

Ultrafast Laser Physics

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Chapter 5: Relaxation oscillations in lasers

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Diode-pumped solid-state lasers

Transversal pumping

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Diode-pumped solid-state lasers

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Diode-pumped solid-state lasers

Three-level system

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nm

Laser rate equation

$$
\frac{dn}{dt} = K(n+1)N - \gamma_c n
$$

$$
\frac{dN}{dt} = R_p - KnN - \gamma_L N
$$

- ideal four-level system
- single mode laser

• homogeneous broadened gain material *N* inversion $N = N_2 - N_1$, $N_1 \approx 0$ *n* photon number (intracavity) $W^{stim} \equiv W_{12} = W_{21} = Kn$ *KN* spontaneous emission rate *KnN* stimulated emission rate γ_c cavity photon decay rate γ_L spontaneous decay rate of level 2

Steady-state solution

Pump threshold Photon number n_s Above threshold, i.e. $r > 1$ $\Leftrightarrow \begin{cases} n_s \approx \frac{\tau_c}{\tau_L} N_{th}(r-1) = \frac{\gamma_L}{K}(r-1) \\ N_s \approx N_{th} = \frac{\gamma_c}{K} \end{cases}$ $r<1$ $n_s{\sim}(\mathit{r}\text{-}1)$ n_s <<1 $\overline{0}$ $\overline{2}$ Pump rate r "clamping" N_{th} Below threshold, i.e. $r < 1$ $N_s = N_{th}$ Inversion $N_{\rm s}$ $\Leftrightarrow \begin{cases} n_s = \frac{r}{1-r} \\ N_s \approx r N_{th} \end{cases}$ $N_s = r \cdot N_{th}$ Pump threshold $\sim R_D$ $\mathbf 0$ $\mathbf{1}$ Pump rate r **ETH**zürich ∧M∿ Ultrafast Lasers book, Subsection 5.1.3

Linearized rate equations

$$
n(t) = n_s + n_1(t), \qquad n_1(t) \ll n_s
$$

$$
N(t) = N_s + N_1(t), \quad N_1(t) \ll N_s
$$

$$
\frac{dn_1(t)}{dt} = \gamma_L(r-1)N_1(t)
$$

$$
\frac{dN_1(t)}{dt} = -\gamma_L r N_1(t) - \gamma_c n_1(t)
$$

$$
\frac{dn_s}{dt} = 0, \quad \frac{dN_s}{dt} = 0
$$

Solution: Ansatz

$$
n_1(t) \propto e^{st}
$$
, $N_1(t) \propto e^{st}$

$$
s = s_{1,2} = -\frac{r\gamma_L}{2} \pm \sqrt{\left(\frac{r\gamma_L}{2}\right)^2 - \gamma_L \gamma_c (r-1)}
$$

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ETH züric/Over-critical damping (no relaxation oscillations)

 $n(t) = n_s + n_1(t), \quad n_1(t) \ll n_s$ $N(t) = N_s + N_1(t), N_1(t) \ll N_s$ Solution: Ansatz

$$
n_1(t) \propto e^{st}, \quad N_1(t) \propto e^{st}
$$

$$
s = s_{1,2} = -\frac{r\gamma_L}{2} \pm \sqrt{\left(\frac{r\gamma_L}{2}\right)^2 - \gamma_L \gamma_c (r-1)}
$$

s is real and negative: $\begin{pmatrix} 7 \\ -1 \end{pmatrix}$

$$
\frac{r\gamma_L}{2}\Big)^2>\gamma_L\gamma_c(r-1)
$$

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two solutions: two over-critically damped relaxation constants:

 $r\gamma_L \gg \gamma_c \Rightarrow \begin{cases} s_1 = -r\gamma_L & \text{Stimulated decay rate of excited atoms} \ s_2 = -\gamma_c \frac{r-1}{r} & \text{Photon decay rate inside laser cavity} \end{cases}$ $s_2 = -\gamma_c \frac{r-1}{r}$ \Rightarrow $s_2 \approx -\gamma_c$, for $r \gg 1$

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Ultrafast Lasers book, Subsection 5.2.3

ETH züric/Over-critical damping (no relaxation oscillations)

Example HeNe Laser (p. 200)

 λ = 632.8 nm, $\tau_{\text{I}} \approx 100$ ns, 2*l* = 0.02 (2% output coupler), T_{R} = 2 ns

$$
\Rightarrow \quad \tau_c = \frac{T_R}{2l} \approx 100 \,\text{ns} \quad \Rightarrow \quad \tau_L \approx \tau_c, \quad \text{or} \quad \gamma_L \approx \gamma_c
$$

For *r* >>1 (when HeNe is pumped sufficiently far above threshold) condition for over-critical damping is fulfilled:

 $r\gamma_L \gg \gamma_c$

HeNe laser falls back into steady state after about 100 ns

ETH zürich Under-critical damping (relaxation oscillations)

 $n(t) = n_s + n_1(t),$ $n_1(t) \ll n_s$
 $N(t) = N_s + N_1(t),$ $N_1(t) \ll N_s$

Solution: Ansatz

$$
n_1(t) \propto e^{st}, \quad N_1(t) \propto e^{st}
$$

$$
s = s_{1,2} = -\frac{r\gamma_L}{2} \pm \sqrt{\left(\frac{r\gamma_L}{2}\right)^2 - \gamma_L \gamma_c (r-1)}
$$

s is complex:
$$
\begin{pmatrix} T \\ - \end{pmatrix}
$$

$$
\left(\frac{r\gamma_L}{2}\right)^2 < \gamma_L\gamma_c(r-1)
$$

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relaxation oscillations with an attenuation factor:

$$
n(t) = n_s + n_1 e^{-\gamma_{relax}t} \cos(\omega_{relax}t) \qquad \boxed{\gamma_{relax}}
$$

$$
\gamma_c \gg r\gamma_L: \quad \omega_{relax} \approx \sqrt{\frac{r-1}{\tau_L}} \frac{1}{\tau_c} = \sqrt{\frac{1}{\tau_{stim}} \frac{1}{\tau_c}}
$$

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Relaxation oscillations in the time domain and frequency domain (microwave spectrum analyzer - not optical frequencies!)

$$
n(t) = n_s + n_1 e^{-\gamma_{relax}t} \cos(\omega_{relax}t)
$$

$$
\gamma_{relax} = \frac{r\gamma_L}{2}
$$

$$
\gamma_c \gg r\gamma_L: \quad \omega_{relax} \approx \sqrt{\frac{r-1}{\tau_L} \frac{1}{\tau_c}} = \sqrt{\frac{1}{\tau_{stim}} \frac{1}{\tau_c}}
$$

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Measurement of small signal gain

$$
f_{relax} \approx \frac{1}{2\pi} \sqrt{\gamma_L \gamma_c (r-1)} = \frac{1}{2\pi} \sqrt{\frac{r-1}{\tau_L \tau_c}} = \frac{1}{2\pi} \sqrt{\frac{\frac{g_0}{l} - 1}{\tau_L \frac{T_R}{2l}}} = \frac{1}{2\pi} \sqrt{\frac{2g_0 - 2l}{\tau_L T_R}}
$$

$$
\tau_c = \frac{T_R}{2l}
$$
\n
$$
2g_0 \approx 4\pi^2 \tau_L T_R f_{relax}^2 + 2l
$$
\n
$$
r = \frac{g_0}{l}
$$
\n
$$
G_0 \approx e^{2g_0} = \exp\left(4\pi^2 \tau_L T_R f_{relax}^2 + 2l\right)
$$

nm

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Measurement of small signal gain

Measured relaxation oscillationfrequency for different outputcouplers, for $g_0 \gg l$

$$
f_{relax} \approx \frac{1}{2\pi} \sqrt{\frac{2g_0 - 2l}{\tau_L T_R}} \xrightarrow{g_0 \gg l} \approx \frac{1}{2\pi} \sqrt{\frac{2g_0}{\tau_L T_R}}
$$

relaxation oscillations independent of output coupler

$$
2g_0 \approx 4\pi^2 \tau_L T_R f_{relax}^2 + 2l
$$

$$
G_0 \approx e^{2g_0} = \exp\left(4\pi^2 \tau_L T_R f_{relax}^2 + 2l\right)
$$

Example: diode-pumped Nd:YLF laser

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