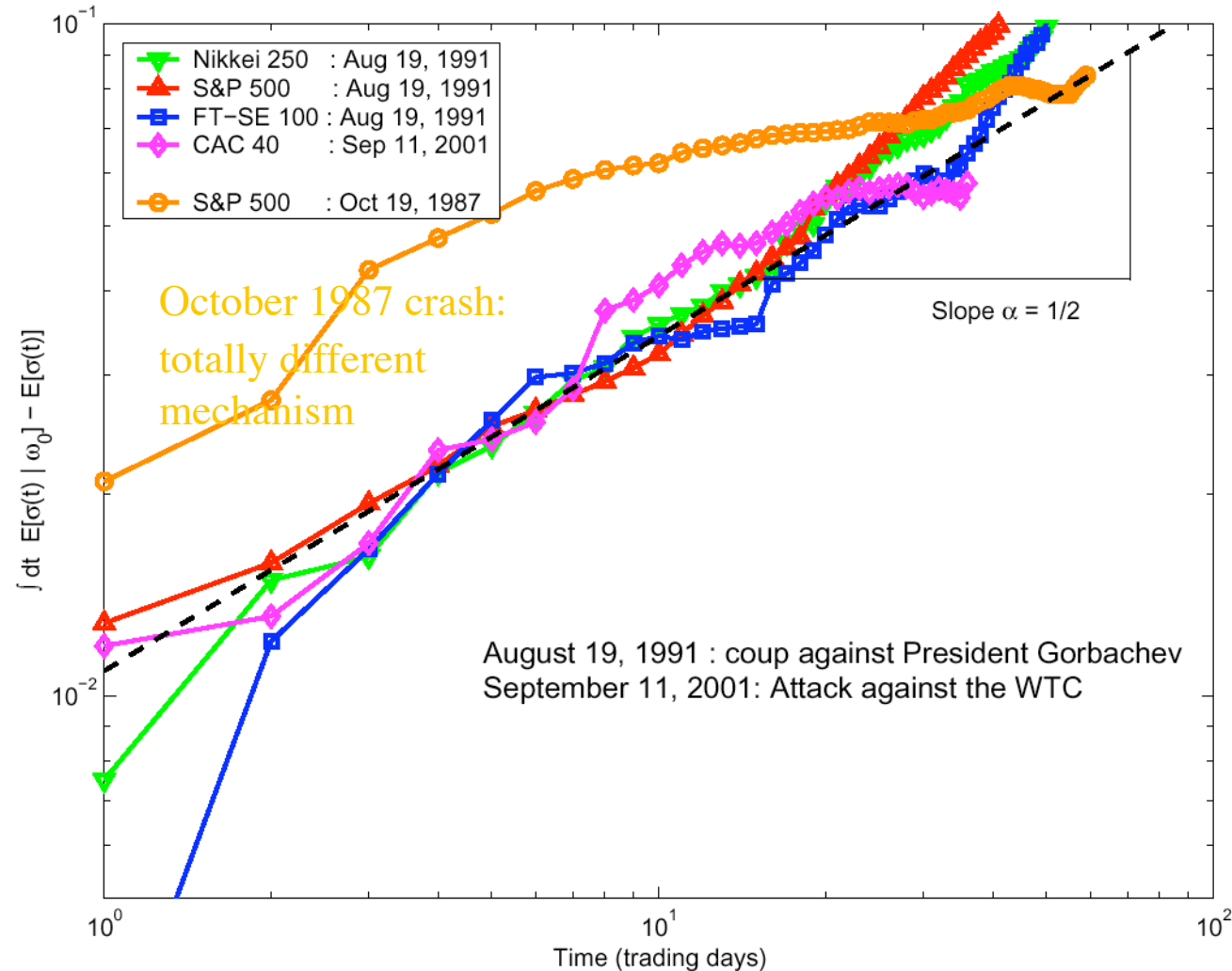
$$C_{\Delta t}(\tau) = \int_0^{\infty} dt K_{\Delta t}(t) K_{\Delta t}(t + |\tau|)$$

$$\hat{K}_{\Delta t}(f)^2 = \hat{C}_{\Delta t}(f) = 2\lambda^2 f^{-1} \left[ \int_0^{Tf} \frac{\sin(t)}{t} dt + O(f\Delta t \ln(f\Delta t)) \right]$$

$$K_{\Delta t}(\tau) \sim K_0 \sqrt{\frac{\lambda^2 T}{\tau}} \quad \text{for } \Delta t \ll \tau \ll T$$

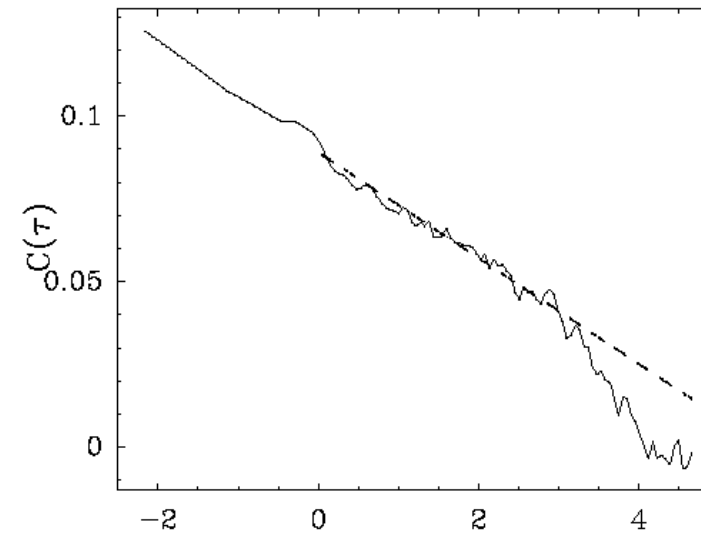
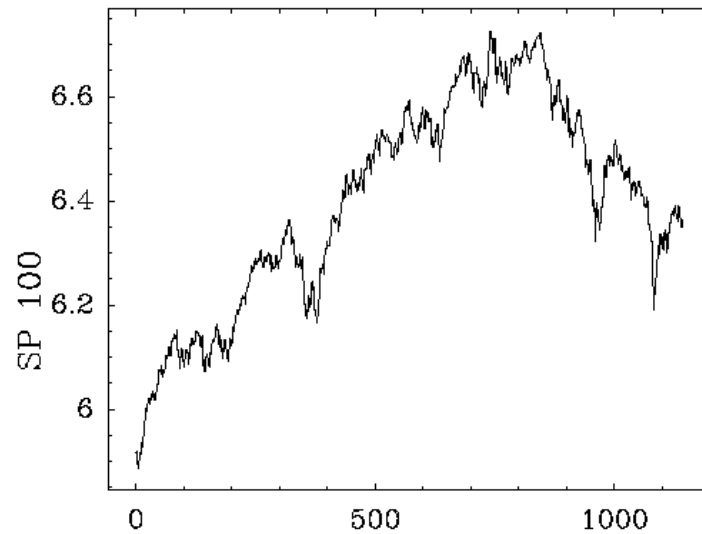
# Linear response to an external shock (Multifractal Random Walk model)

$$E_{\text{exo}}[\sigma^2(t) | \omega_0] - \overline{\sigma^2(t)} \propto e^{2K_0 t^{-1/2}} - 1 \approx \frac{2K_0}{\sqrt{t}}$$

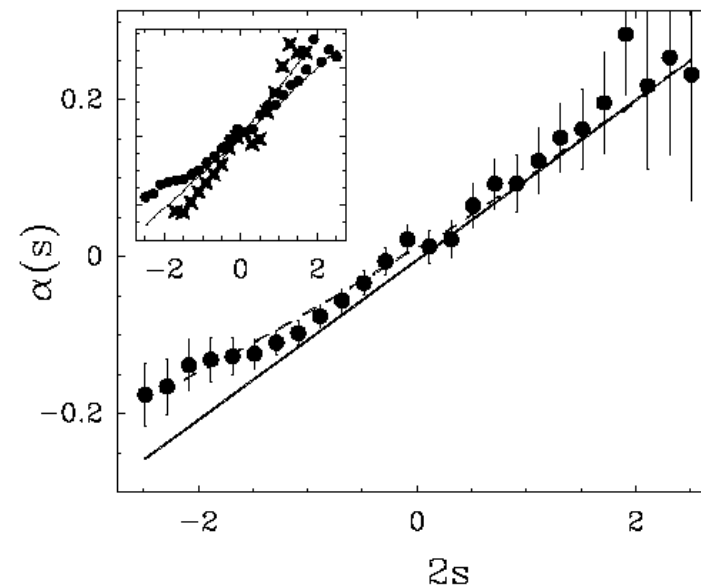
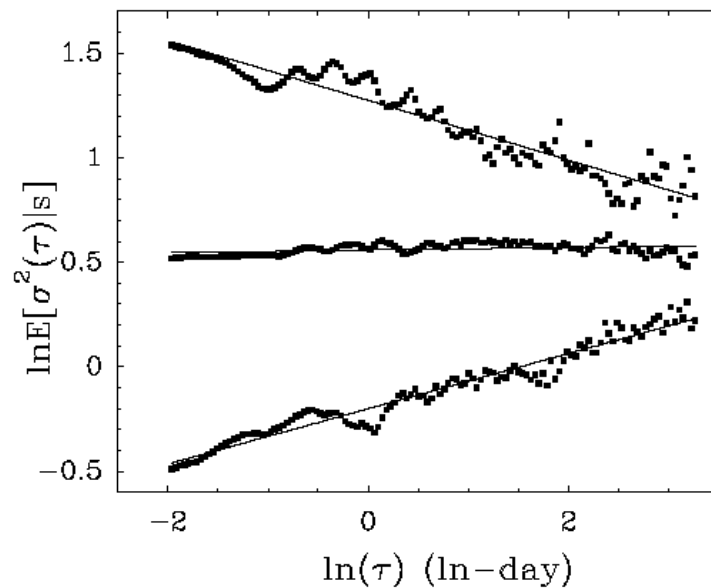


D. Sornette, Y.  
Malevergne and J.F.  
Muzy. Risk 16 (2),  
67-71 (2003)

# Endogenous shocks and Multifractal Random Walk model

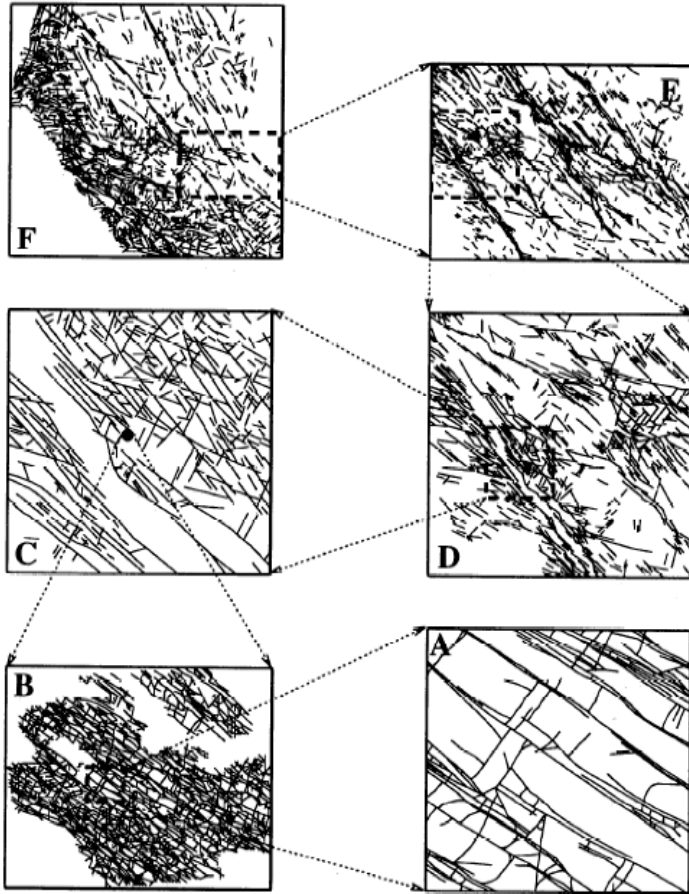


$$E_{\text{endo}}[\sigma^2(t) | \omega_0] \sim t^{-\alpha(s)} \quad \ln(\tau) \text{ (ln-day)}$$



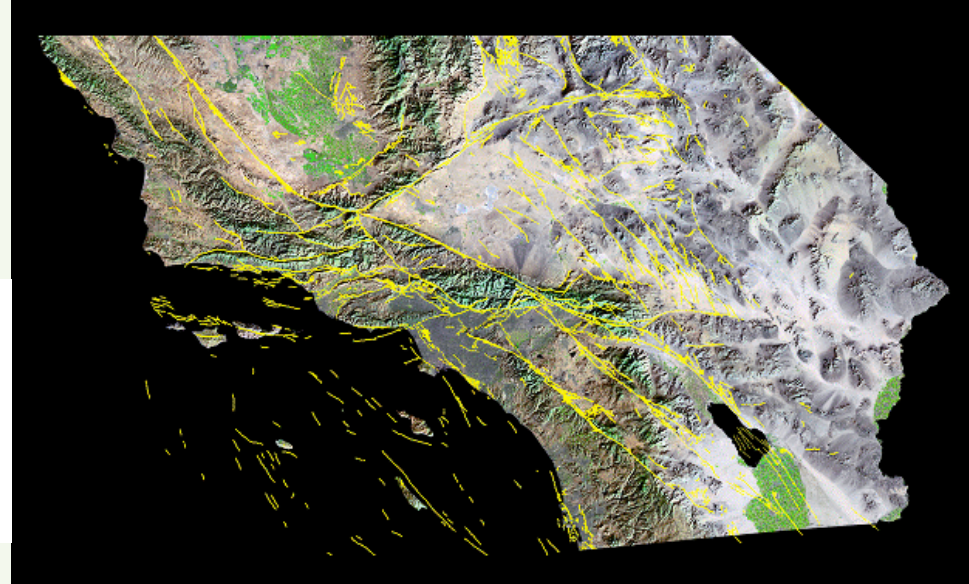
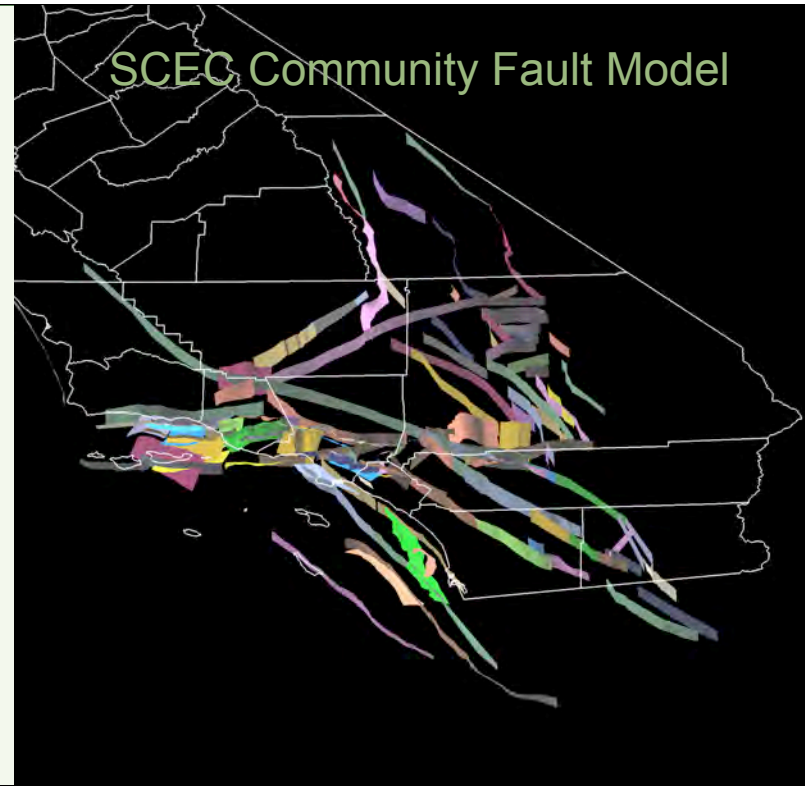
# Hierarchical geometry of faulting

Ouillon, Castaing, Sornette (JGR 1996)



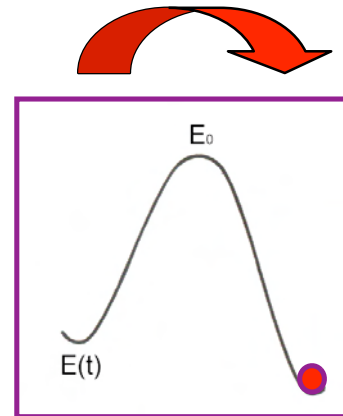
Map A: linear size=10 m, orig. scale=1:1  
Map B: linear size=60 m, orig. scale=1:220  
Map C: linear size=11 km, orig. scale=1:62,500  
Map D: linear size=45 km, orig. scale=1:125,000  
Map E: linear size=150 km, orig. scale=1:250,000  
Map F: linear size=400 km, orig. scale=1:1,000,000

## SCEC Community Fault Model



# The physical model : thermal activation driven by stress

Before the shock



After the shock

Energy barrier =  $E_0 - E(t)$

Arrhenius law for the activation rate:

$$\lambda(t) = \lambda_0 \exp\left(-\frac{E_0 - E(t)}{kT}\right)$$

stress barrier =  $\sigma_0 - \sigma(t)$

$$\lambda(t) = \lambda_0 \exp\left(-\frac{\sigma_0 - \sigma(t)}{kT} V\right)$$

**Compatible with state-and-rate friction, stress corrosion, ...**

$\lambda(t)$  : instantaneous rate

$\lambda_0 \sim$  average nucleation rate

$\sigma_0$  : material strength

$\sigma(t)$  : applied stress

$V$  : activation volume

$T$  : temperature

$k$  : Boltzmann constant



Experiments by Zhurkov Int. J. Fract. Mech. 1, 311 (1965)

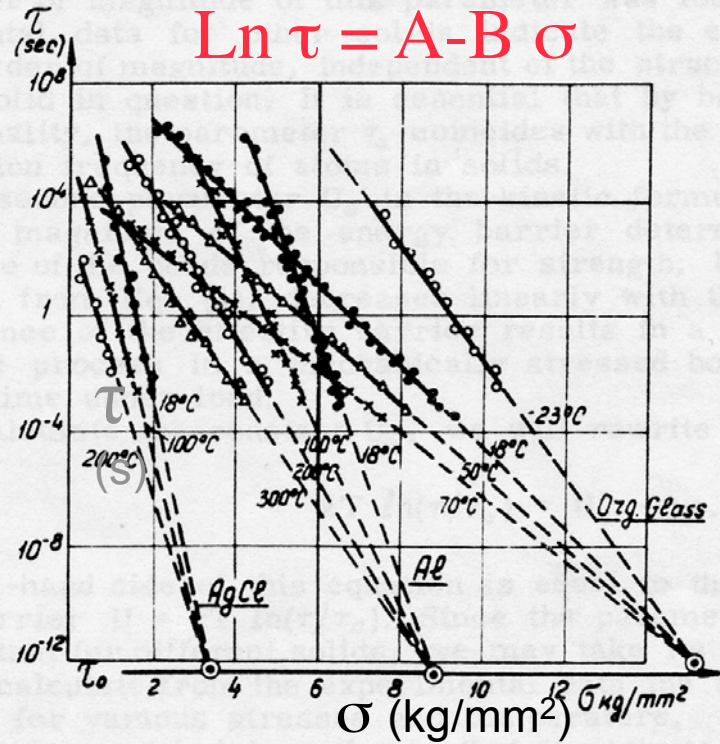


Fig. 5. Time and temperature dependence of the lifetime of solids on stress.  
 1. Silver chloride (Reference 4)  
 2. Aluminum (Reference 5)  
 3. Plexiglas (Reference 6)

$$\tau = \tau_0 \exp\left(\frac{U}{k_B T}\right)$$

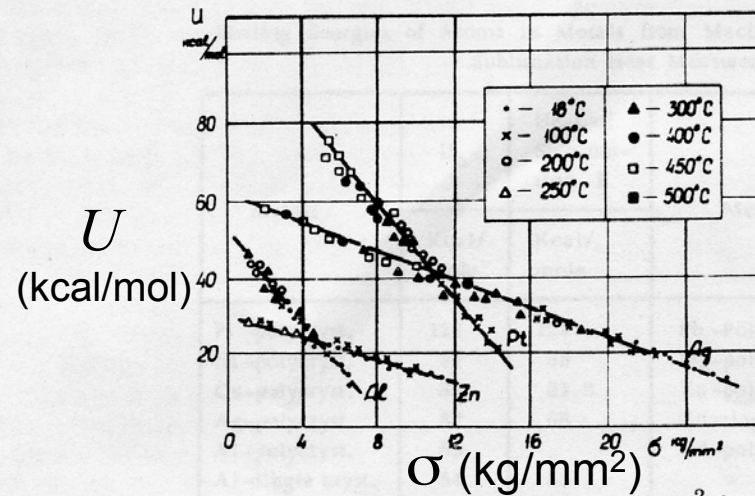


Fig. 6. Effective barrier  $U$  kcal/mol vs. tensile stress  $\sigma$  kg/mm<sup>2</sup> for polycrystalline

**Empirical energy barrier**

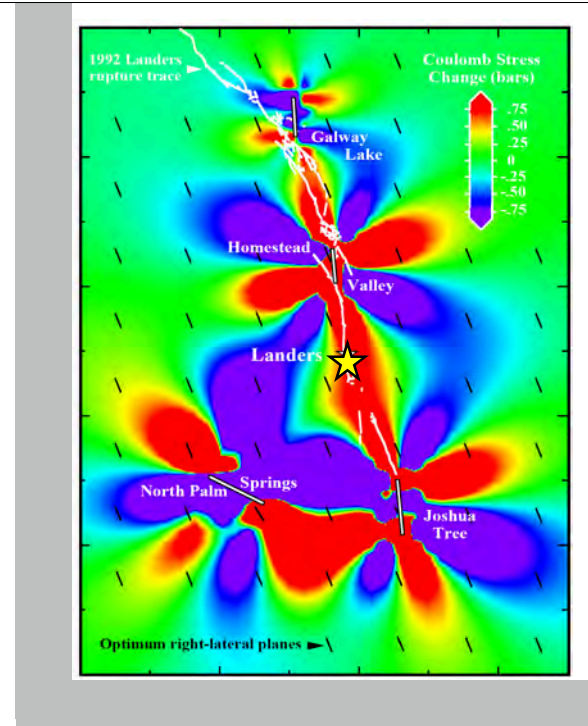
$$U = U_0 - \alpha \sigma$$

où  $U_0$ : énergie de sublimation

A possible mechanism : thermal activated process

# Taking account of history and boundary conditions

$$\lambda(\vec{r}, t) = \lambda'_0 \exp\left(\frac{\sigma(\vec{r}, t)}{kT} V\right)$$



Stress is assumed to be a scalar for the sake of simplicity

$$\Sigma(\vec{r}, t) = \Sigma_{\text{far field}}(\vec{r}, t) + \int_{-\infty}^t \int dN[d\vec{r}' \times d\tau] \Delta\sigma(\vec{r}', \tau) g(\vec{r} - \vec{r}', t - \tau)$$

local  
stress

tectonic  
loading

**Stress fluctuations induced by all past events  
in the system**

D. Sornette and G. Ouillon, Multifractal Scaling of Thermally-Activated Rupture Processes, Phys. Rev. Lett. 94, 038501 (2005)  
G. Ouillon and D. Sornette, Magnitude-Dependent Omori Law: Theory and Empirical Study, J. Geophys. Res., 110, B04306, doi:10.1029/2004JB003311 (2005).

$$\Sigma(\vec{r}, t) = \Sigma_{\text{far field}}(\vec{r}, t) + \int_{-\infty}^t \int dN[d\vec{r}' \times d\tau] \Delta\sigma(\vec{r}', \tau) g(\vec{r} - \vec{r}', t - \tau)$$

$$g(\vec{r}, t) = f(\vec{r}) \times h(t)$$

The **rheology is viscoplastic**, with a relaxation function featuring a very large relaxation time  $\tau_M$  :

$$h(t) = \frac{h_0}{(t + t_1)^{1+\theta}} \exp\left(-\frac{t}{\tau_M}\right)$$

At each location, **stress fluctuations** due to previous events are distributed as:

$$P(\sigma) d\sigma \approx \frac{C}{(\sigma + \sigma_0)^{1+\mu}} d\sigma$$

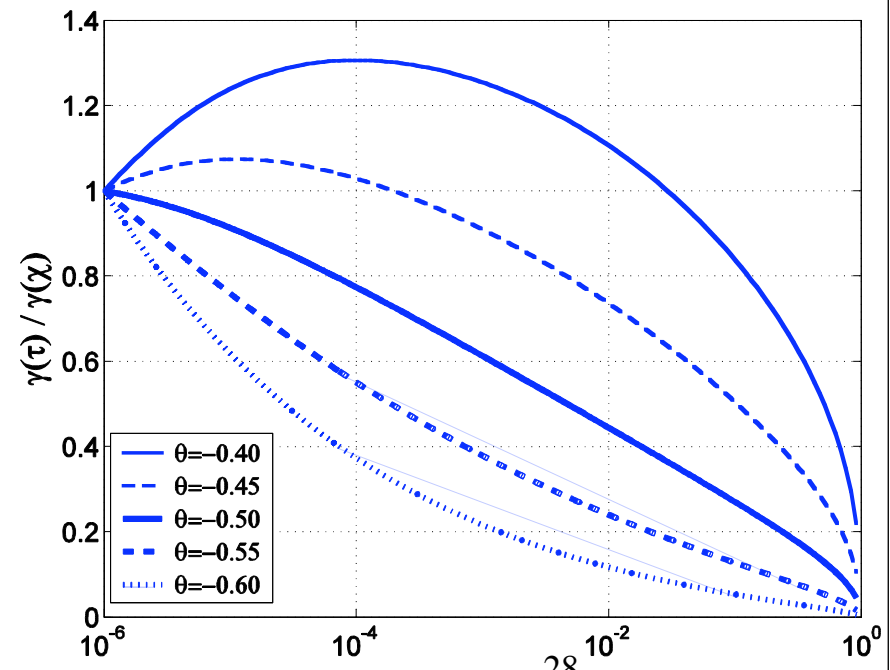
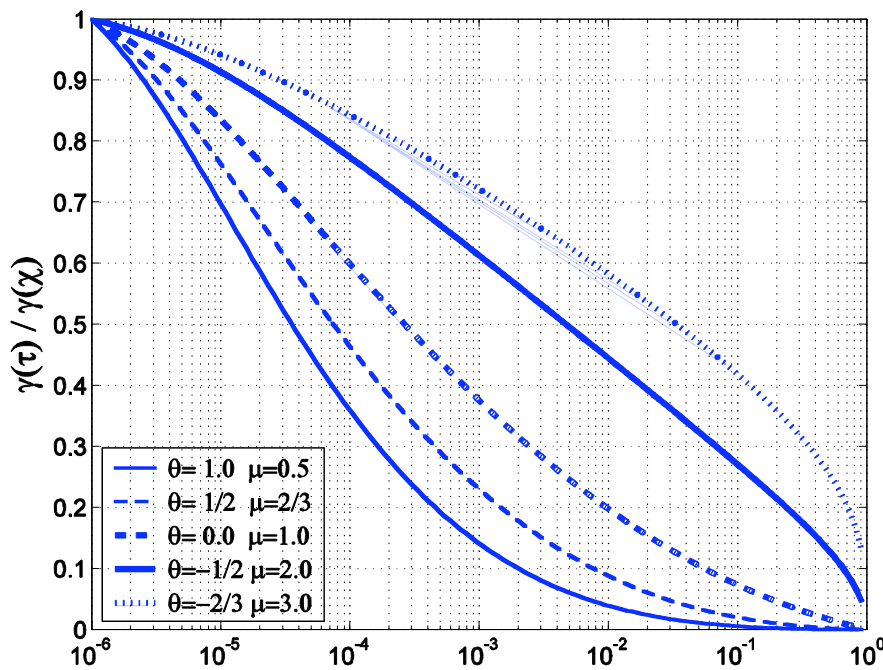
Theoretical predictions using tail covariance (Ide-Sornette, 2001)

$$\Pr[\lambda(t) > \lambda_q | \lambda_M] = \Pr[e^{\beta\omega(t)} > \frac{\lambda_q}{\lambda_{\text{tec}}} | \omega_M] = \Pr[\omega(t) > (1/\beta) \ln \left( \frac{\lambda_q}{\lambda_{\text{tec}}} \right) | \omega_M]$$

$$\lambda_q(t) = A_q \lambda_{\text{tec}} e^{\beta\gamma(t)\omega_M} \quad \gamma(t) = \frac{h_0^2}{\Delta t^{2/\mu}} \left( \frac{1}{t^{2m-1}} \int_{c/t}^{T+c} dy \frac{1}{(y+1)^m} \frac{1}{y^m} \right)^{\frac{2}{\mu}}$$

$$m = (1 + \theta)\mu/2.$$

Since  $\gamma(t) \sim \ln(t)$  and  $\omega_m \sim M$ , we obtain  $p(M) = a M + b$



We obtain an exact multifractality if  $\mu(1+\theta) \sim 1$

$$p(M) = aM + b$$

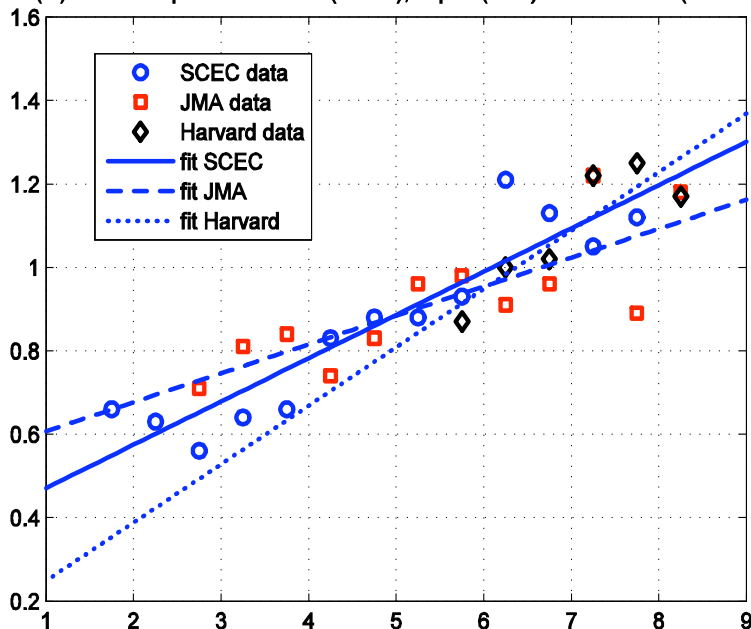
We processed three catalogs, that we pre-processed to check for their completeness and its evolution with time.

We then computed stacked aftershocks time series, sorting them within intervals of 0.5 magnitude amplitudes.

**We clearly observed a linear dependence of  $p$  with magnitude  $M$ .**

Statistical tests have been performed using a bootstrap strategy, and we were able to show that all slopes were significantly different from 0, and that all linear relationships were significantly different from each other.

P(M) relationships for California (SCEC), Japan (JMA) and the world (Harvard)



*Ribeiro et al, 2006*

**For Southern California (SCEC catalog):**

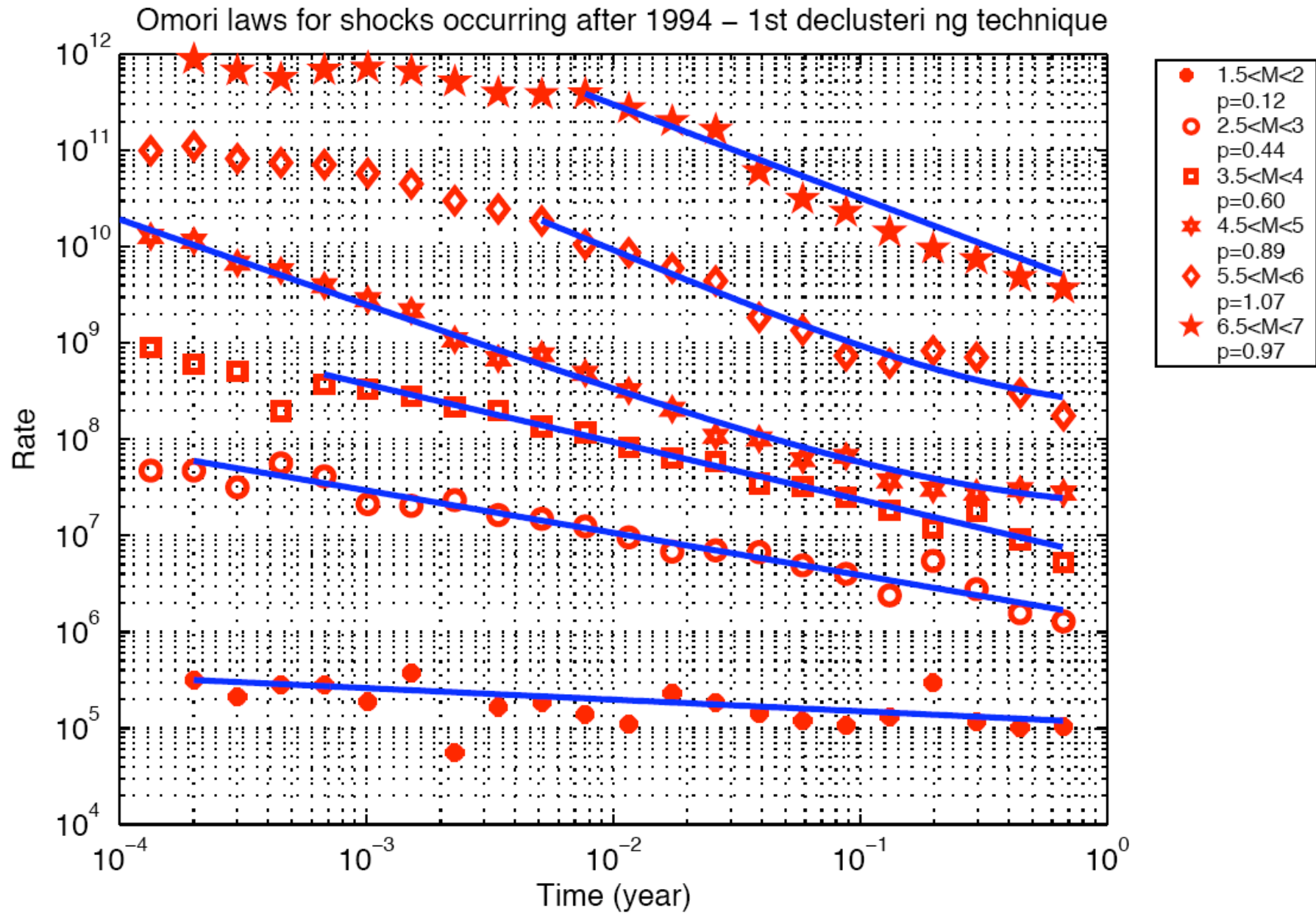
$$p(M) = 0.10M + 0.37$$

**For Japan (JMA catalog):**

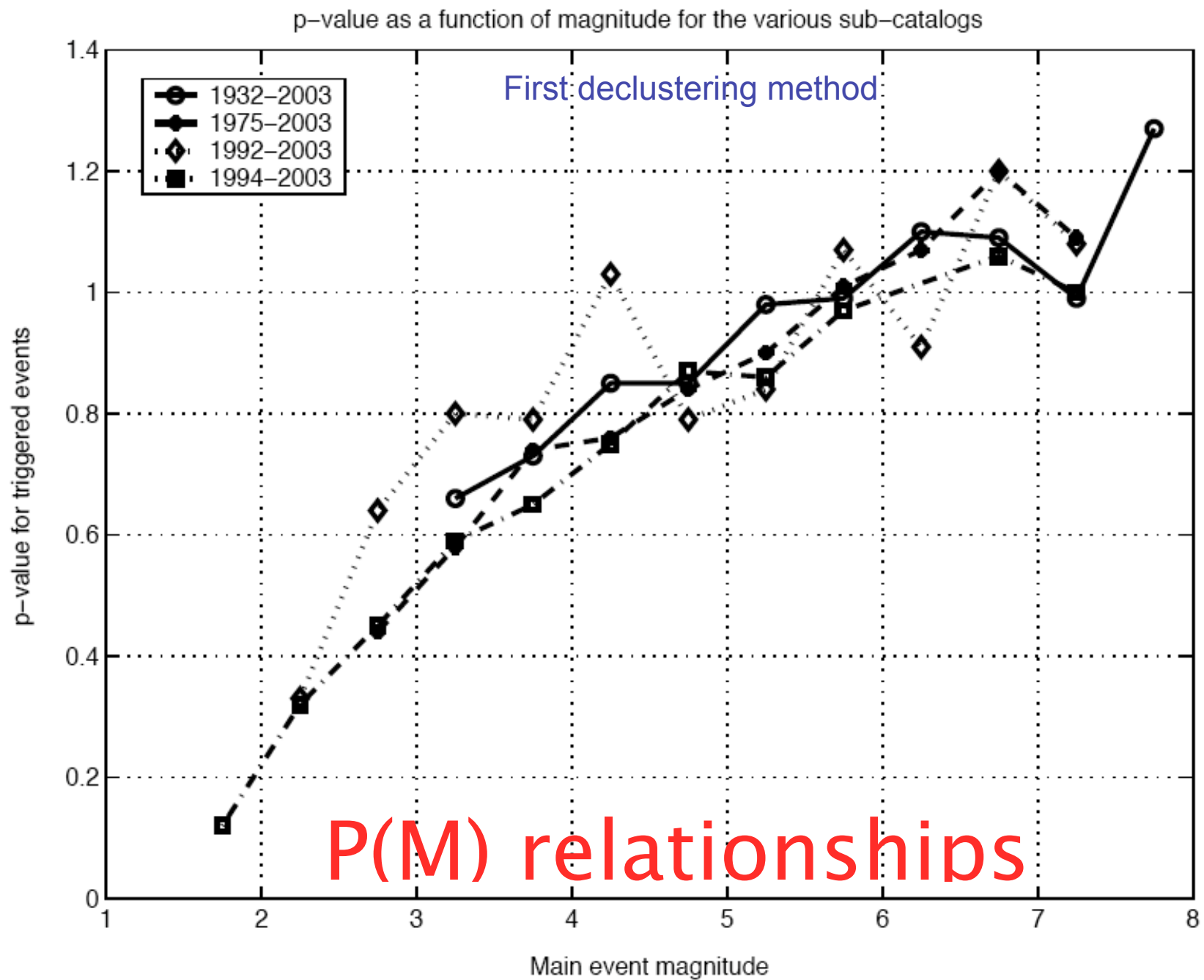
$$p(M) = 0.07M + 0.54$$

**For the World (Harvard catalog):**

$$p(M) = 0.14M + 0.11$$



D. Sornette and G. Ouillon, Multifractal Scaling of Thermally-Activated Rupture Processes, Phys. Rev. Lett. 94, 038501 (2005)



# Epileptic Seizures – Quakes of the Brain?

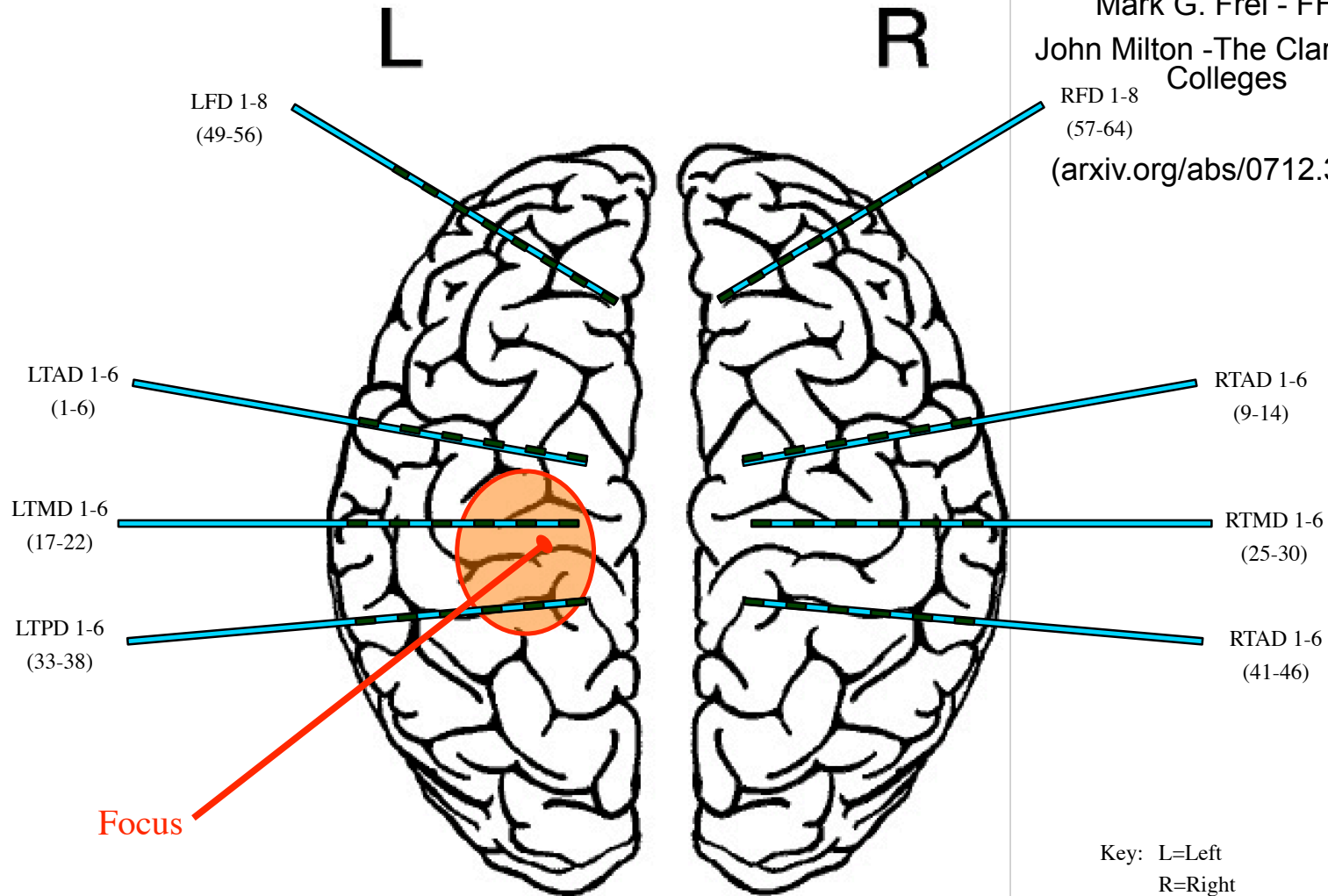
with Ivan Osorio – KUMC & FHS

Mark G. Frei - FHS

John Milton -The Claremont  
Colleges

RFD 1-8  
(57-64)

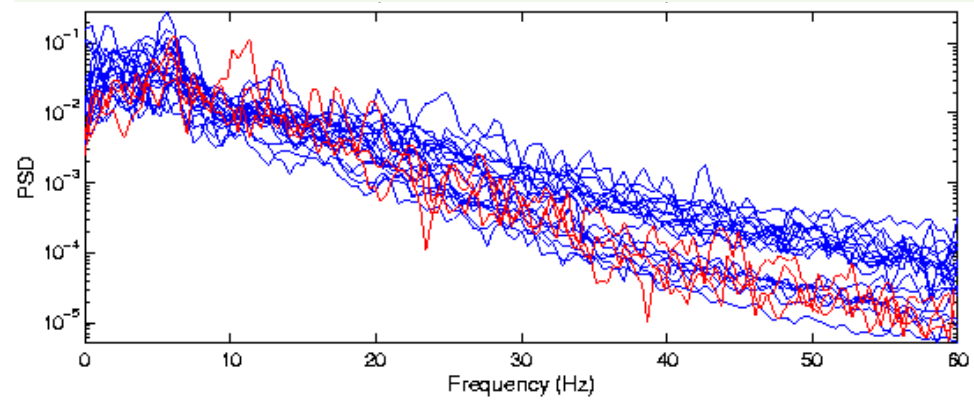
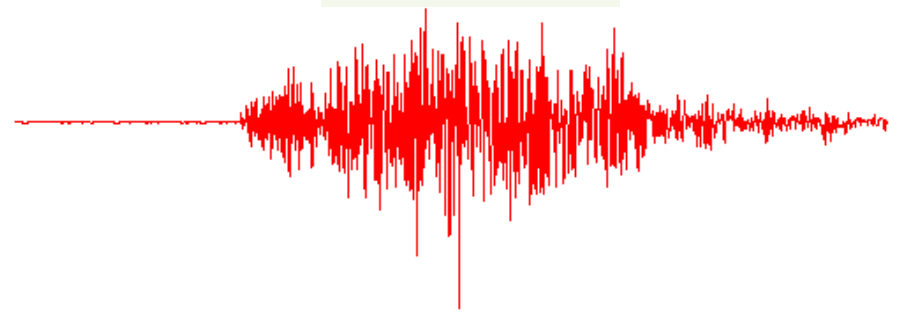
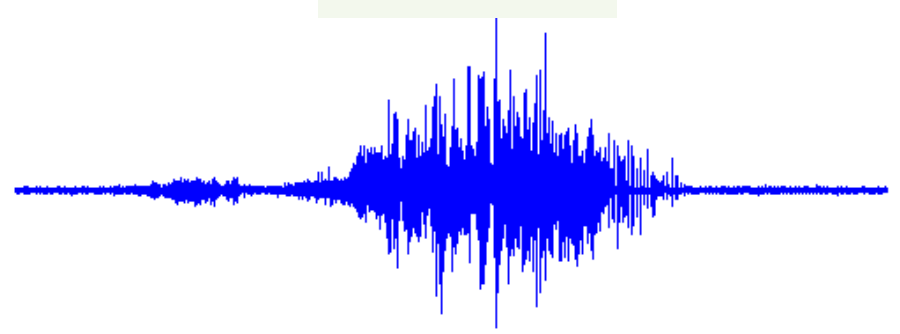
([arxiv.org/abs/0712.3929](https://arxiv.org/abs/0712.3929))



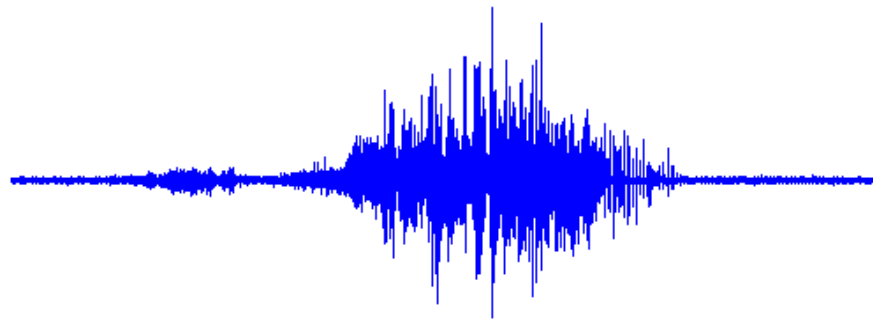
Depth Needle Electrodes Contact Numbering: N ... 3 2 1

Key: L=Left  
R=Right  
A=Anterior  
M=Mesial  
P=Posterior  
D=Depth  
T=Temporal

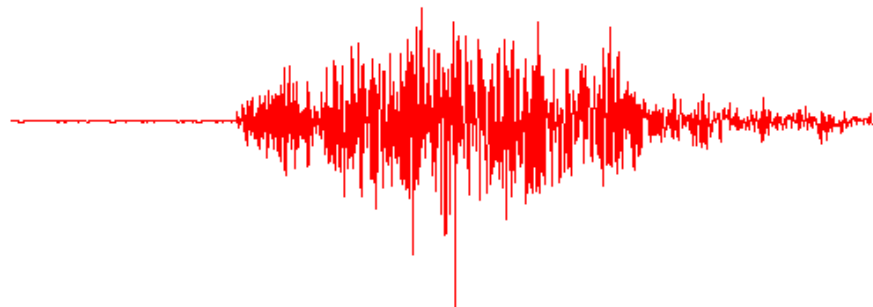




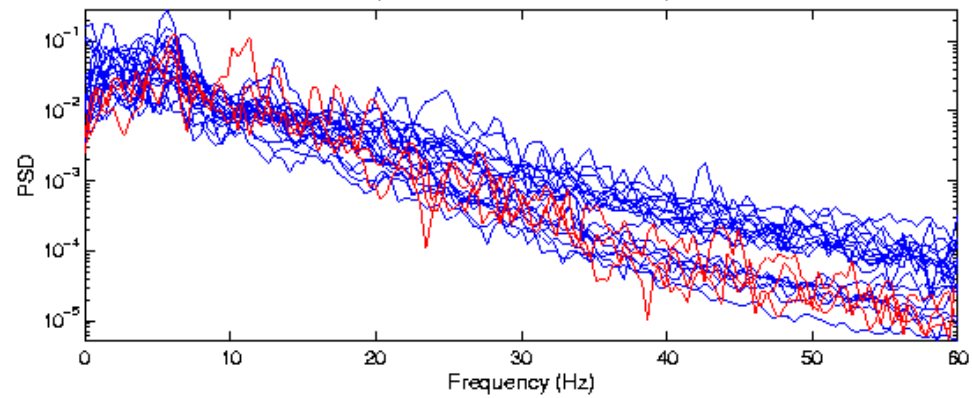
Seizure



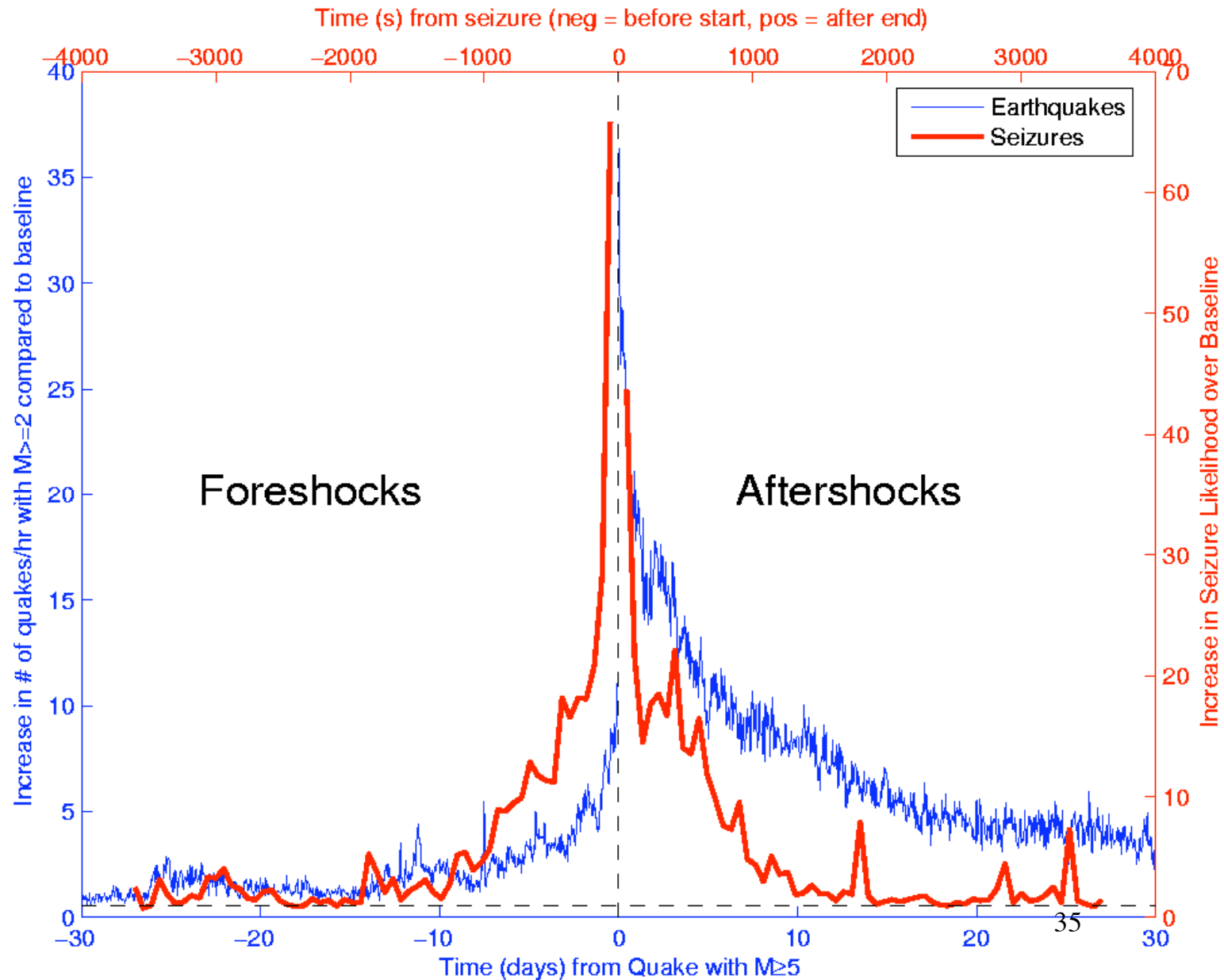
Earthquake

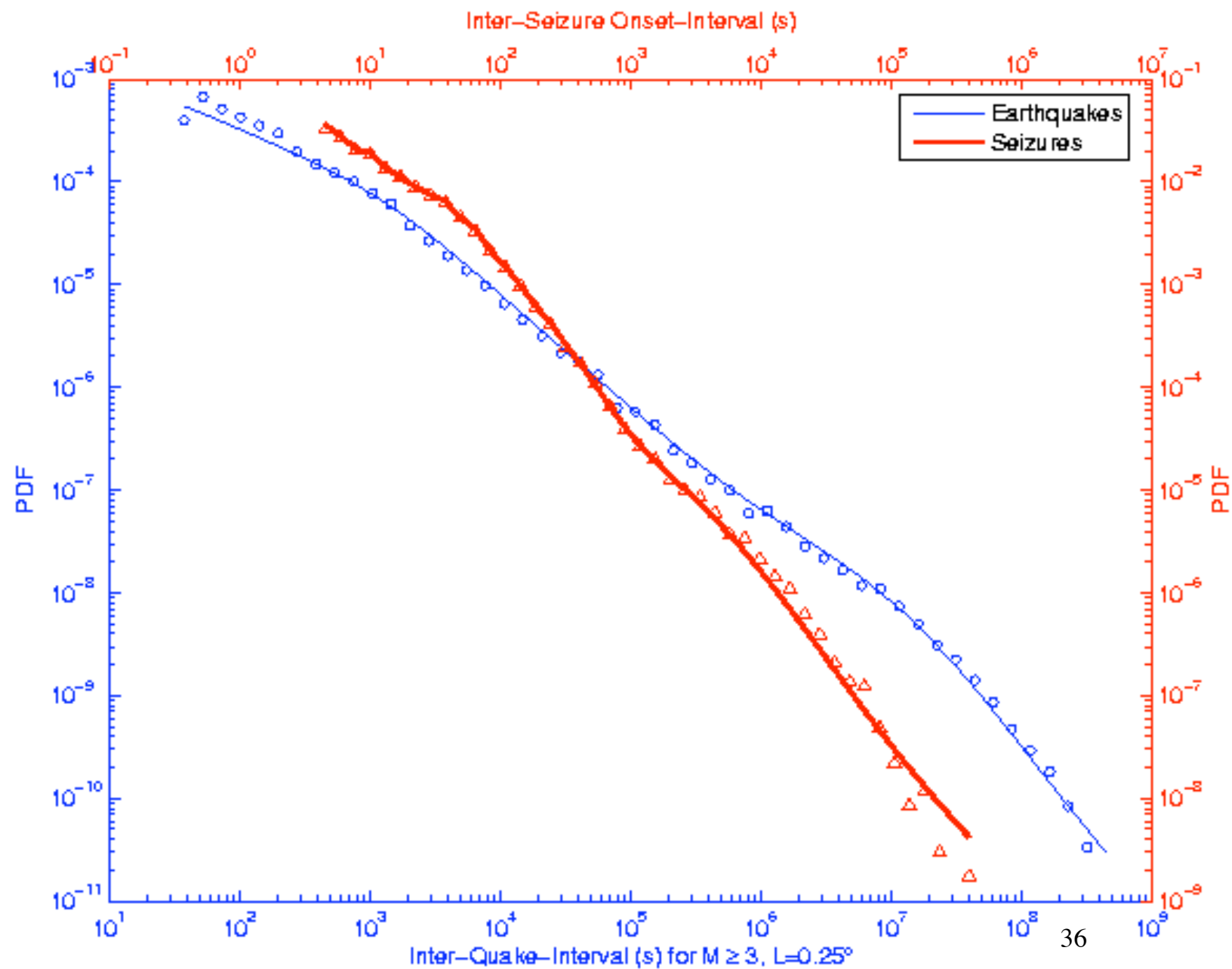


PSD estimates for 20 seizures (blue) and triaxial acceleration components for Loma Prieta Quake (red)

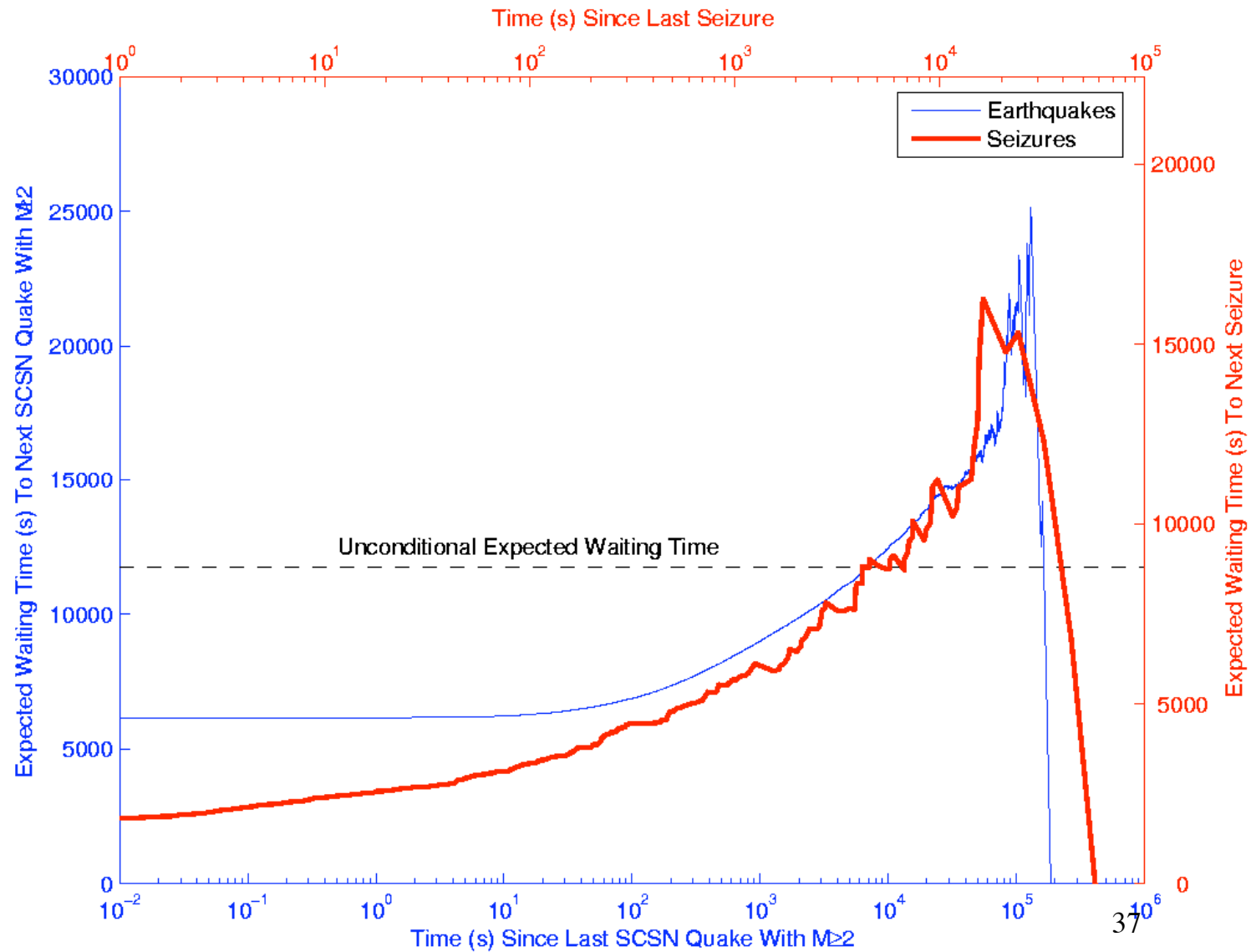


# Omori law: Direct and Inverse

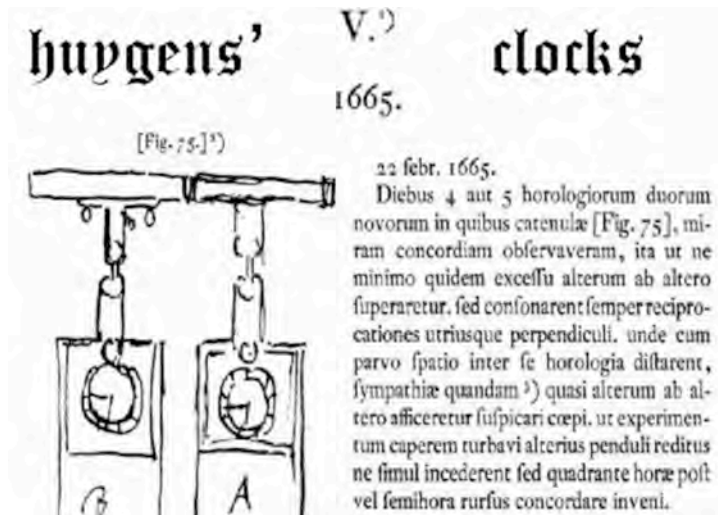




The longer it has been since the last event,  
the longer it will be since the next one! (Sornette&Knopoff, 1997)



# SYNCHRONISATION AND COLLECTIVE EFFECTS IN EXTENDED STOCHASTIC SYSTEMS



Fireflies

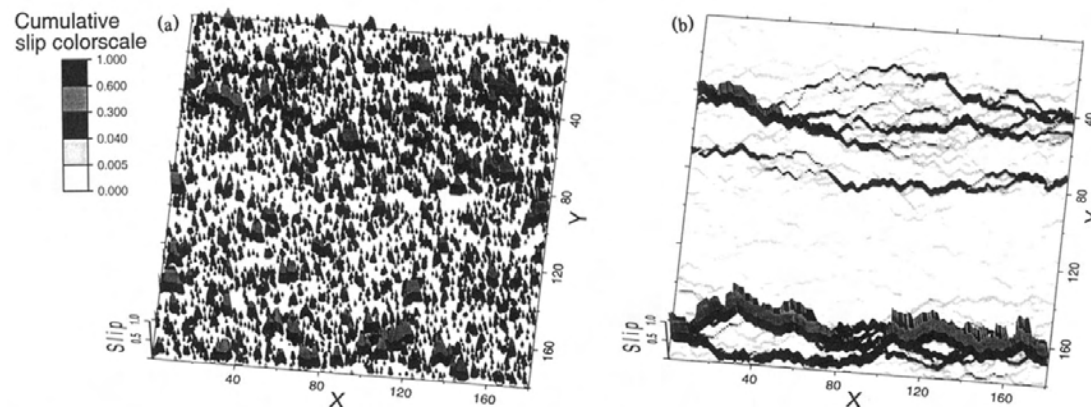
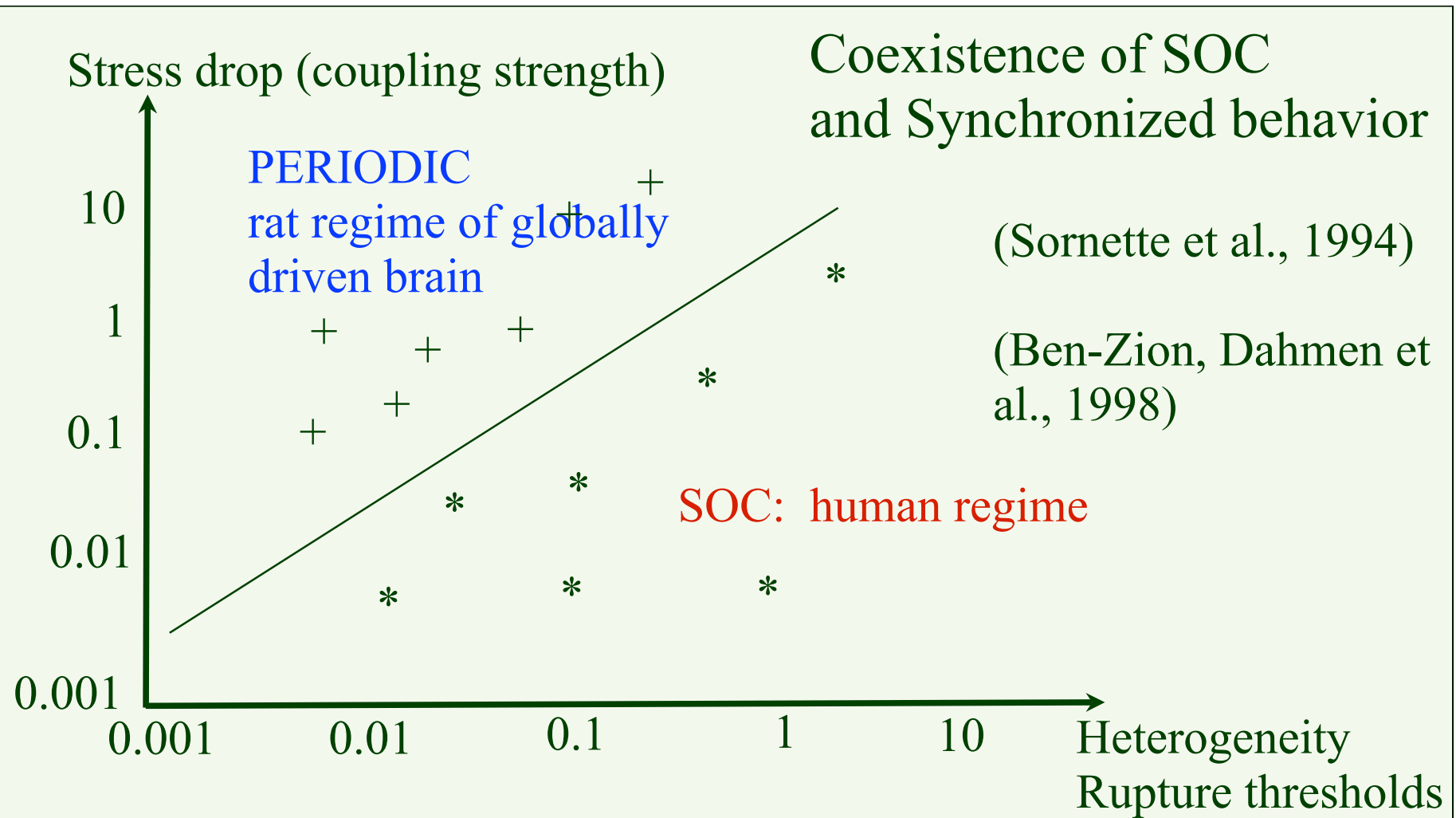


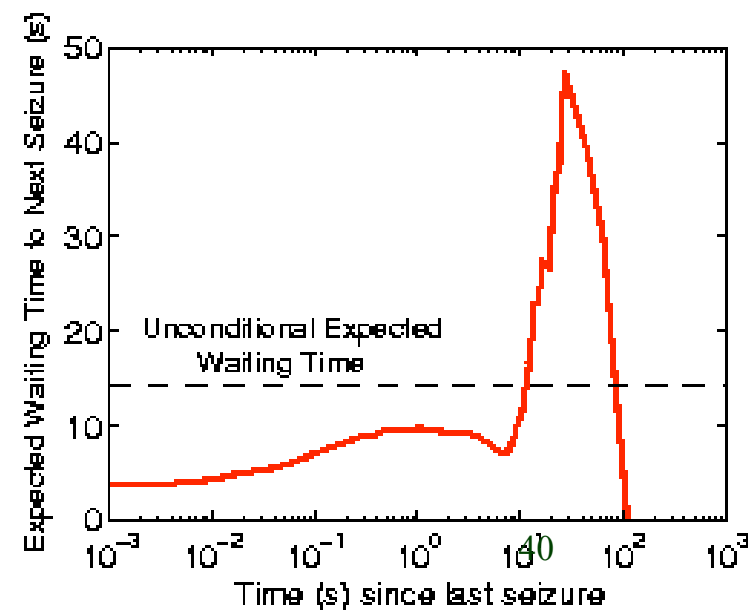
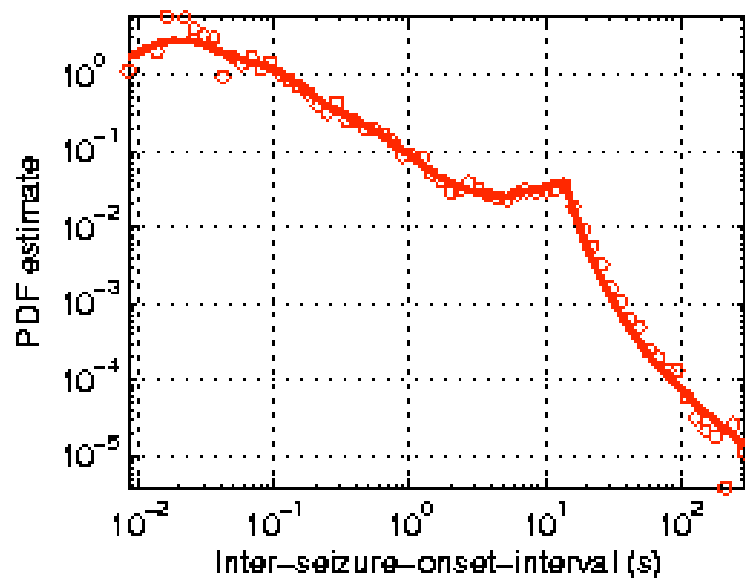
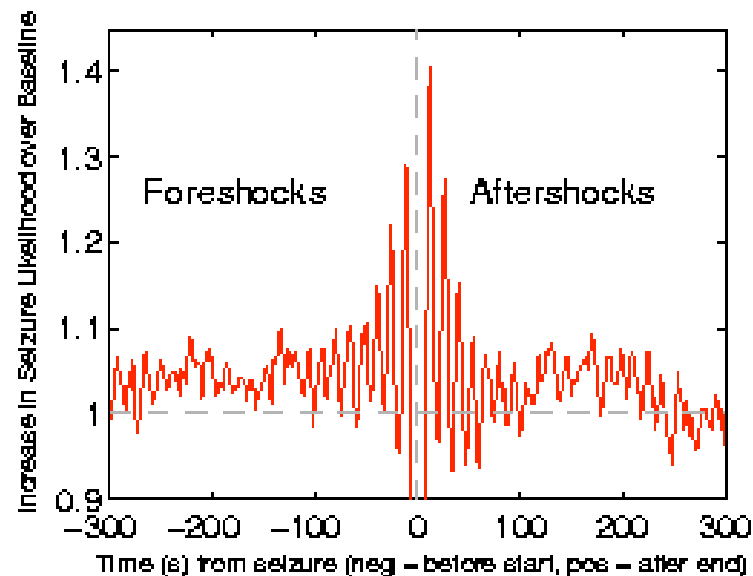
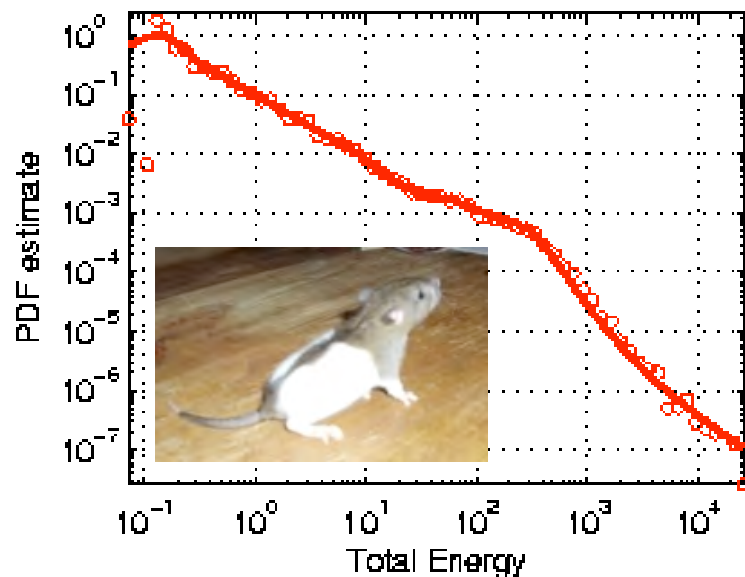
FIG. 1. Evolution of the cumulative earthquake slip, represented along the vertical axis in the white to black color code shown above the picture, at two different times: (a) early time and (b) long time, in a system of size  $L=90$  by  $L=90$ , where  $\Delta\sigma=1.9$  and  $\beta=0.1$ .

Miltenberger et al. (1993)

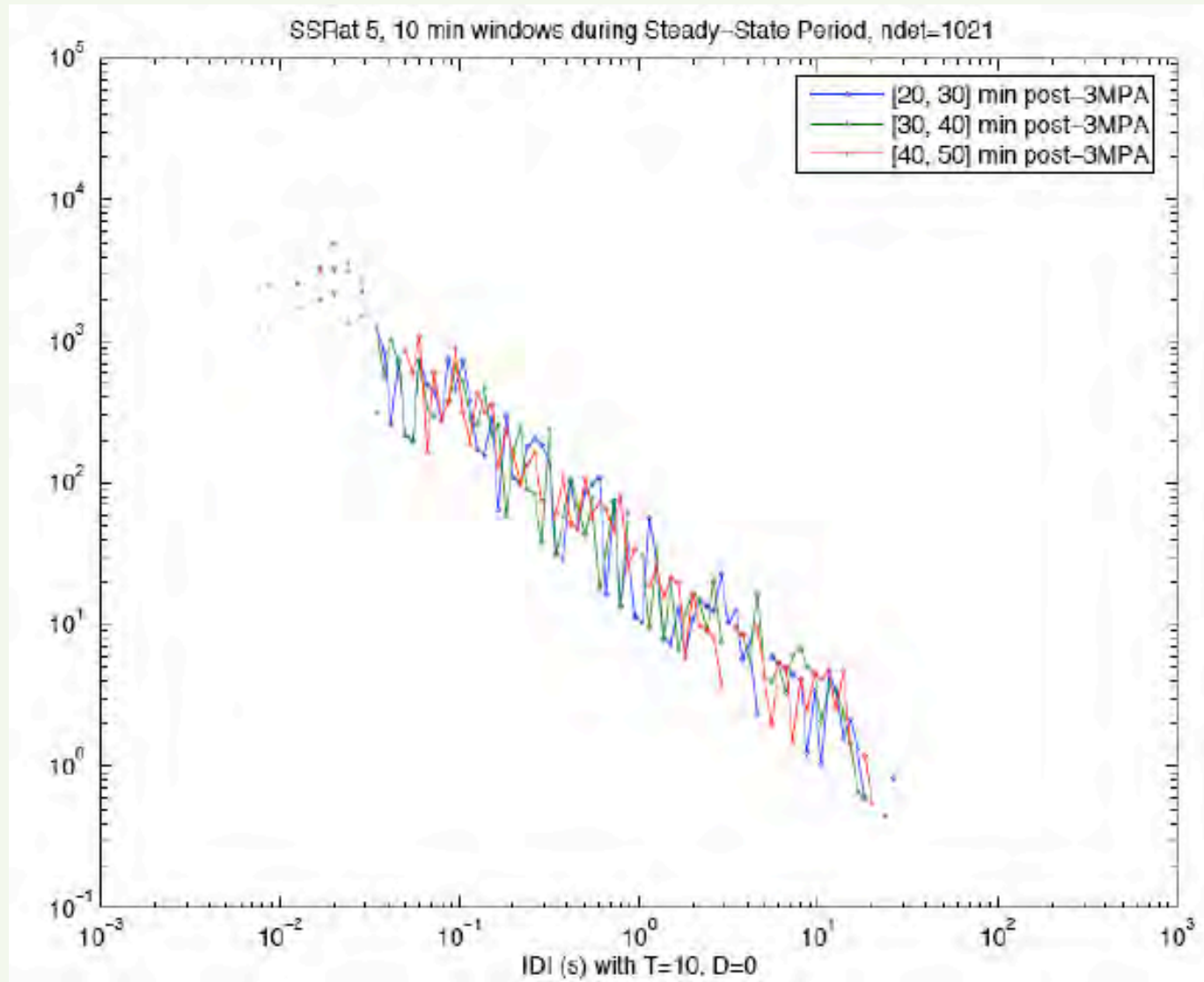


“Phase diagram” for the model in the space (heterogeneity, stress drop). Crosses (+) correspond to systems which exhibit a periodic time evolution. Stars \* corresponds to systems that are self-organized critical, with a Gutenberg-Richter earthquake size distribution and fault localization whose geometry is well-described by the geometry of random directed polymers.

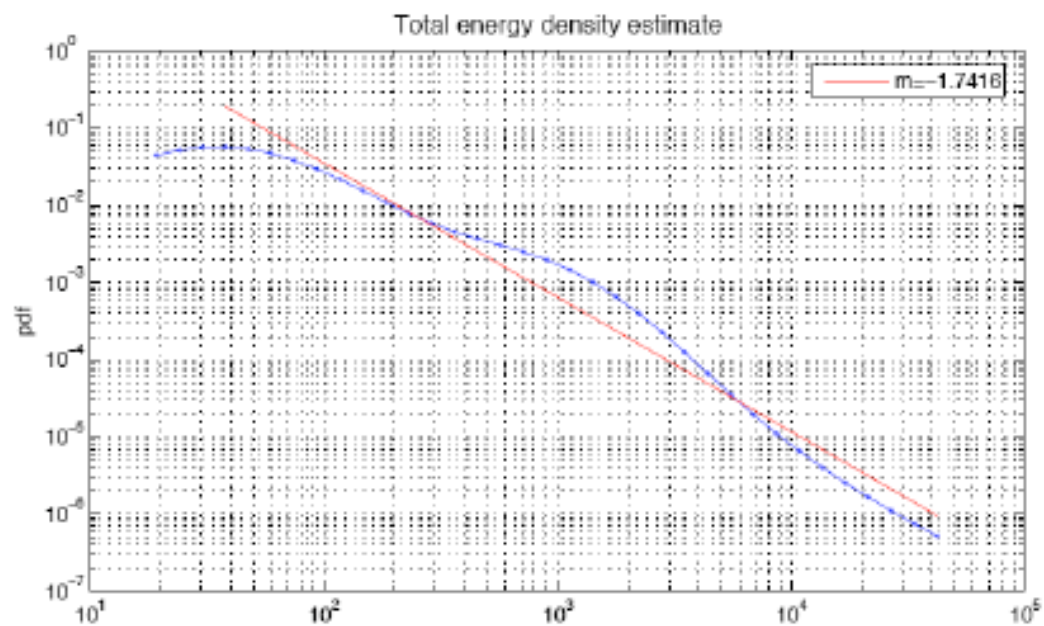
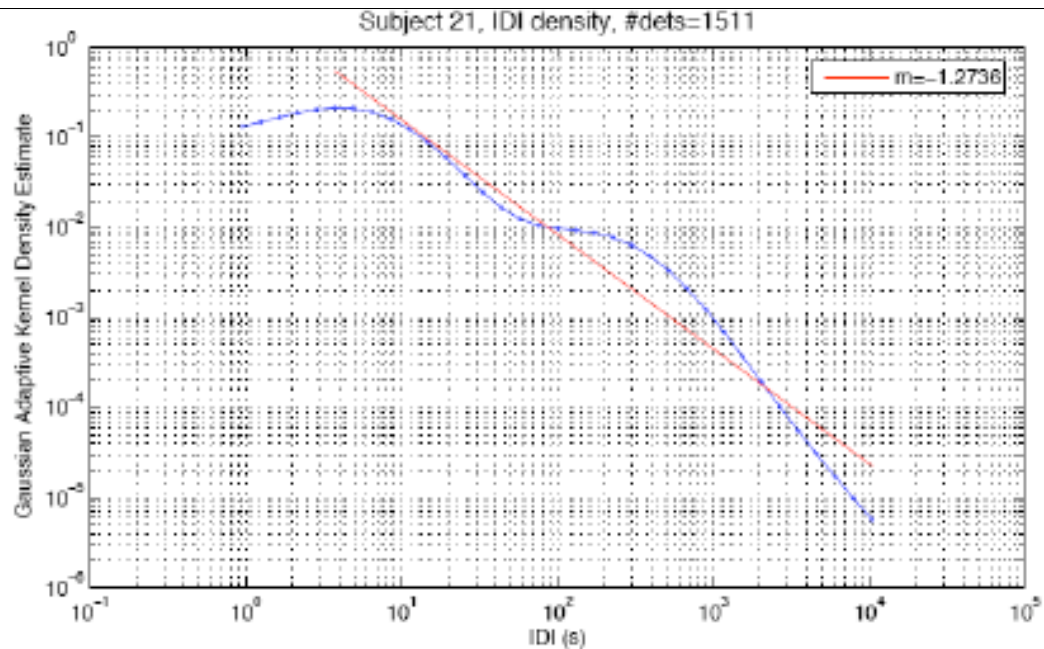
19 rats treated intravenously (2) with the convulsant 3-mercapto-proprionic acid (3-MPA)







Distribution of inter-seizure time intervals for rat 5, demonstrating a pure power law, which is characteristic of the SOC state. This scale-free distribution should be contrasted with the pdf's obtained for the other rats, which are marked by a strong shoulder associated with a characteristic time scale, which reveals the periodic regime.

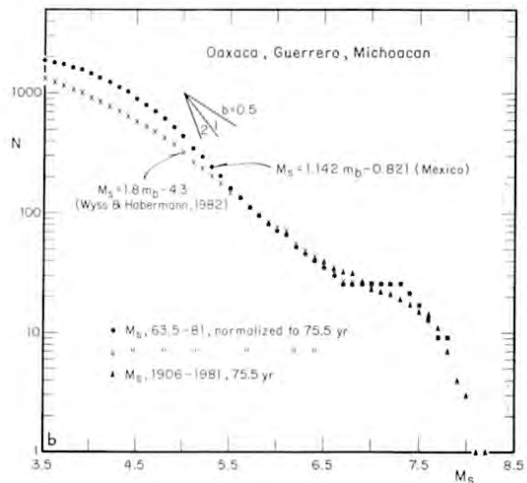


The pdf's of the seizure energies and of the inter-seizure waiting times for subject 21.

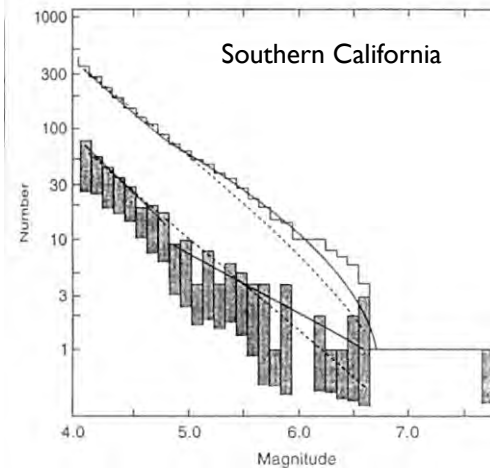
Note the shoulder in each distribution, demonstrating the presence of a characteristic size and time scale, qualifying the periodic regime.

# Complex magnitude distributions

## Characteristic earthquakes?

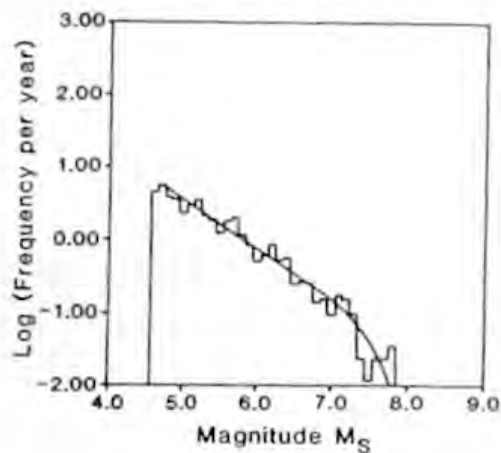


*Singh, et. al.,  
1983, BSSA 73,  
1779-1796*

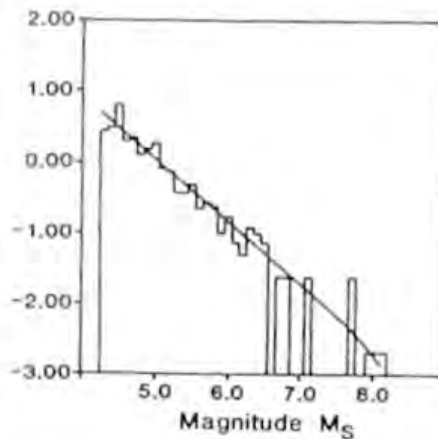


*Knopoff, 2000,  
PNAS 97,  
11880-11884*

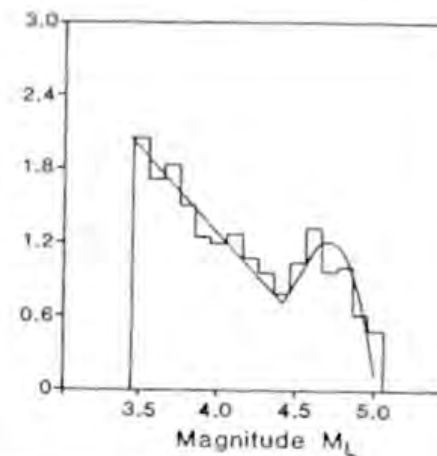
(a) Eastern Mediterranean



(b) Southern California

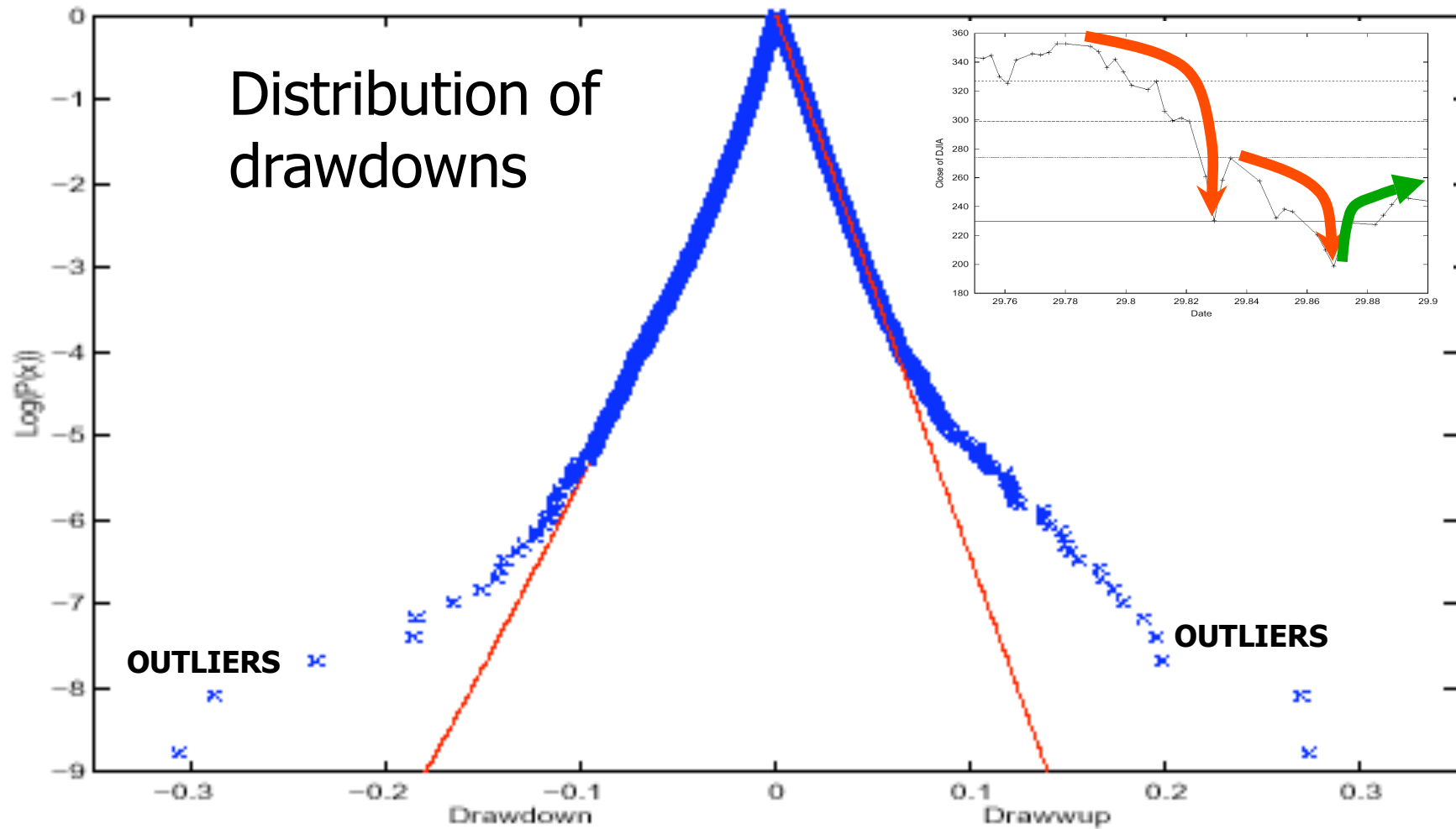


(c) Mount St. Helens



*Main, 1995,  
BSSA 85,  
1299-1308*

# Dow Jones Industrial Average



A. Johansen and D. Sornette, Stock market crashes are outliers,  
European Physical Journal B 1, 141-143 (1998)

A. Johansen and D. Sornette, Large Stock Market Price Drawdowns Are Outliers,  
Journal of Risk 4(2), 69-110, Winter 2001/02

# Securitization of credit risks: is it the next “systemic collapse”?

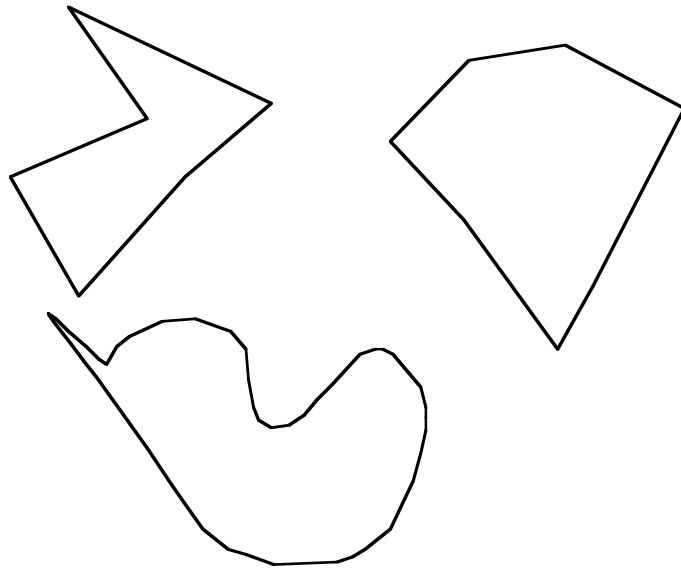
- Securitization of credit risks leads to smaller risks
- But more inter-connected  
⇒ global risk?

CDS and CDO: form of insurance contracts linked to underlying debt that protects the buyer in case of default.

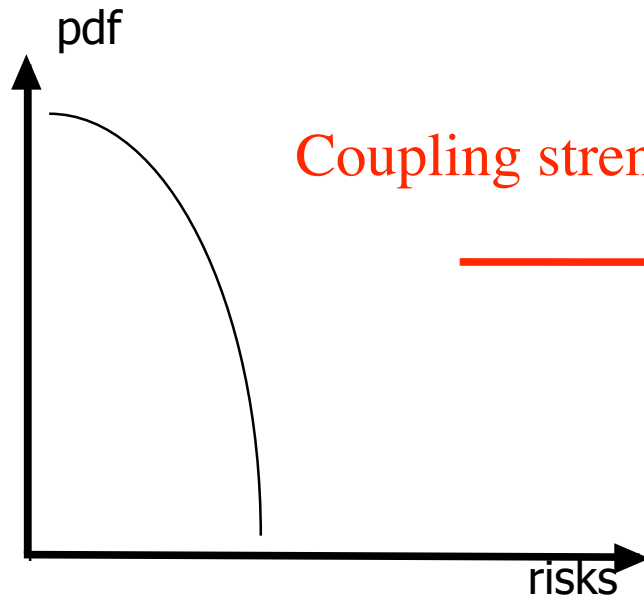
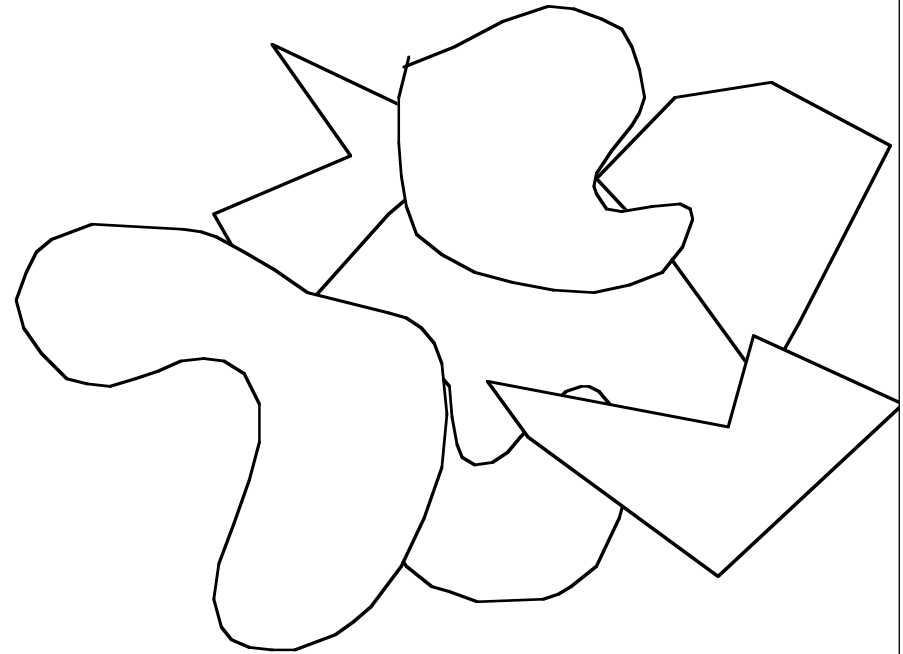
The market has almost doubled in size every year for the past five years, reaching \$20 trillion in notional amounts outstanding last June 2007, according to the Bank for International Settlements.

Bundling of indexes of CDSs together and slicing them into tranches, based on riskiness and return. The most toxic tranche at the bottom exposes the holder to the first 3% of losses but also gives him a large portion of the returns. At the top, the risks and returns are much smaller-unless there is a systemic failure.

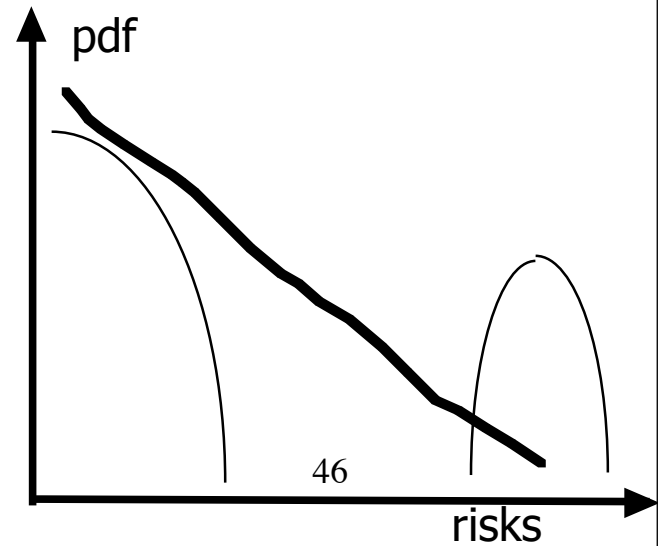
Separation of financial and credit risks

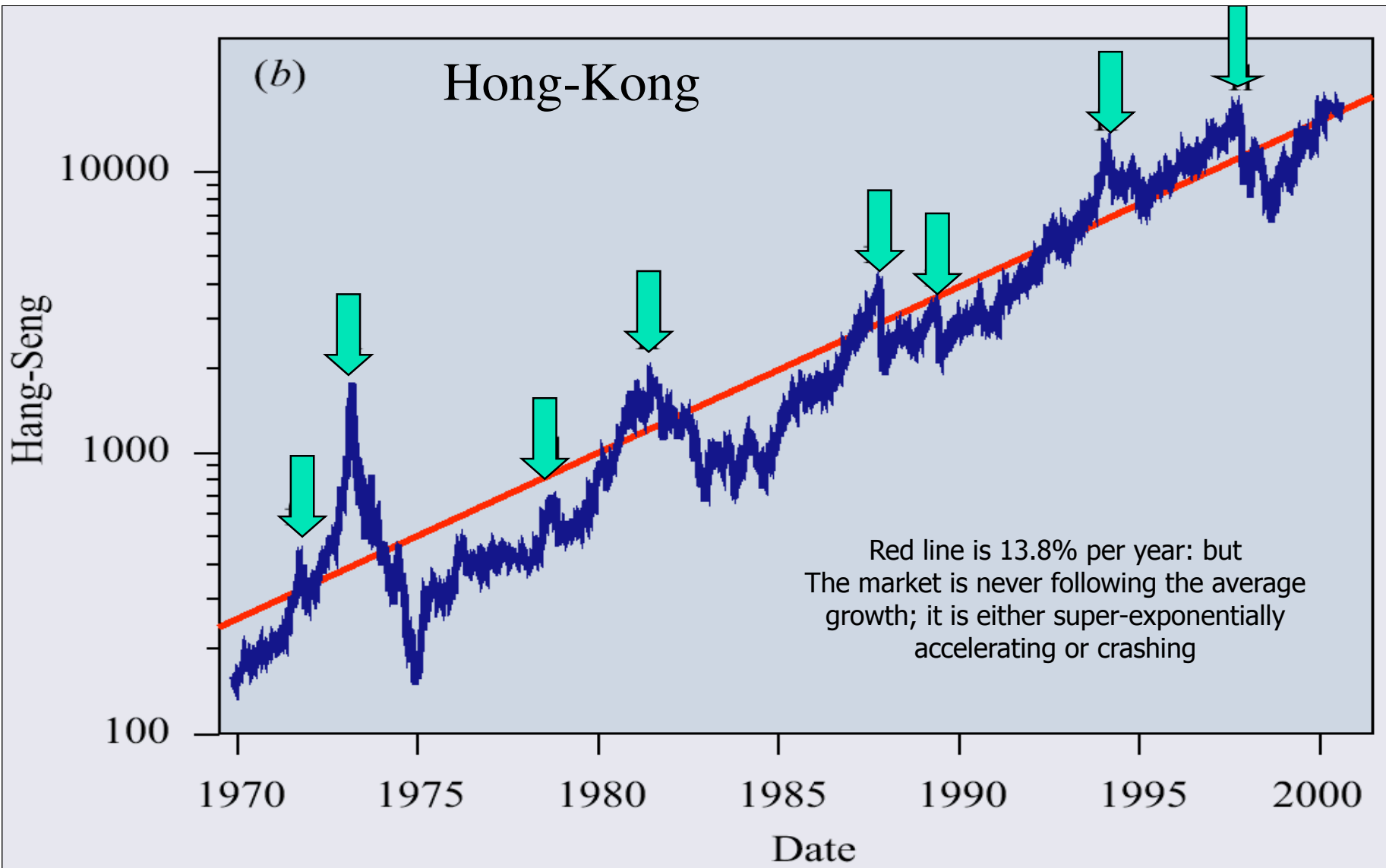


Securitization leads to larger inter-connectivity



Coupling strength increases





Patterns of price trajectory during 0.5-1 year before each peak: Log-periodic power law

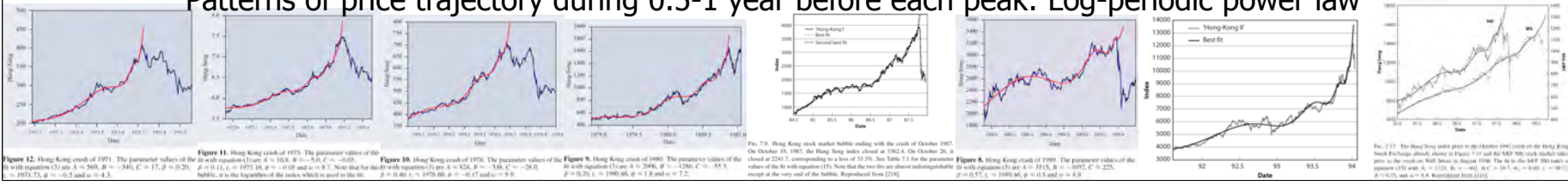


Figure 12. Hong Kong crash of 1971. The parameter values of the fit with equation (1) are:  $A = 560$ ,  $B = -340$ ,  $C = 17$ ,  $\beta = 0.20$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ .  
 Figure 13. Hong Kong crash of 1975. The parameter values of the fit with equation (1) are:  $A = 10.0$ ,  $B = -5.0$ ,  $C = -0.03$ ,  $\beta = 0.11$ ,  $\alpha = 0.75$ ,  $\delta = -0.02$ , and  $\omega = 8.7$ . Note that for the fit with equation (1) are:  $A = 524$ ,  $B = -350$ ,  $C = -200$ ,  $\beta = 0.20$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ .  
 Figure 14. Hong Kong crash of 1976. The parameter values of the fit with equation (1) are:  $A = 2000$ ,  $B = -1200$ ,  $C = -35$ ,  $\beta = 0.20$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ .  
 Figure 15. Hong Kong crash of 1980. The parameter values of the fit with equation (1) are:  $A = 2000$ ,  $B = -1200$ ,  $C = -35$ ,  $\beta = 0.20$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ .  
 Figure 16. Hong Kong stock market bubble ending with the crash of October 1987. On October 19, 1987, the Hong Kong index closed at 3324. On October 26, it closed at 2247, corresponding to a loss of 33.3%. See Table 7 for the parameter values of the fit with equation (1). Note that the two fits are almost indistinguishable except at the very end of the bubble. Reprinted from [216].  
 Figure 17. Hong Kong crash of 1989. The parameter values of the fit with equation (1) are:  $A = 2015$ ,  $B = -1072$ ,  $C = -225$ ,  $\beta = 0.21$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ .  
 Figure 18. Hong Kong stock market bubble ending with the crash of October 1997. The parameter values of the fit with equation (1) are:  $A = 1125$ ,  $B = -400$ ,  $C = -10$ ,  $\beta = 0.21$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ .  
 Figure 19. The Hong Kong index price in the October 1997 crash. The Hong Kong index price is shown in blue. The best fit is shown in red. The fit is the best fit with equation (1) with  $A = 1125$ ,  $B = -400$ ,  $C = -10$ ,  $\beta = 0.21$ ,  $\alpha = 0.72$ ,  $\delta = -0.2$ , and  $\omega = 4.2$ . Reprinted from [217].

# What is the cause of the crash?



- ✓ Proximate causes: many possibilities
- ✓ Fundamental cause: maturation towards an **instability**



An instability is characterized by

- large or diverging susceptibility to external perturbations or influences
- exponential growth of random perturbations leading to a change of regime, or selection of a new attractor of the dynamics.