

# **Ultrafast Laser Physics**

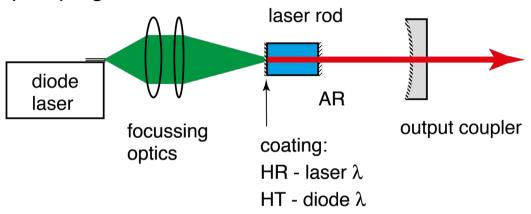
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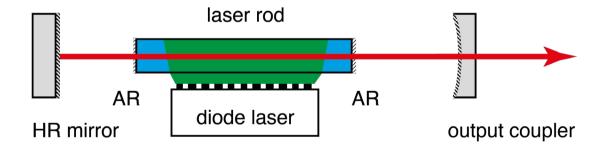
Chapter 5: Relaxation oscillations in lasers

## Diode-pumped solid-state lasers

#### Longitudinal pumping

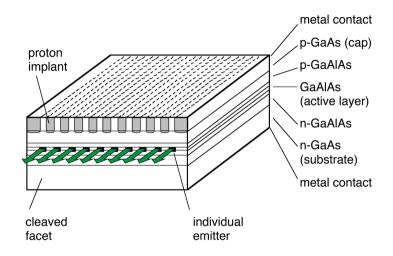


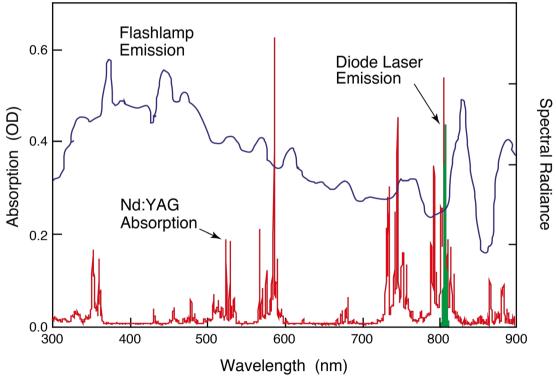
#### Transversal pumping



### **ETH**zürich

## Diode-pumped solid-state lasers

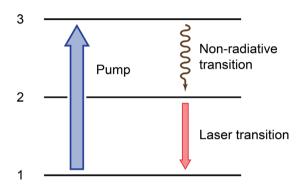


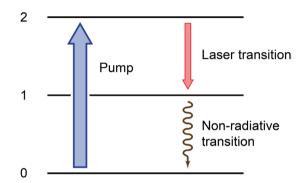


## Diode-pumped solid-state lasers

#### Three-level system

Yb-doped

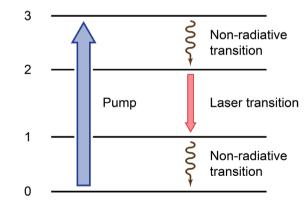




#### Four-level system

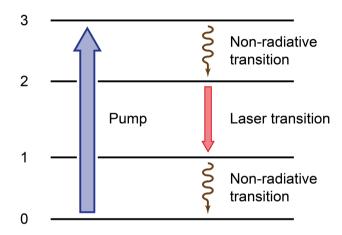
Nd-doped

Ti:sapphire



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### Laser rate equation



$$\frac{dn}{dt} = K(n+1)N - \gamma_c n$$

$$\frac{dN}{dt} = R_p - KnN - \gamma_L N$$

- · ideal four-level system
- single mode laser
- homogeneous broadened gain material

N inversion  $N = N_2 - N_1$ ,  $N_1 \approx 0$ 

*n* photon number (intracavity)

$$W^{stim} \equiv W_{12} = W_{21} = Kn$$

KN spontaneous emission rate

KnN stimulated emission rate

 $\gamma_c$  cavity photon decay rate

 $\gamma_L$  spontaneous decay rate of level 2

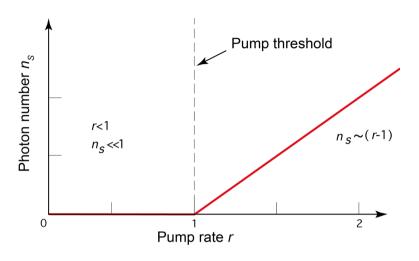
### Steady-state solution

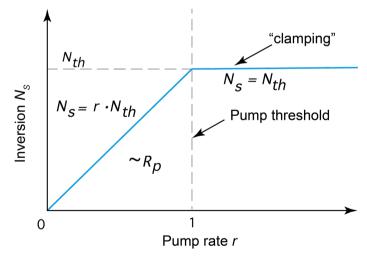
Above threshold, i.e. r > 1

$$\Leftrightarrow \begin{cases} n_s \approx \frac{\tau_c}{\tau_L} N_{th}(r-1) = \frac{\gamma_L}{K}(r-1) \\ N_s \approx N_{th} = \frac{\gamma_c}{K} \end{cases}$$

Below threshold, i.e. r < 1

$$\Leftrightarrow \begin{cases} n_s = \frac{r}{1-r} \\ N_s \approx r N_{th} \end{cases}$$





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### Linearized rate equations

$$n(t) = n_s + n_1(t),$$
  $n_1(t) \ll n_s$   
 $N(t) = N_s + N_1(t),$   $N_1(t) \ll N_s$ 

$$\frac{dn_1(t)}{dt} = \gamma_L(r-1)N_1(t)$$

$$\frac{dN_1(t)}{dt} = -\gamma_L r N_1(t) - \gamma_c n_1(t)$$

$$\frac{dn_s}{dt} = 0, \quad \frac{dN_s}{dt} = 0$$

Solution: Ansatz

$$n_1(t) \propto e^{st}, \quad N_1(t) \propto e^{st}$$

$$s = s_{1,2} = -\frac{r\gamma_L}{2} \pm \sqrt{\left(\frac{r\gamma_L}{2}\right)^2 - \gamma_L \gamma_c(r-1)}$$

## ETHzürich Over-critical damping (no relaxation oscillations)

$$n(t) = n_s + n_1(t),$$
  $n_1(t) \ll n_s$   
 $N(t) = N_s + N_1(t),$   $N_1(t) \ll N_s$ 

$$s = s_{1,2} = -\frac{r\gamma_L}{2} \pm \sqrt{\left(\frac{r\gamma_L}{2}\right)^2 - \gamma_L \gamma_c(r-1)}$$

Solution: Ansatz

$$n_1(t) \propto e^{st}, \quad N_1(t) \propto e^{st}$$

s is real and negative:  $\left(\frac{r\gamma_L}{2}\right)^2 > \gamma_L\gamma_c(r-1)$ 

two solutions: two over-critically damped relaxation constants:

$$r\gamma_L \gg \gamma_c \quad \Rightarrow \quad \begin{cases} s_1 = -r\gamma_L & \text{Stimulated decay rate of excited atoms} \\ s_2 = -\gamma_c \frac{r-1}{r} & \text{Photon decay rate inside laser cavity} \end{cases}$$

$$s_2 = -\gamma_c \frac{r-1}{r} \quad \Rightarrow \quad s_2 \approx -\gamma_c, \quad \text{for} \quad r \gg 1$$

### ETH zürich Over-critical damping (no relaxation oscillations)

Example HeNe Laser (p. 200)

 $\lambda = 632.8 \text{ nm}, \ \tau_{L} \approx 100 \text{ ns}, \ 2l = 0.02 \ (2\% \text{ output coupler}), \ T_{R} = 2 \text{ ns}$ 

$$\Rightarrow$$
  $\tau_c = \frac{T_R}{2l} \approx 100 \,\mathrm{ns} \quad \Rightarrow \quad \tau_L \approx \tau_c, \quad \mathrm{or} \quad \gamma_L \approx \gamma_c$ 

For r >> 1 (when HeNe is pumped sufficiently far above threshold) condition for over-critical damping is fulfilled:

$$r\gamma_L \gg \gamma_c$$

HeNe laser falls back into steady state after about 100 ns

## ETHzürich Under-critical damping (relaxation oscillations)

$$n(t) = n_s + n_1(t),$$
  $n_1(t) \ll n_s$   
 $N(t) = N_s + N_1(t),$   $N_1(t) \ll N_s$ 

$$s = s_{1,2} = -\frac{r\gamma_L}{2} \pm \sqrt{\left(\frac{r\gamma_L}{2}\right)^2 - \gamma_L \gamma_c(r-1)}$$

Solution: Ansatz

$$n_1(t) \propto e^{st}, \quad N_1(t) \propto e^{st}$$

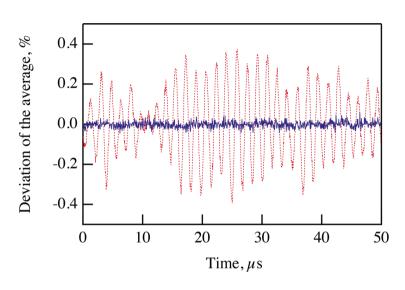
s is complex: 
$$\left(\frac{r\gamma_L}{2}\right)^2 < \gamma_L\gamma_c(r-1)$$

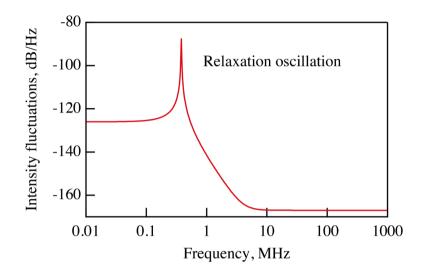
relaxation oscillations with an attenuation factor:

$$n(t) = n_s + n_1 e^{-\gamma_{relax}t} \cos(\omega_{relax}t)$$
  $\gamma_{relax} = \frac{r\gamma_L}{2}$ 

$$\gamma_c \gg r\gamma_L: \quad \omega_{relax} \approx \sqrt{\frac{r-1}{\tau_L} \frac{1}{\tau_c}} = \sqrt{\frac{1}{\tau_{stim}} \frac{1}{\tau_c}}$$

#### Relaxation oscillations





Relaxation oscillations in the time domain and frequency domain (microwave spectrum analyzer - not optical frequencies!)

$$n(t) = n_s + n_1 e^{-\gamma_{relax}t} \cos(\omega_{relax}t)$$

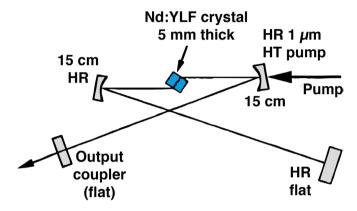
$$\gamma_{relax} = \frac{r\gamma_L}{2}$$

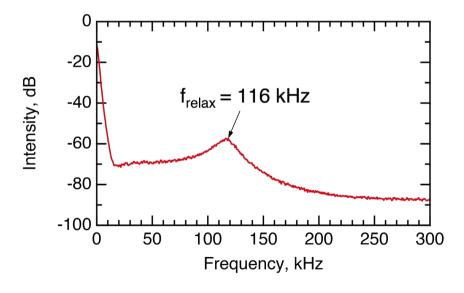
$$\gamma_c \gg r \gamma_L : \quad \omega_{relax} \approx \sqrt{\frac{r-1}{\tau_L} \frac{1}{\tau_c}} = \sqrt{\frac{1}{\tau_{stim}} \frac{1}{\tau_c}}$$

### Relaxation oscillations

Example:

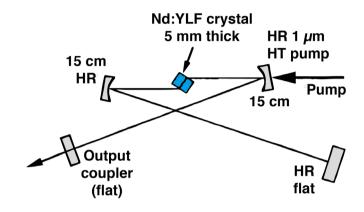
diode-pumped cw Nd:YLF laser





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### Measurement of small signal gain



$$g = \frac{g_0}{1 + 2I/I_{sat}}$$

$$f_{relax} \approx \frac{1}{2\pi} \sqrt{\gamma_L \gamma_c(r-1)} = \frac{1}{2\pi} \sqrt{\frac{r-1}{\tau_L \tau_c}} = \frac{1}{2\pi} \sqrt{\frac{\frac{g_0}{l} - 1}{\tau_L \frac{T_R}{2l}}} = \frac{1}{2\pi} \sqrt{\frac{2g_0 - 2l}{\tau_L T_R}}$$

$$\tau_c = \frac{T_R}{2l}$$

$$2g_0 \approx 4\pi^2 \tau_L T_R f_{relax}^2 + 2l$$

$$r = \frac{g_0}{l}$$

$$G_0 \approx e^{2g_0} = \exp\left(4\pi^2 \tau_L T_R f_{relax}^2 + 2l\right)$$

### Measurement of small signal gain

Measured relaxation oscillationfrequency for different outputcouplers, for  $g_0 \gg l$ 

$$f_{relax} \approx \frac{1}{2\pi} \sqrt{\frac{2g_0 - 2l}{\tau_L T_R}} \xrightarrow{g_0 \gg l} \approx \frac{1}{2\pi} \sqrt{\frac{2g_0}{\tau_L T_R}}$$

relaxation oscillations independent of output coupler

$$2g_0 \approx 4\pi^2 \tau_L T_R f_{relax}^2 + 2l$$

$$G_0 \approx e^{2g_0} = \exp\left(4\pi^2 \tau_L T_R f_{relax}^2 + 2l\right)$$

Example: diode-pumped Nd:YLF laser

